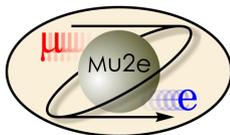


The Mu2e Experiment

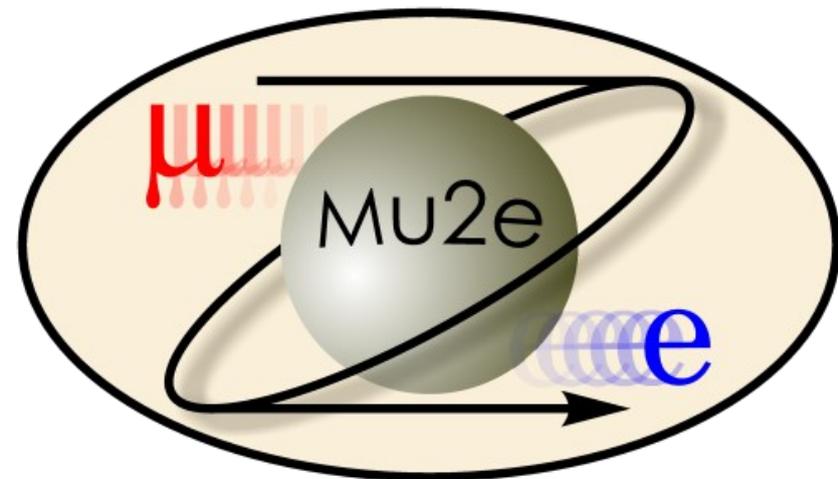
A High-sensitivity Charged Lepton Flavor Violating Search at Fermilab

For the Mu2e Collaboration
Kevin Lynch
Boston University
INPC, Vancouver, BC, Canada
July 4-9, 2010

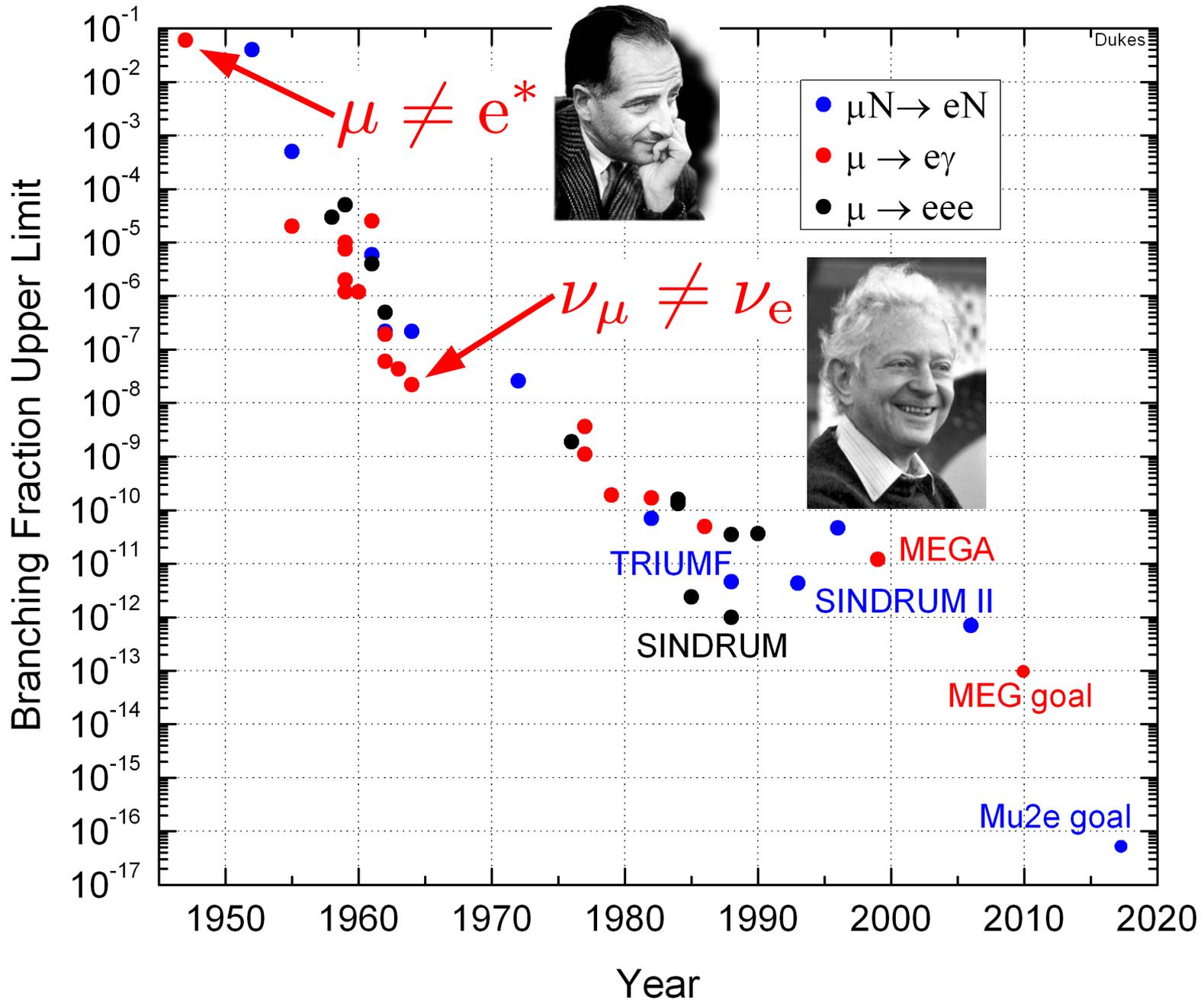


A brief roadmap to this talk...

- Lepton Flavor Violation
- Primary backgrounds to muon conversion
- The design of Mu2e



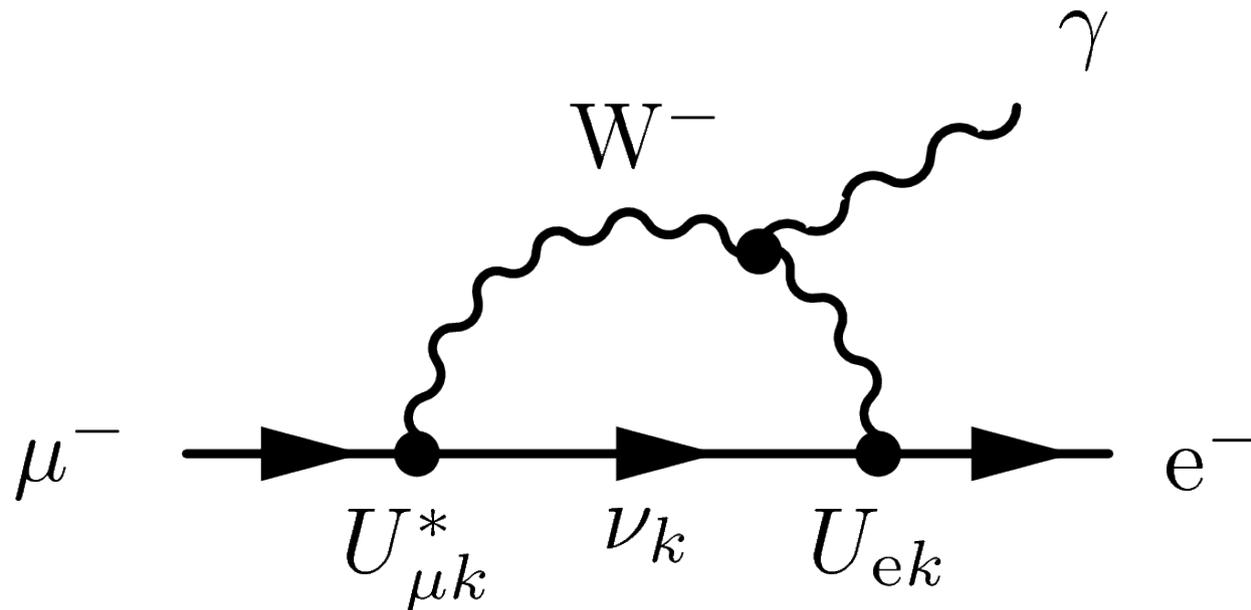
The mystery of lepton flavor has been with us for many decades



While charged lepton flavor violation has never been observed ...

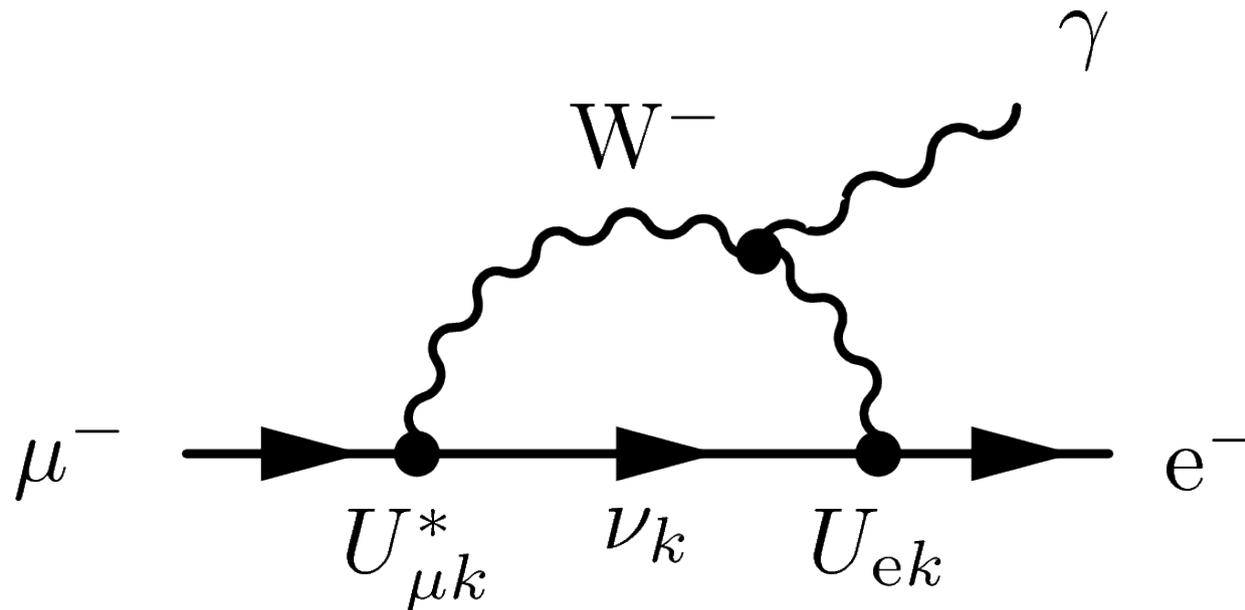
While charged lepton flavor violation has never been observed ...

... today we know it must occur, *even in the Standard Model*, through neutrino loop effects.



While charged lepton flavor violation has never been observed ...

... today we know it must occur, *even in the Standard Model*, through neutrino loop effects.



However, the predicted SM rates are unobservably small:

$$\text{Br}(\mu \rightarrow e\gamma) = \frac{3\alpha}{32\pi} \left| \sum_{k=2,3} U_{\mu k}^* U_{ek} \frac{\Delta m_{1k}^2}{M_W^2} \right|^2 < 10^{-54}$$

Any observation of cLFV is a direct signal of new physics!

Any observation of cLFV is a direct signal of new physics!

$$\mu^{\pm} \rightarrow e^{\pm} \gamma$$

$$\mu^{\pm} \rightarrow e^{\pm} e^{+} e^{-}$$

$$\mu^{-} A(Z, N) \rightarrow e^{-} A(Z, N)$$

Any observation of cLFV is a direct signal of new physics!

$$\mu^{\pm} \rightarrow e^{\pm} \gamma$$

$$\mu^{\pm} \rightarrow e^{\pm} e^{+} e^{-}$$

$$\mu^{-} A(Z, N) \rightarrow e^{-} A(Z, N)$$

The signal kinematics lie in a sea of background from ordinary muon decay.

Any observation of cLFV is a direct signal of new physics!

$$\mu^{\pm} \rightarrow e^{\pm} \gamma$$

$$\mu^{\pm} \rightarrow e^{\pm} e^{+} e^{-}$$

$$\mu^{-} A(Z, N) \rightarrow e^{-} A(Z, N)$$

The signal kinematics lie in a sea of background from ordinary muon decay.

There is no such kinematic limitation for the conversion process.

Any observation of cLFV is a direct signal of new physics!

$$\mu^{\pm} \rightarrow e^{\pm} \gamma$$

MEG at PSI

$$\mu^{\pm} \rightarrow e^{\pm} e^{+} e^{-}$$

The signal kinematics lie in a sea of background from ordinary muon decay.

$$\mu^{-} A(Z, N) \rightarrow e^{-} A(Z, N)$$

COMET/PRISM at JPARC

Mu2e at FNAL

There is no such kinematic limitation for the conversion process.

Any observation of cLFV is a direct signal of new physics!

$$\mu^\pm \rightarrow e^\pm \gamma$$

MEG at PSI

$$\mu^\pm \rightarrow e^\pm e^+ e^-$$

The signal kinematics lie in a sea of background from ordinary muon decay.

$$\mu^- A(Z, N) \rightarrow e^- A(Z, N)$$

COMET/PRISM at JPARC

Mu2e at FNAL

There is no such kinematic limitation for the conversion process.

The discovery of Weak scale SUSY at LHC would imply observable cLFV rates

$$R_{\mu e} \sim 10^{-15}$$

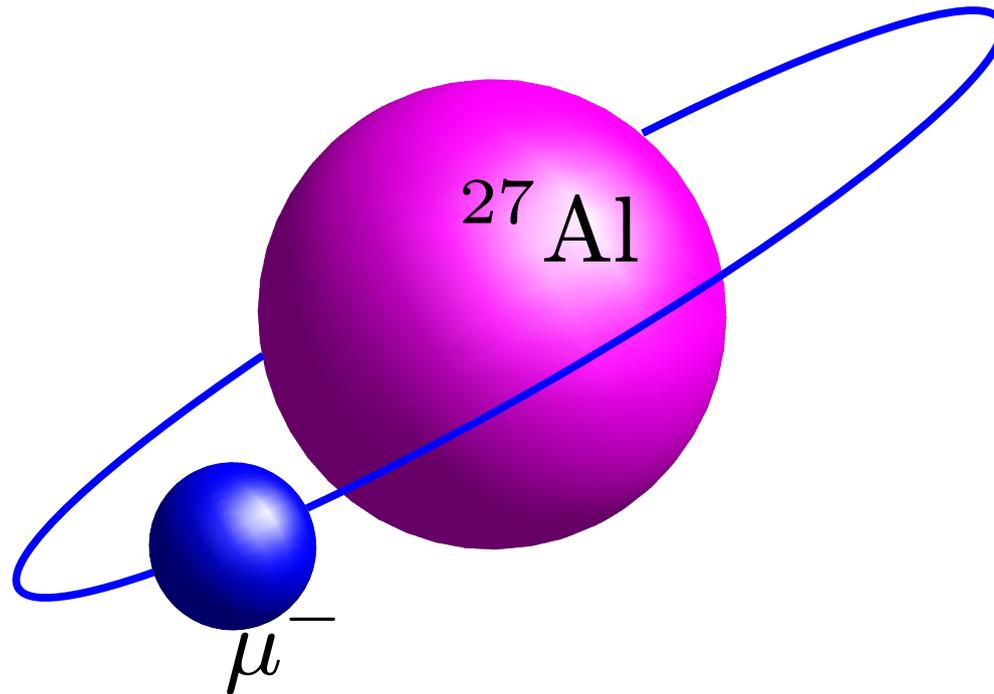
Coherent conversion kinematics

Negative muons rapidly stop, capture on atoms, and cascade to the 1s state...

Coherent conversion kinematics

Negative muons rapidly stop, capture on atoms, and cascade to the 1s state...

$$E_{1s} = -Z^2 2.8 \text{ keV}$$

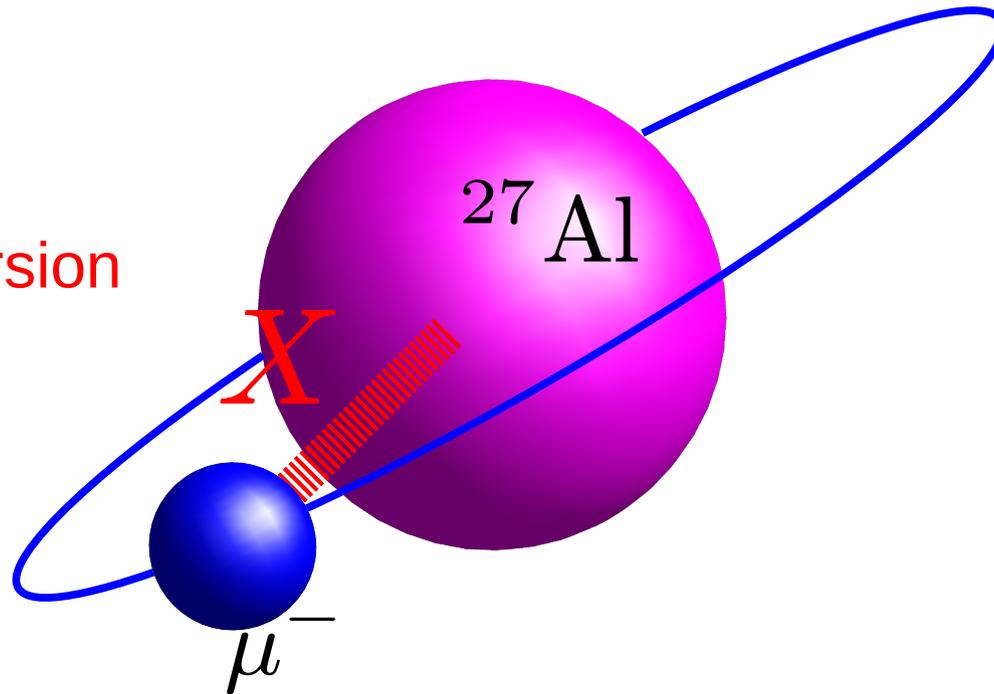


Coherent conversion kinematics

Negative muons rapidly stop, capture on atoms, and cascade to the 1s state...

$$E_{1s} = -Z^2 2.8 \text{ keV}$$

...then conversion happens...

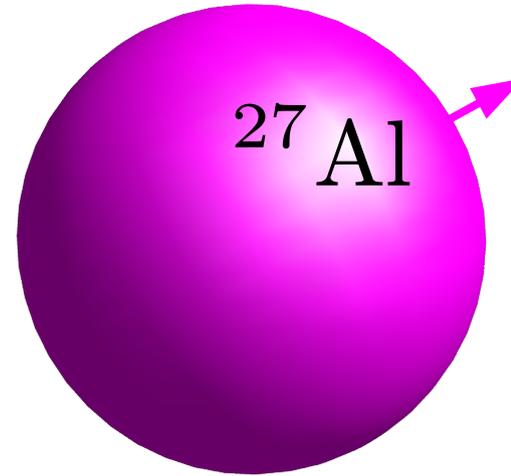
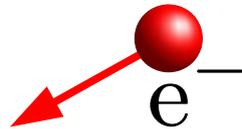


Coherent conversion kinematics

Negative muons rapidly stop, capture on atoms, and cascade to the 1s state...

$$E_{1s} = -Z^2 2.8 \text{ keV}$$

...then conversion happens...



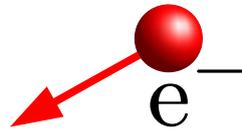
... a 1-to-2 process producing monochromatic electrons!

Coherent conversion kinematics

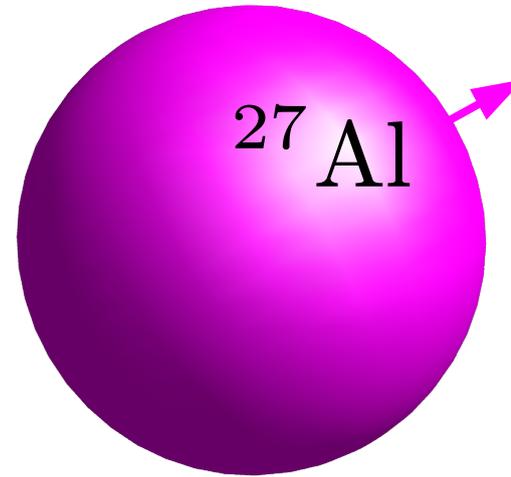
Negative muons rapidly stop, capture on atoms, and cascade to the 1s state...

$$E_{1s} = -Z^2 2.8 \text{ keV}$$

...then conversion happens...



... a 1-to-2 process producing monochromatic electrons!

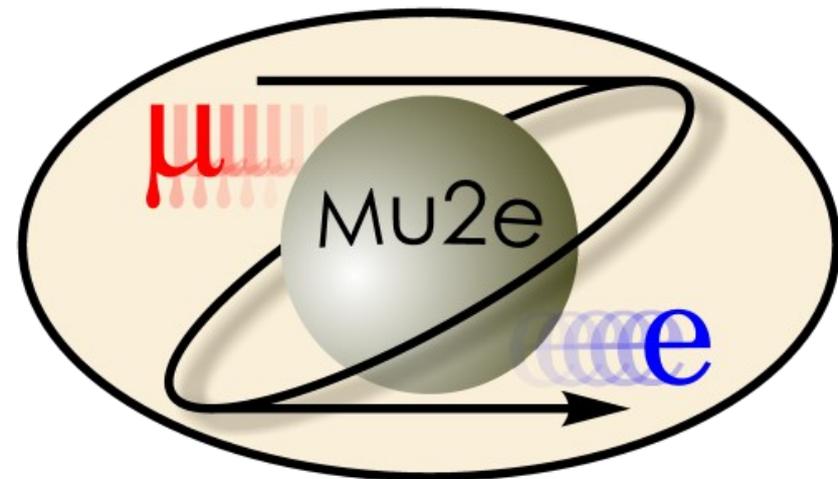


$$E_e = m_\mu - E_{1s} - E_{A(Z,N)}$$

$$E_e(^{27}\text{Al}) = 104.97 \text{ MeV}$$

A brief roadmap to this talk...

- Lepton Flavor Violation
- **Primary backgrounds to muon conversion**
- The design of Mu2e

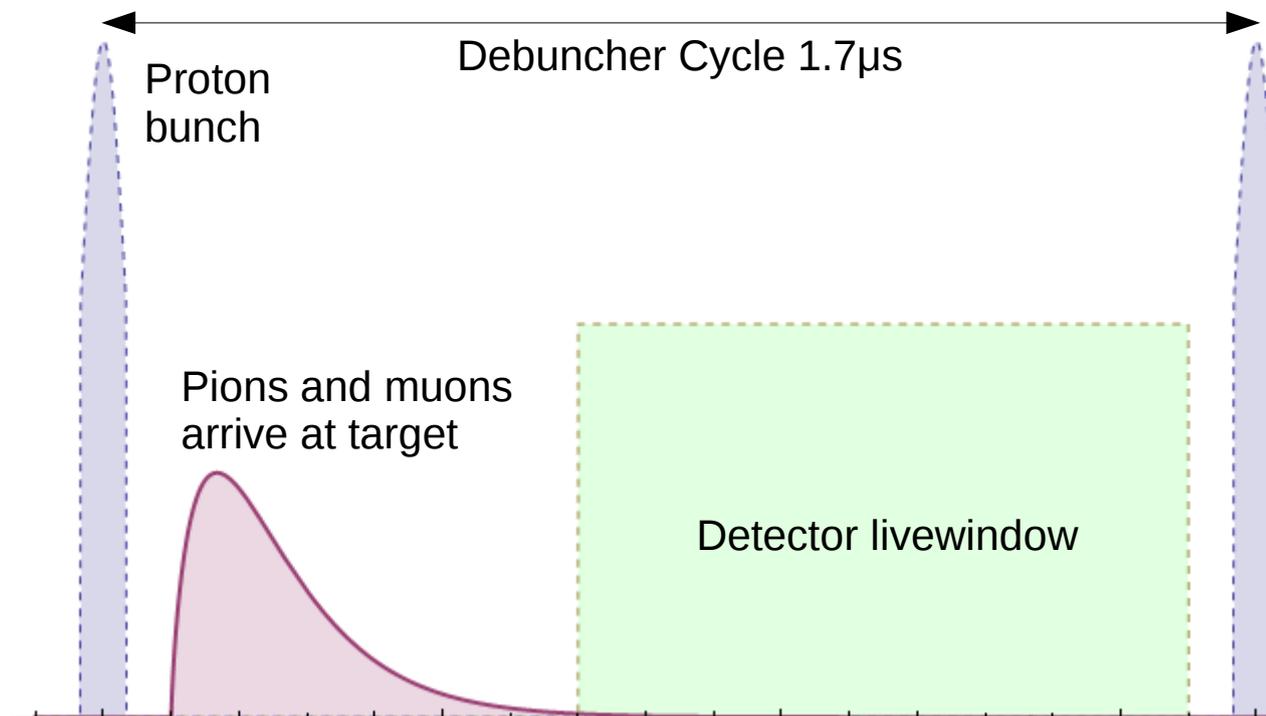


The two most dangerous backgrounds have very different timing properties at FNAL.

The FNAL accelerator complex produces proton beams with a pulsed structure.

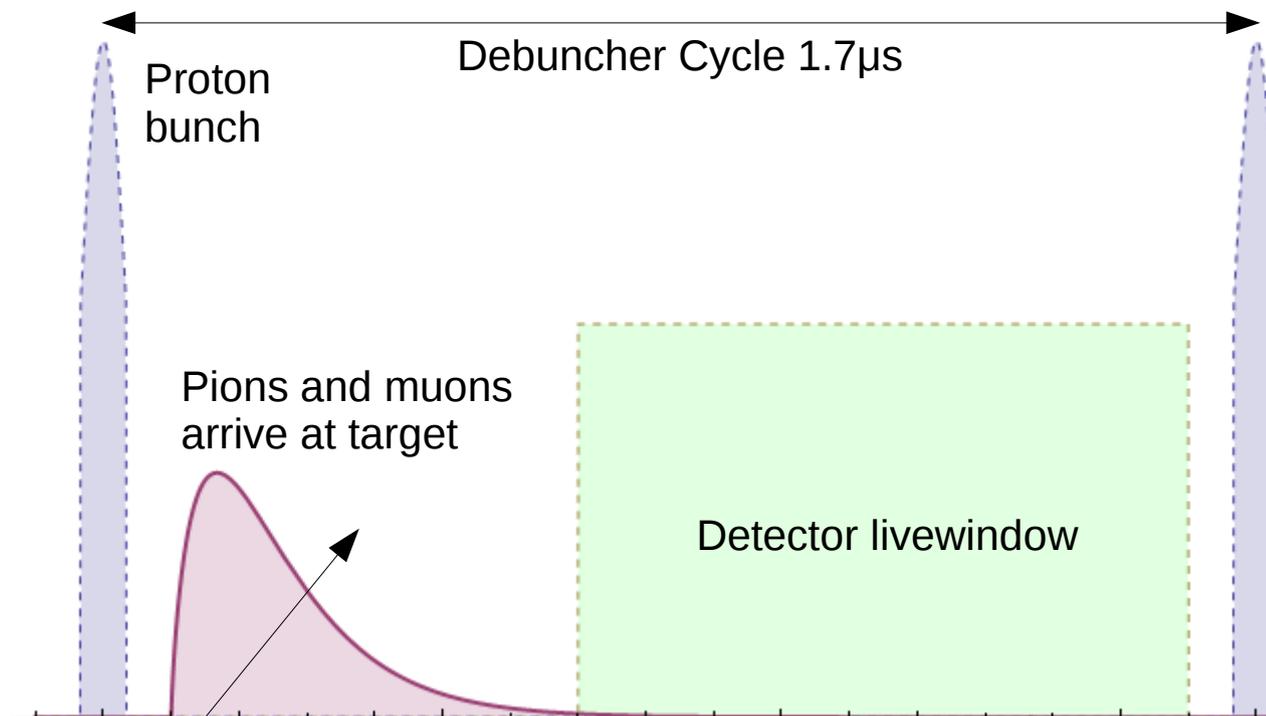
The two most dangerous backgrounds have very different timing properties at FNAL.

The FNAL accelerator complex produces proton beams with a pulsed structure.



The two most dangerous backgrounds have very different timing properties at FNAL.

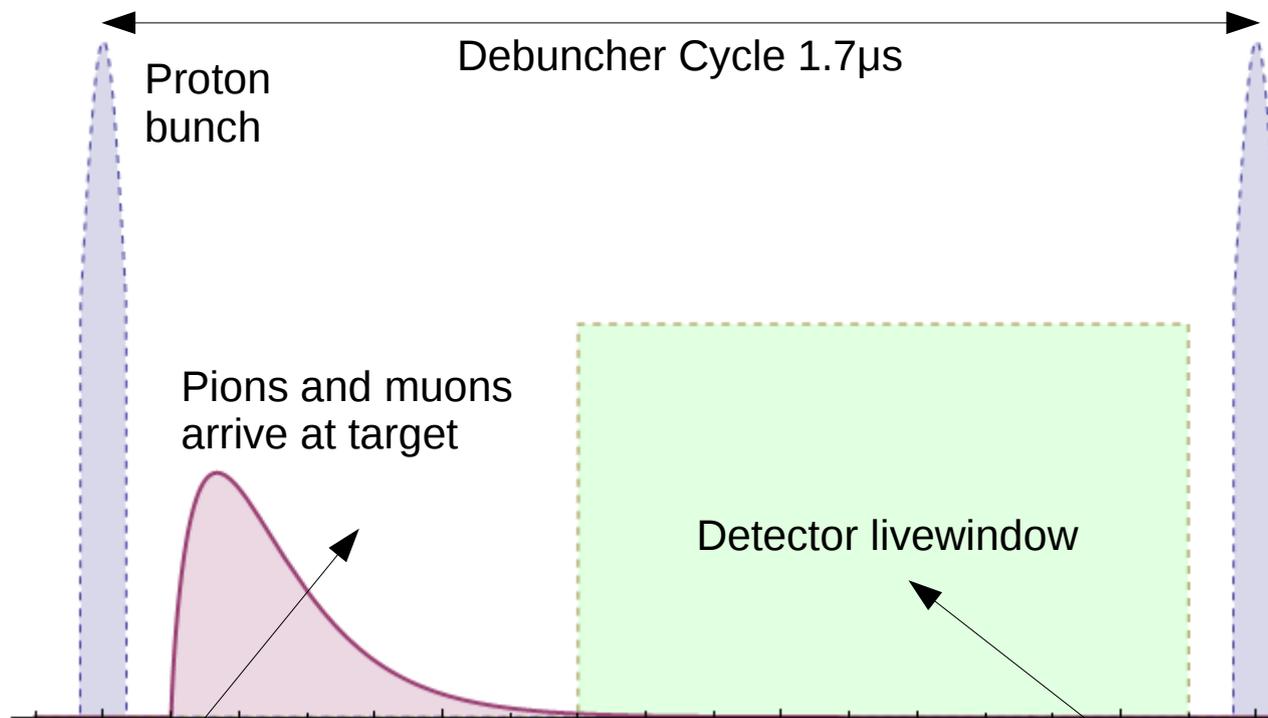
The FNAL accelerator complex produces proton beams with a pulsed structure.



Prompt: Radiative Pion Capture with pair production

The two most dangerous backgrounds have very different timing properties at FNAL.

The FNAL accelerator complex produces proton beams with a pulsed structure.



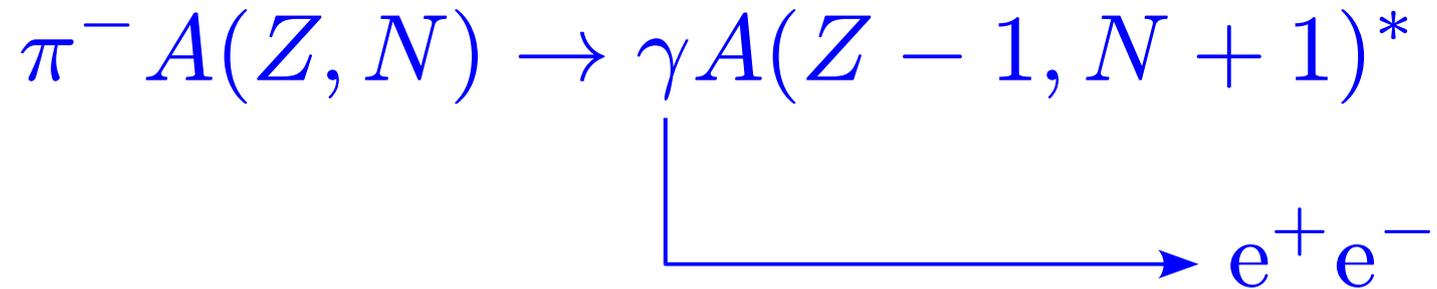
Prompt: Radiative Pion Capture with pair production

Delayed: Muon Decay-in-Orbit

Radiative pion capture can produce electrons near the conversion energy

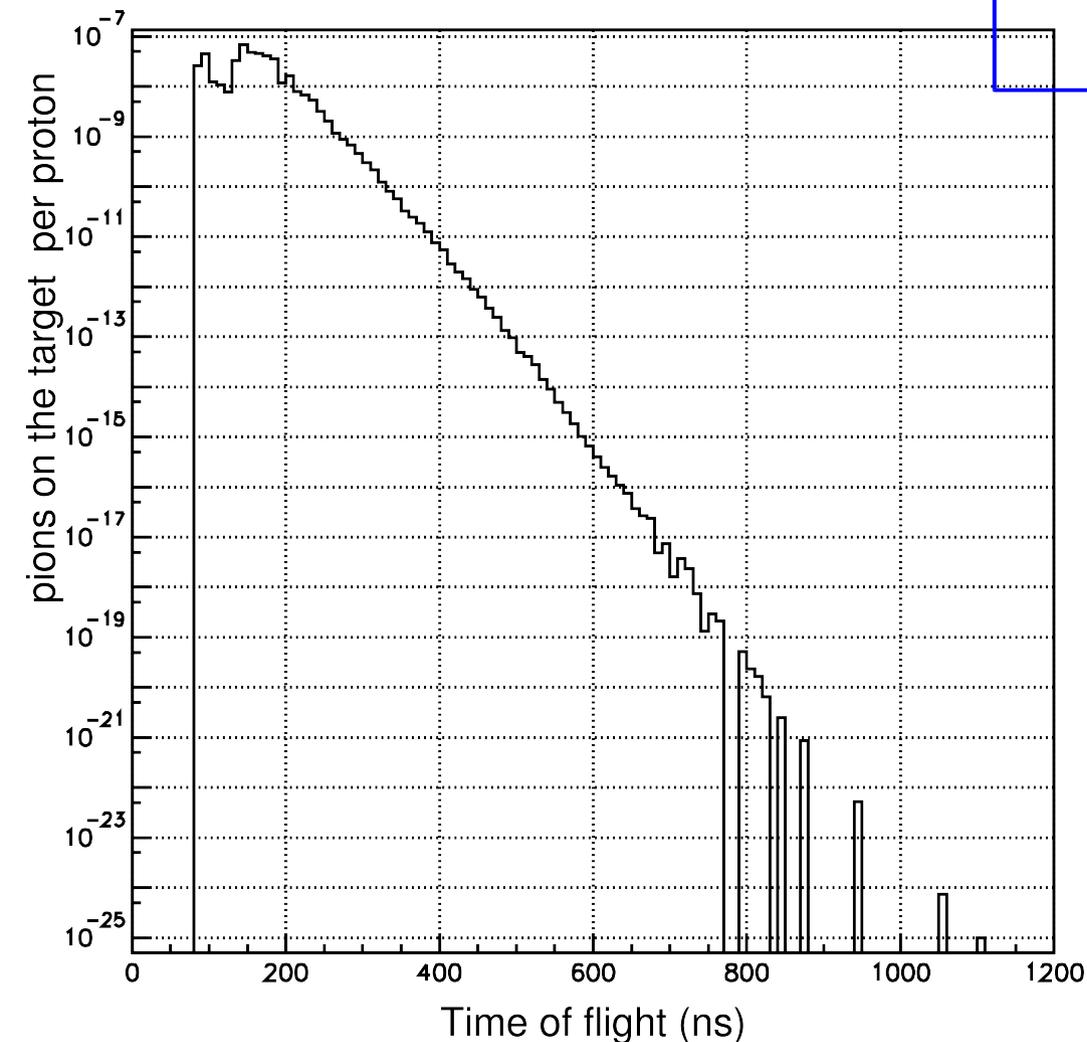


Radiative pion capture can produce electrons near the conversion energy



The electron spectrum extends out to nearly m_{π}

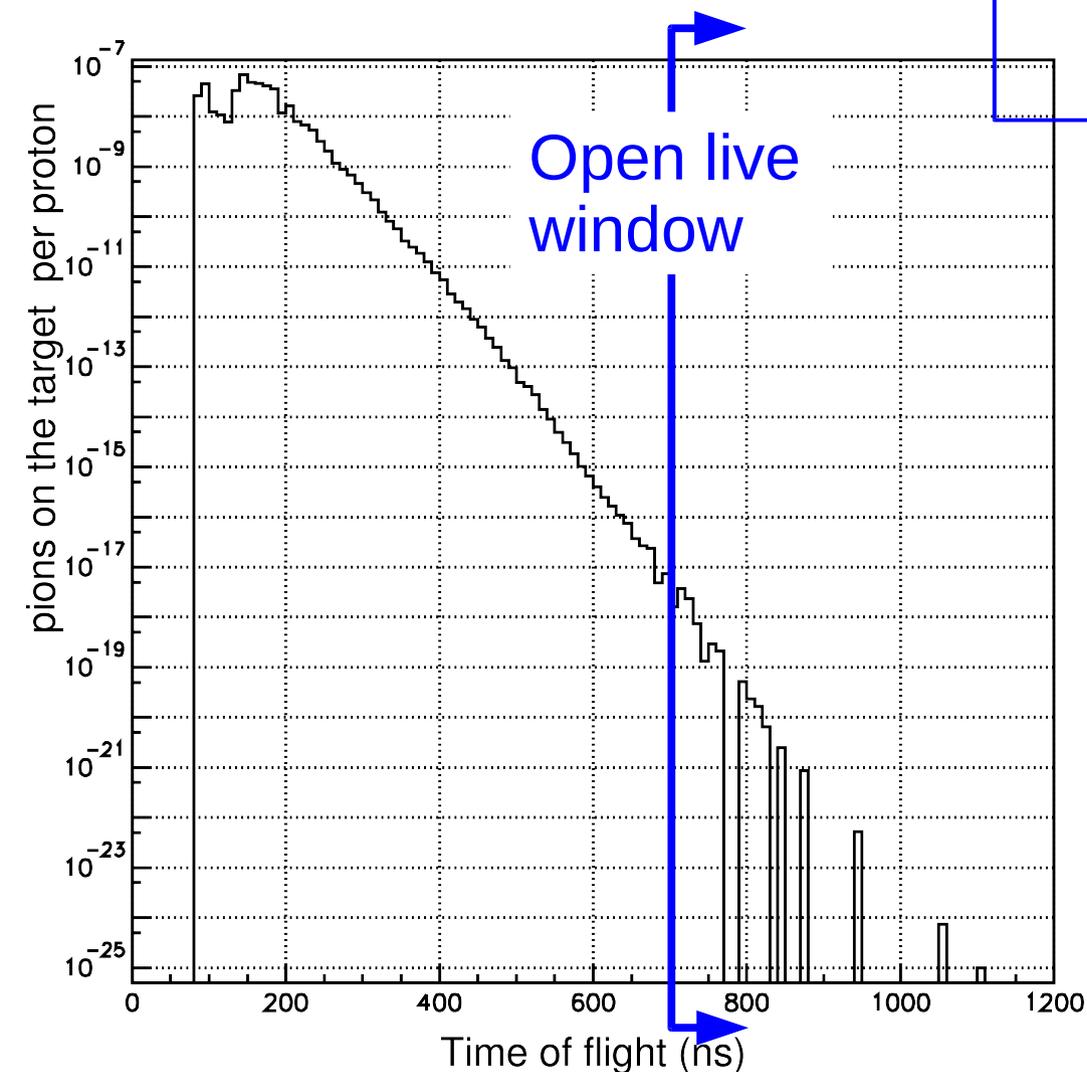
Radiative pion capture can produce electrons near the conversion energy



$e^+ e^-$

The electron spectrum extends out to nearly m_π

Radiative pion capture can produce electrons near the conversion energy

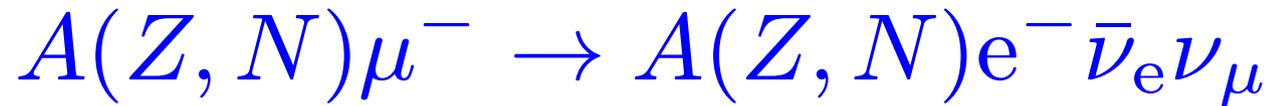


$e^+ e^-$

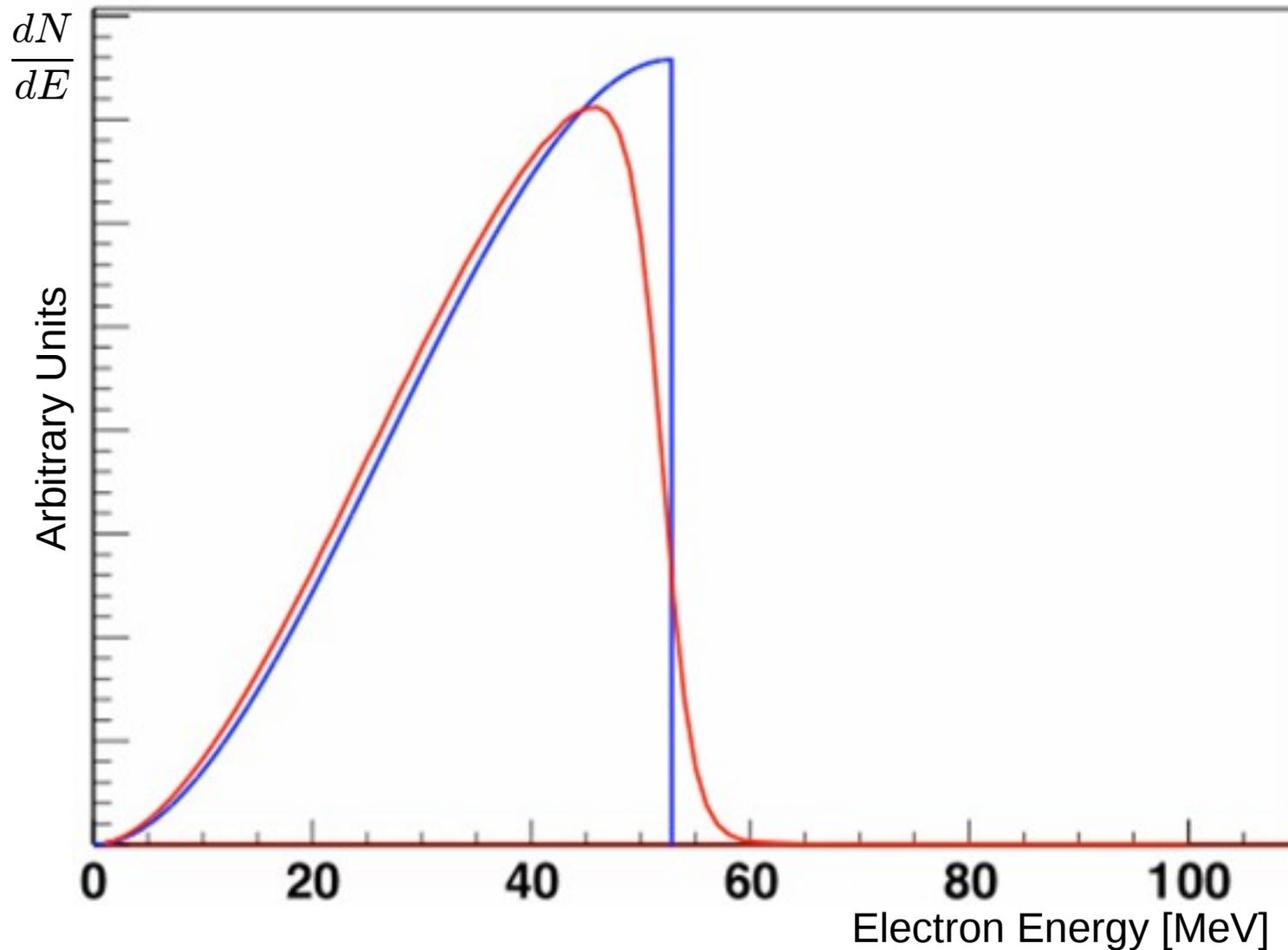
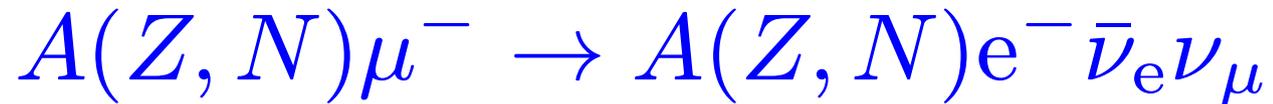
The electron spectrum extends out to nearly m_π

This (and other) prompt backgrounds are eliminated by keeping the DAQ dead until the background rate has fallen below some threshold.

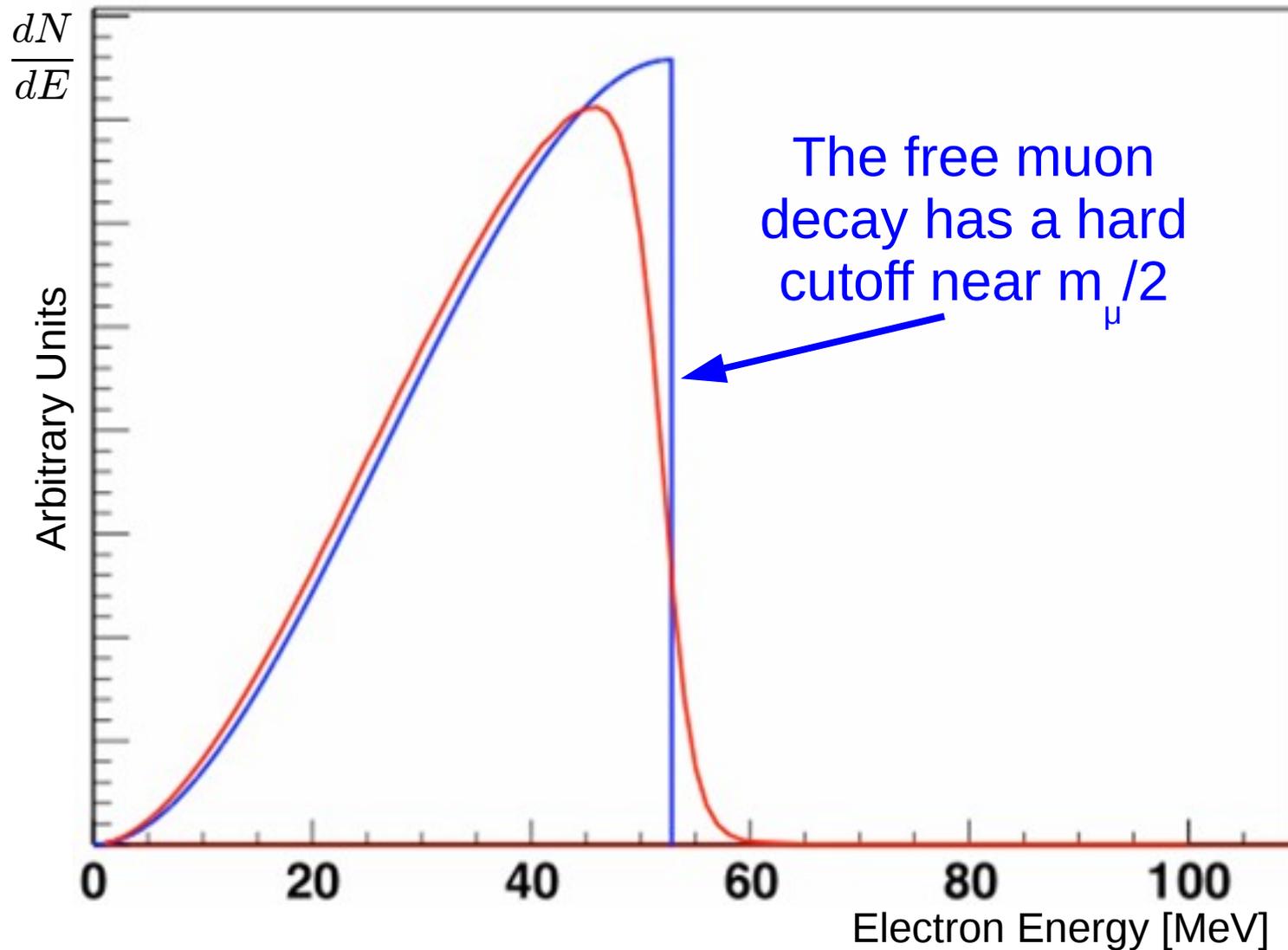
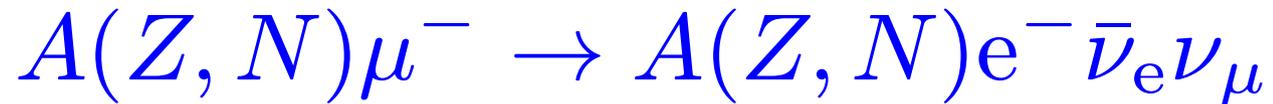
Decay-in-Orbit is the major source of delayed background in the live window



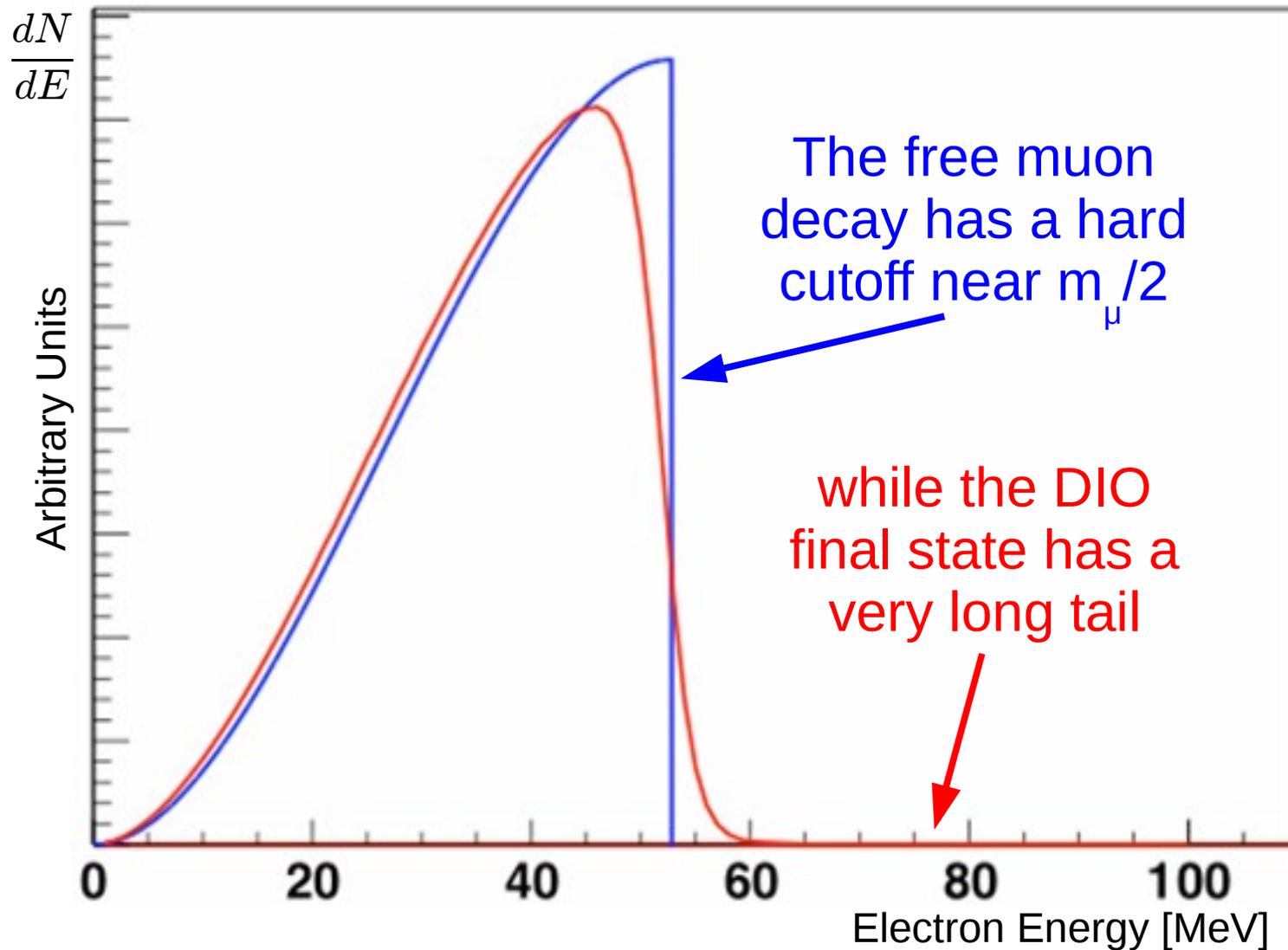
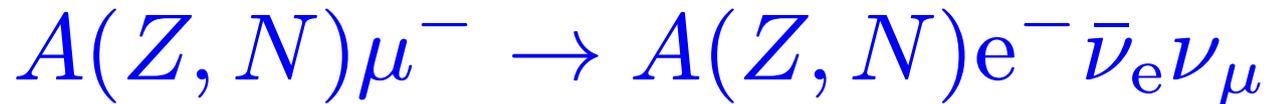
Decay-in-Orbit is the major source of delayed background in the live window



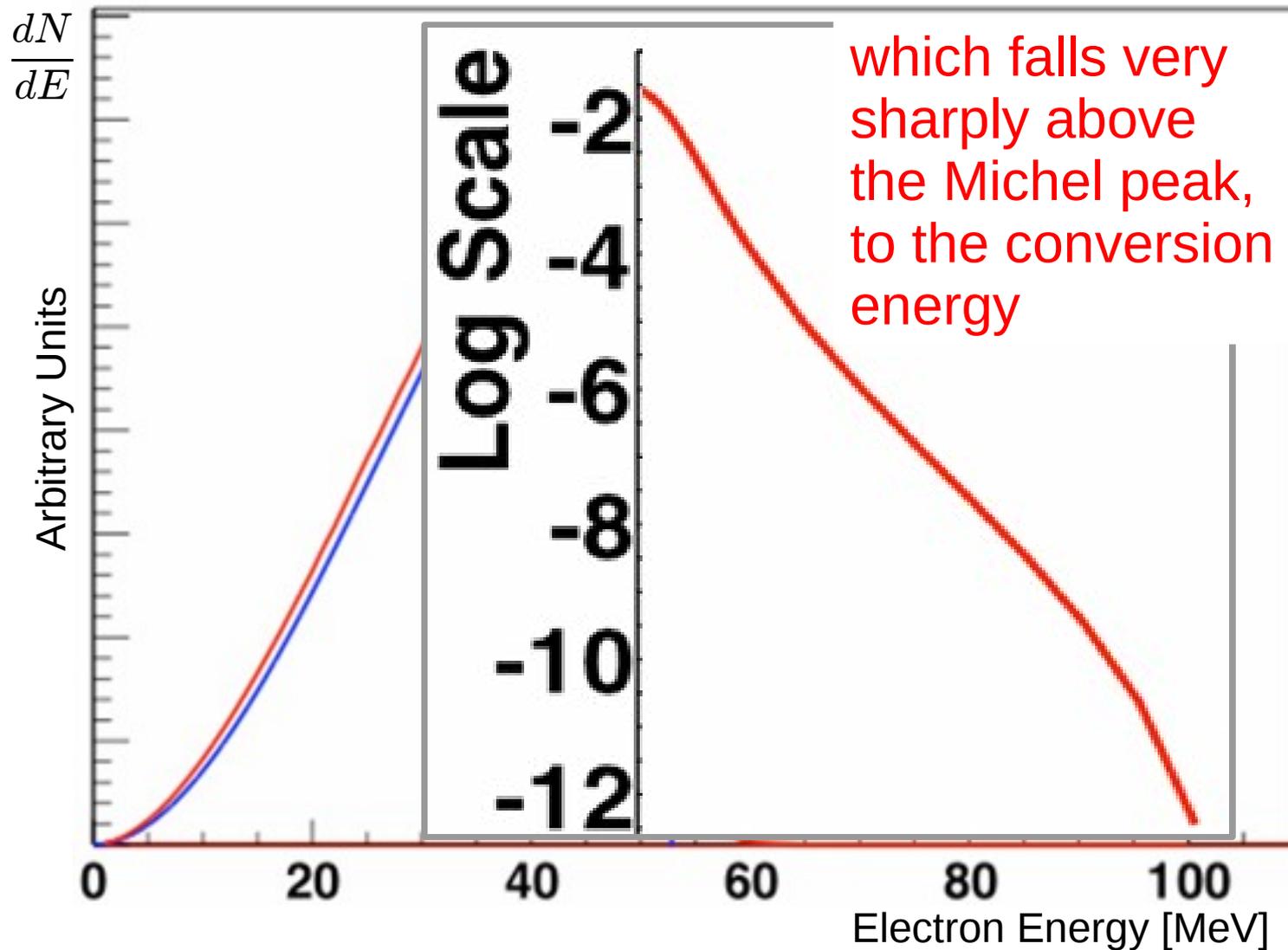
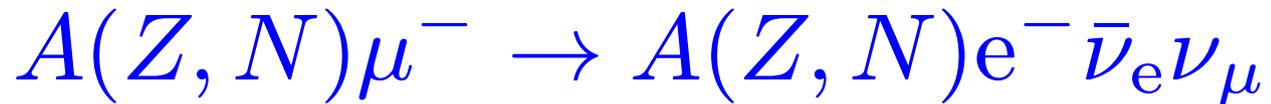
Decay-in-Orbit is the major source of delayed background in the live window



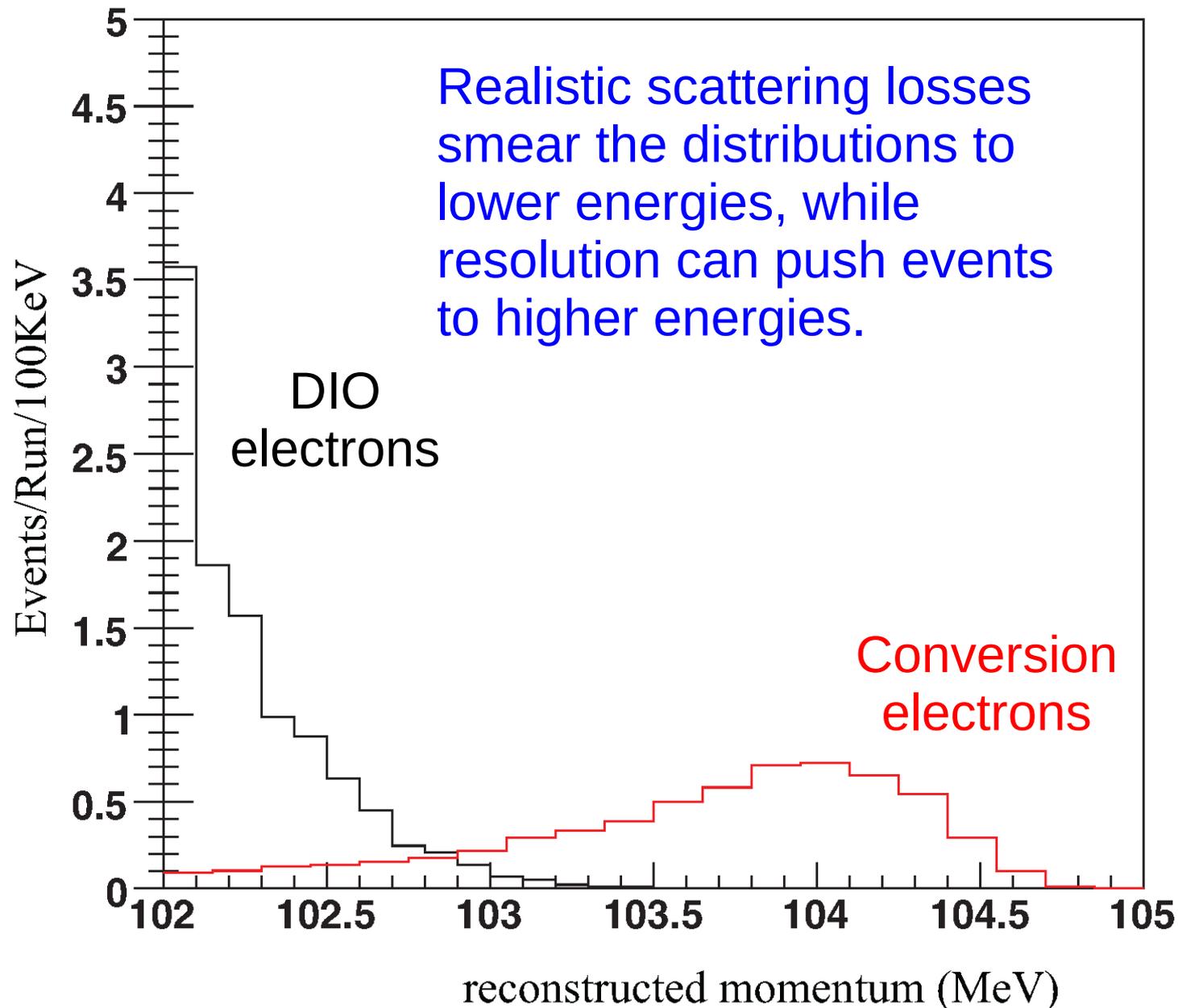
Decay-in-Orbit is the major source of delayed background in the live window



Decay-in-Orbit is the major source of delayed background in the live window



An experimental reality check



The backgrounds and resolutions constrain the design and reach of the experiment

The backgrounds and resolutions constrain the design and reach of the experiment

- Spectrometer design:
 - Track momentum resolution: $O(200\text{keV})$
- Beam transport
 - Out-of-time particle suppression: 10^9

The backgrounds and resolutions constrain the design and reach of the experiment

- Spectrometer design:
 - Track momentum resolution: $O(200\text{keV})$
- Beam transport
 - Out-of-time particle suppression: 10^9

In two years of running

The backgrounds and resolutions constrain the design and reach of the experiment

- Spectrometer design:
 - Track momentum resolution: $O(200\text{keV})$
- Beam transport
 - Out-of-time particle suppression: 10^9

In two years of running

4×10^{20} POT

The backgrounds and resolutions constrain the design and reach of the experiment

- Spectrometer design:
 - Track momentum resolution: $O(200\text{keV})$
- Beam transport
 - Out-of-time particle suppression: 10^9

In two years of running

$$4 \times 10^{20} \text{ POT} \begin{array}{l} \downarrow \\ \rightarrow \end{array} 1 \times 10^{18} \text{ muons}$$

The backgrounds and resolutions constrain the design and reach of the experiment

- Spectrometer design:
 - Track momentum resolution: $O(200\text{keV})$
- Beam transport
 - Out-of-time particle suppression: 10^9

In two years of running

4×10^{20} POT

└─→ 1×10^{18} muons

└─→ $R_{\mu e} \leq 6 \times 10^{-17}$ (90%CL)

The backgrounds and resolutions constrain the design and reach of the experiment

- Spectrometer design:
 - Track momentum resolution: $O(200\text{keV})$
- Beam transport
 - Out-of-time particle suppression: 10^9

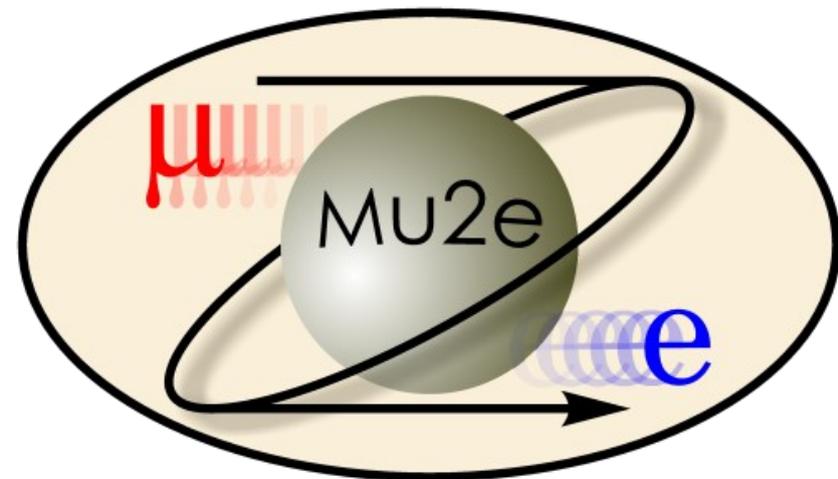
In two years of running

$$\begin{aligned} 4 \times 10^{20} \text{ POT} & \longrightarrow 1 \times 10^{18} \text{ muons} \\ & \longrightarrow R_{\mu e} \leq 6 \times 10^{-17} \text{ (90\%CL)} \end{aligned}$$

This is a factor 10^4 improvement
on the current limit!

A brief roadmap to this talk...

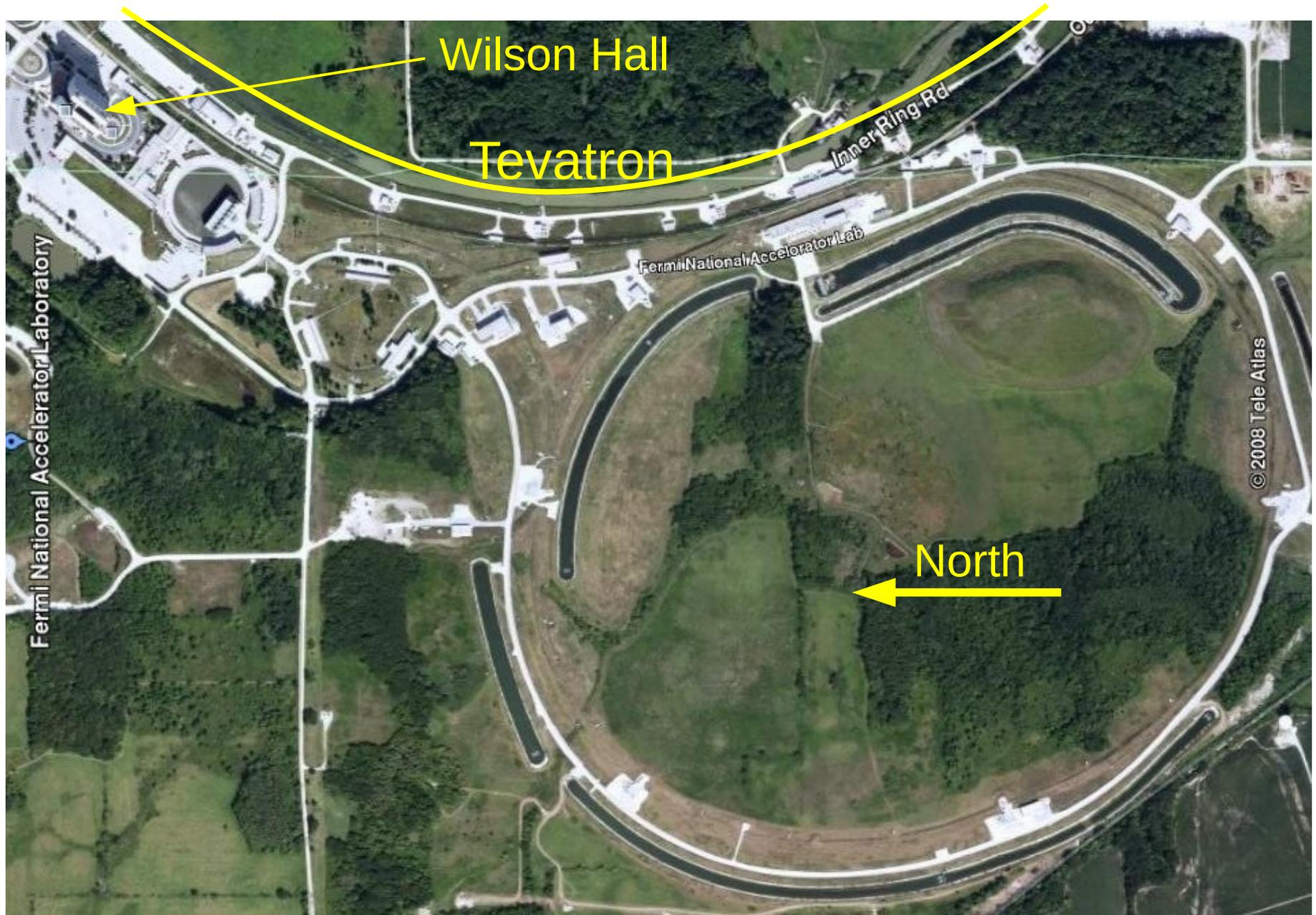
- Lepton Flavor Violation
- Primary backgrounds to muon conversion
- The design of Mu2e



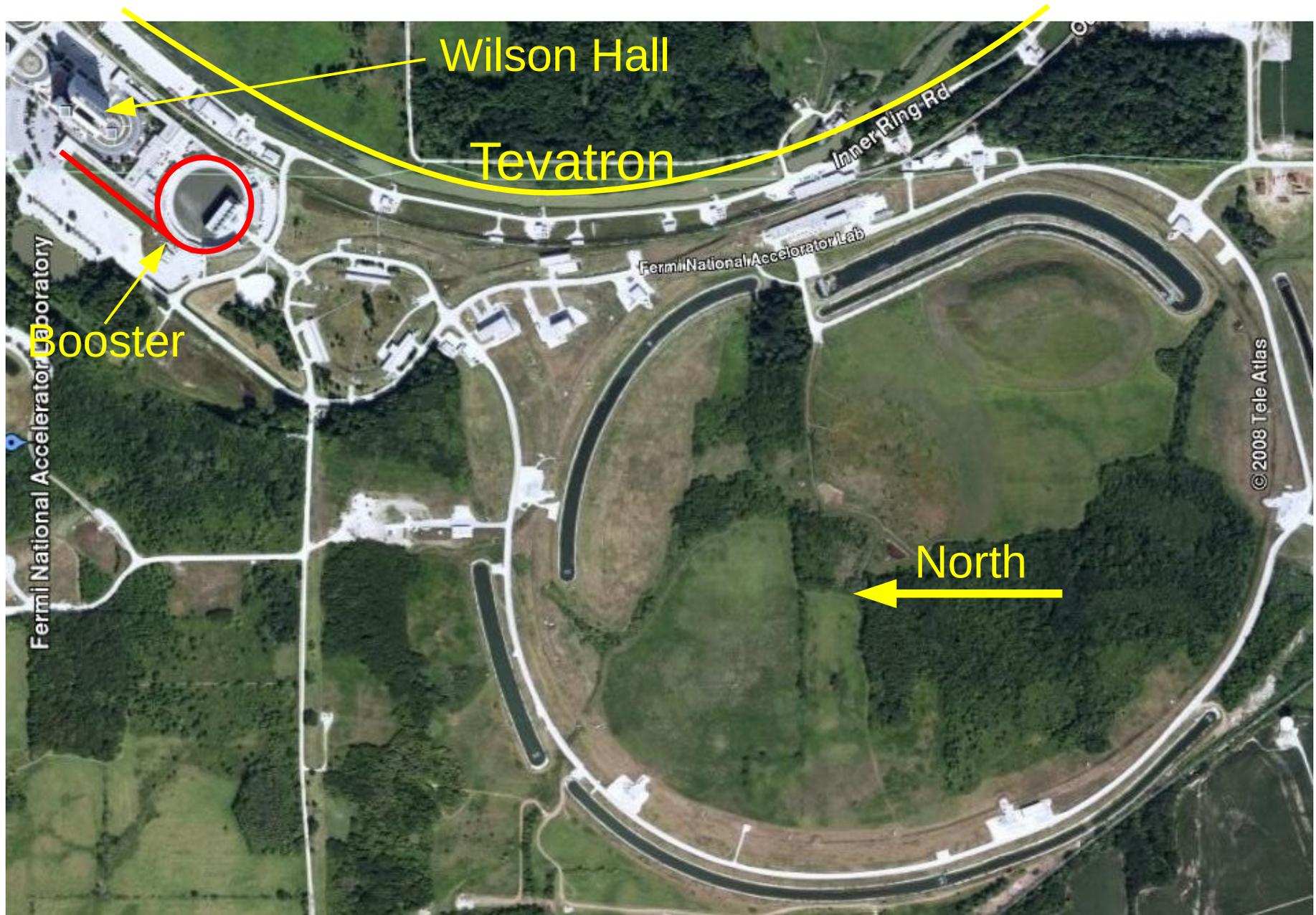
Where are we going to build this thing?



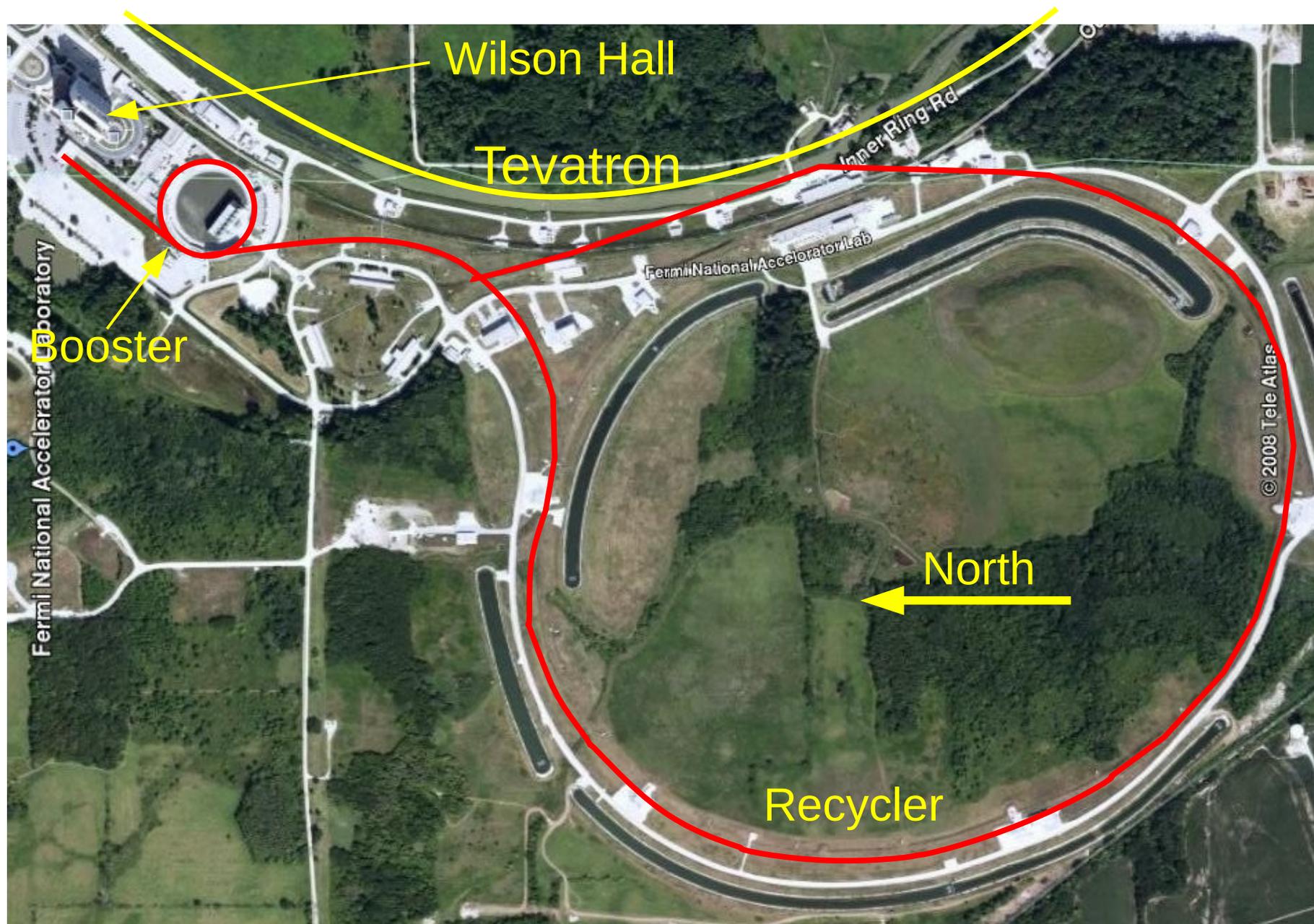
Where are we going to build this thing?



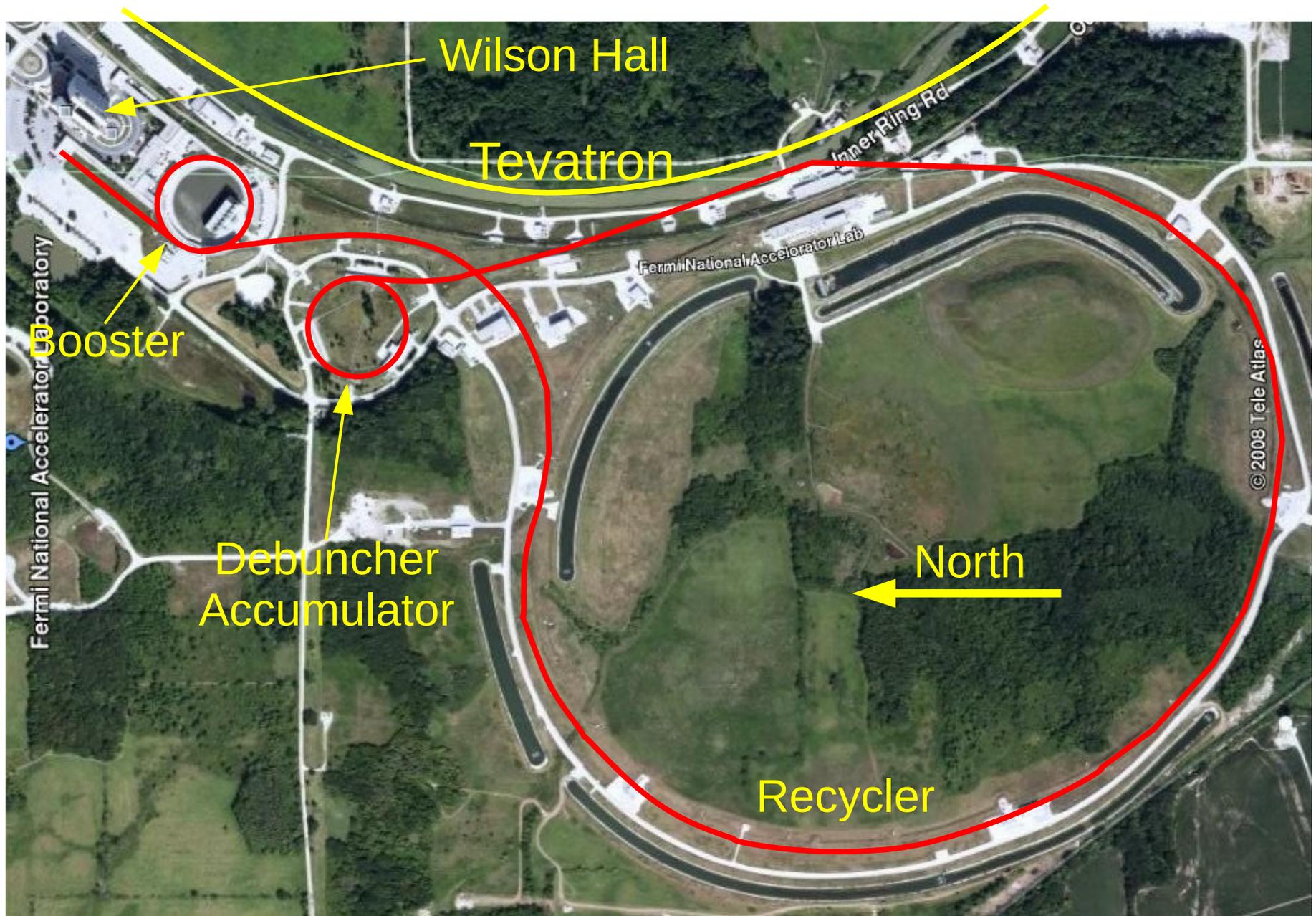
Where are we going to build this thing?



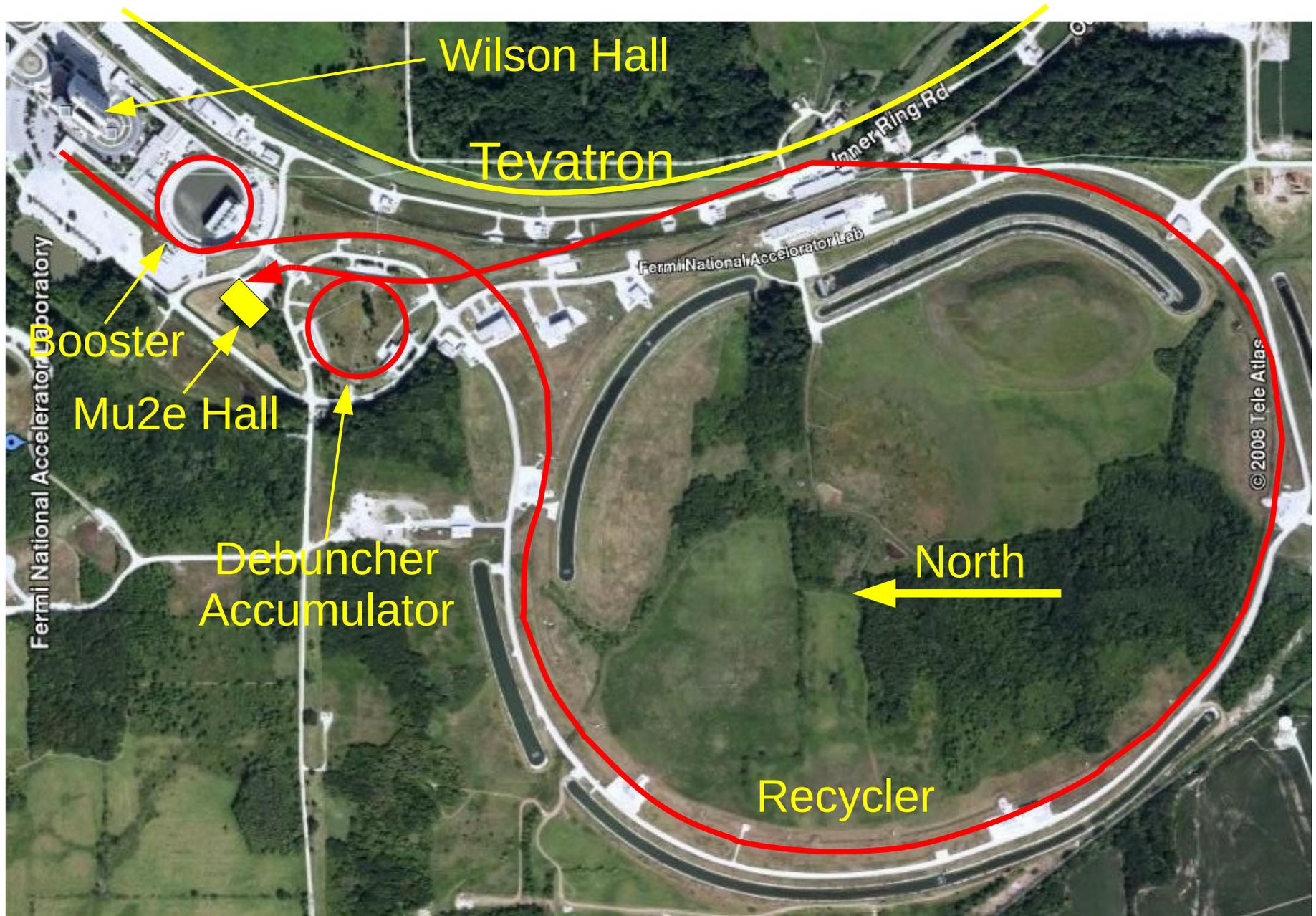
Where are we going to build this thing?



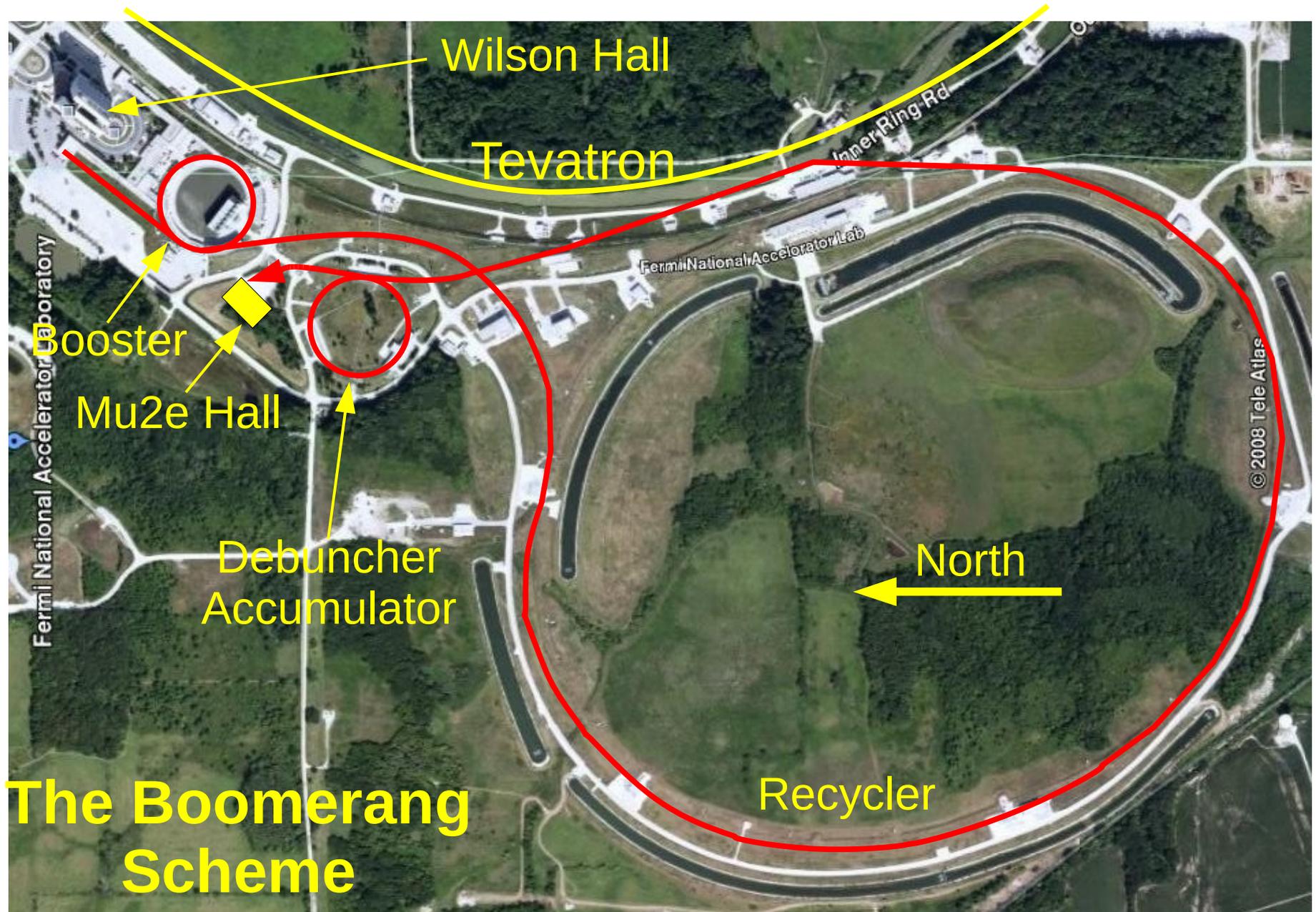
Where are we going to build this thing?



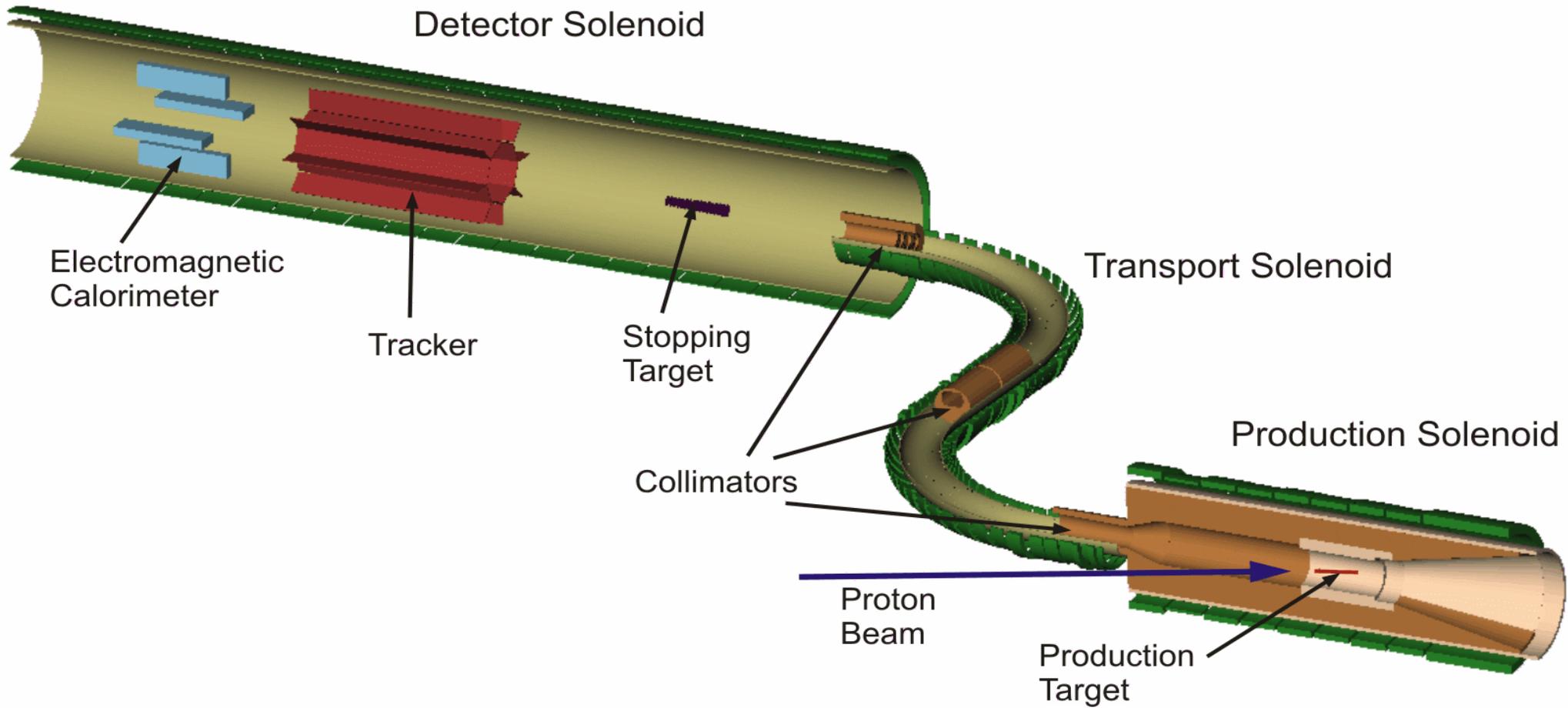
Where are we going to build this thing?



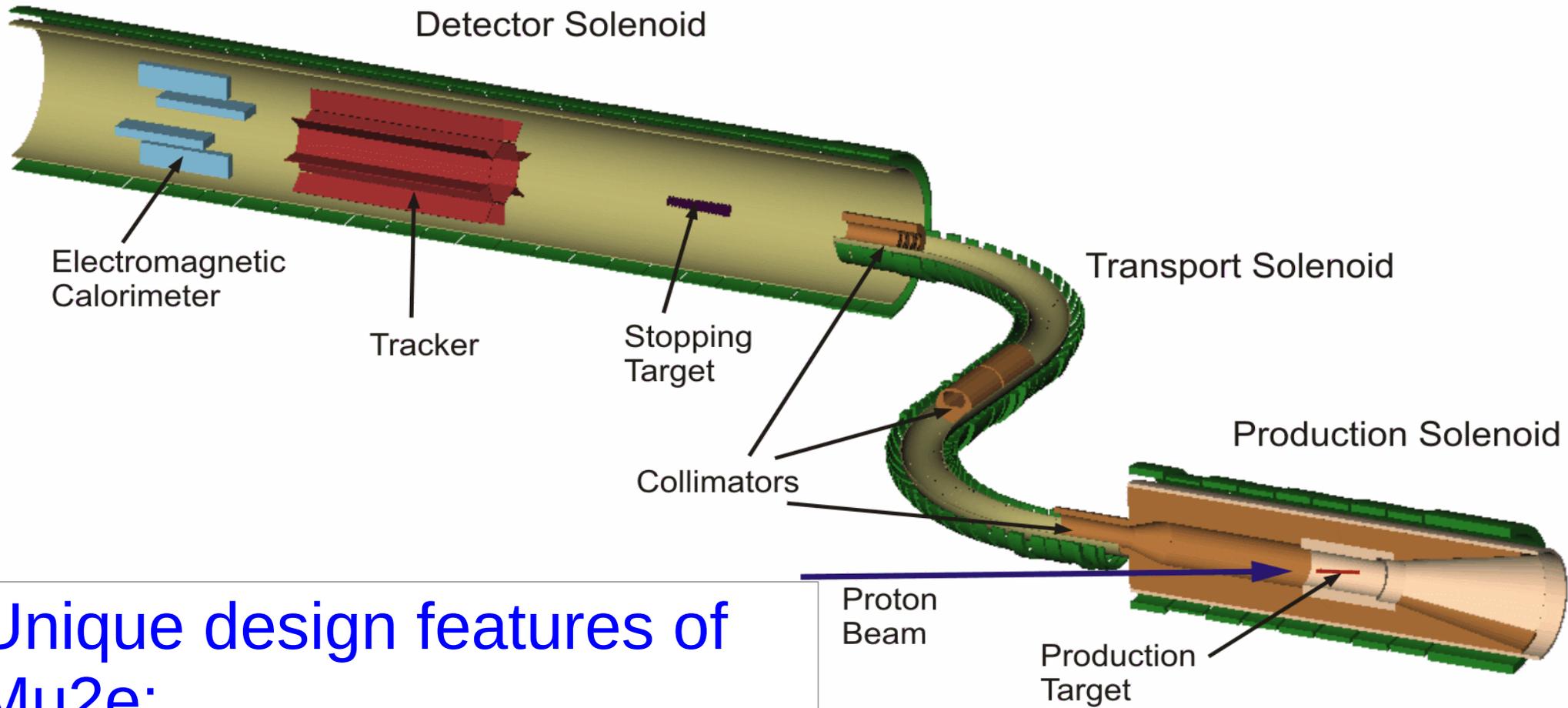
Where are we going to build this thing?



What goes in the hall?



What goes in the hall?

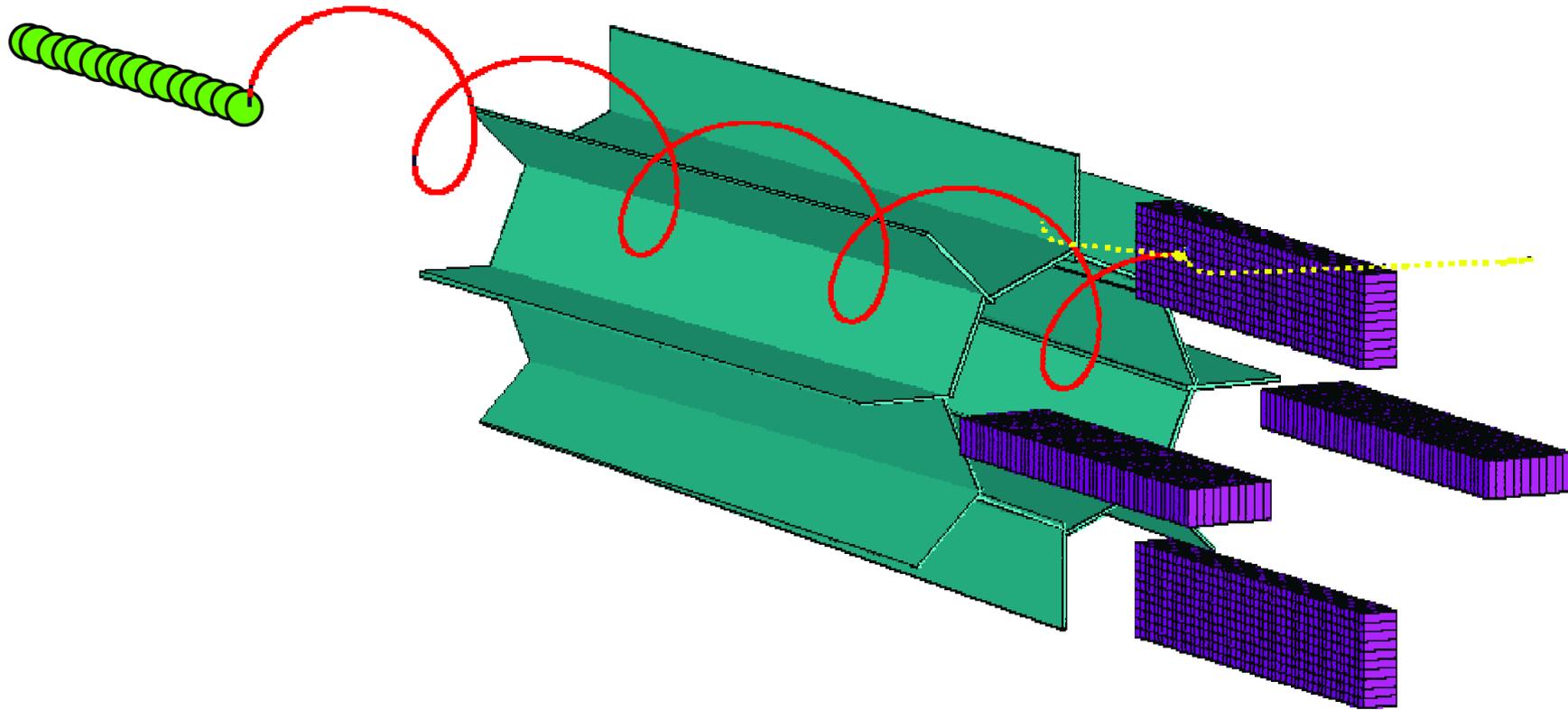


Unique design features of Mu2e:

- Pulsed proton source
- Curved transport solenoid
- Graded detector solenoid fields

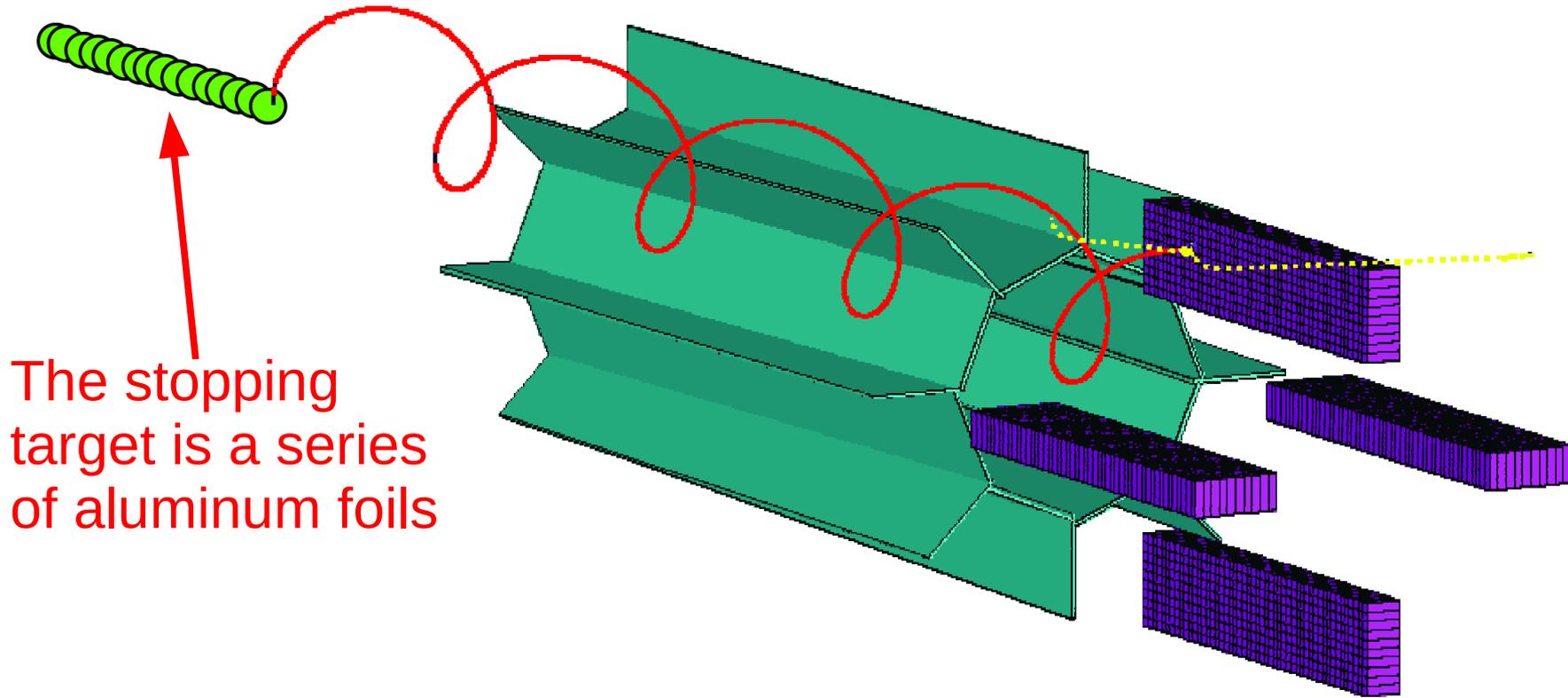
The target and detectors

A conversion electron track



The target and detectors

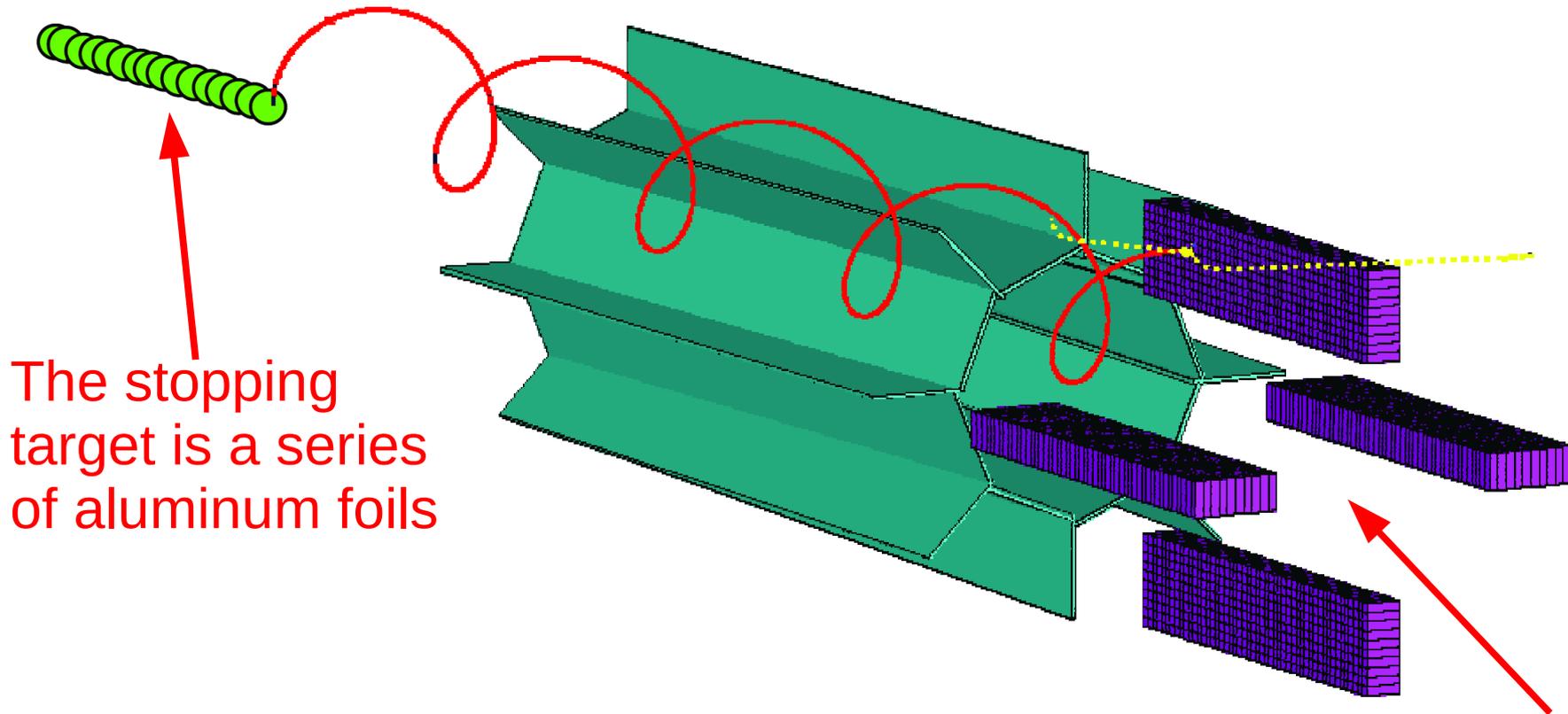
A conversion electron track



The stopping target is a series of aluminum foils

The target and detectors

A conversion electron track

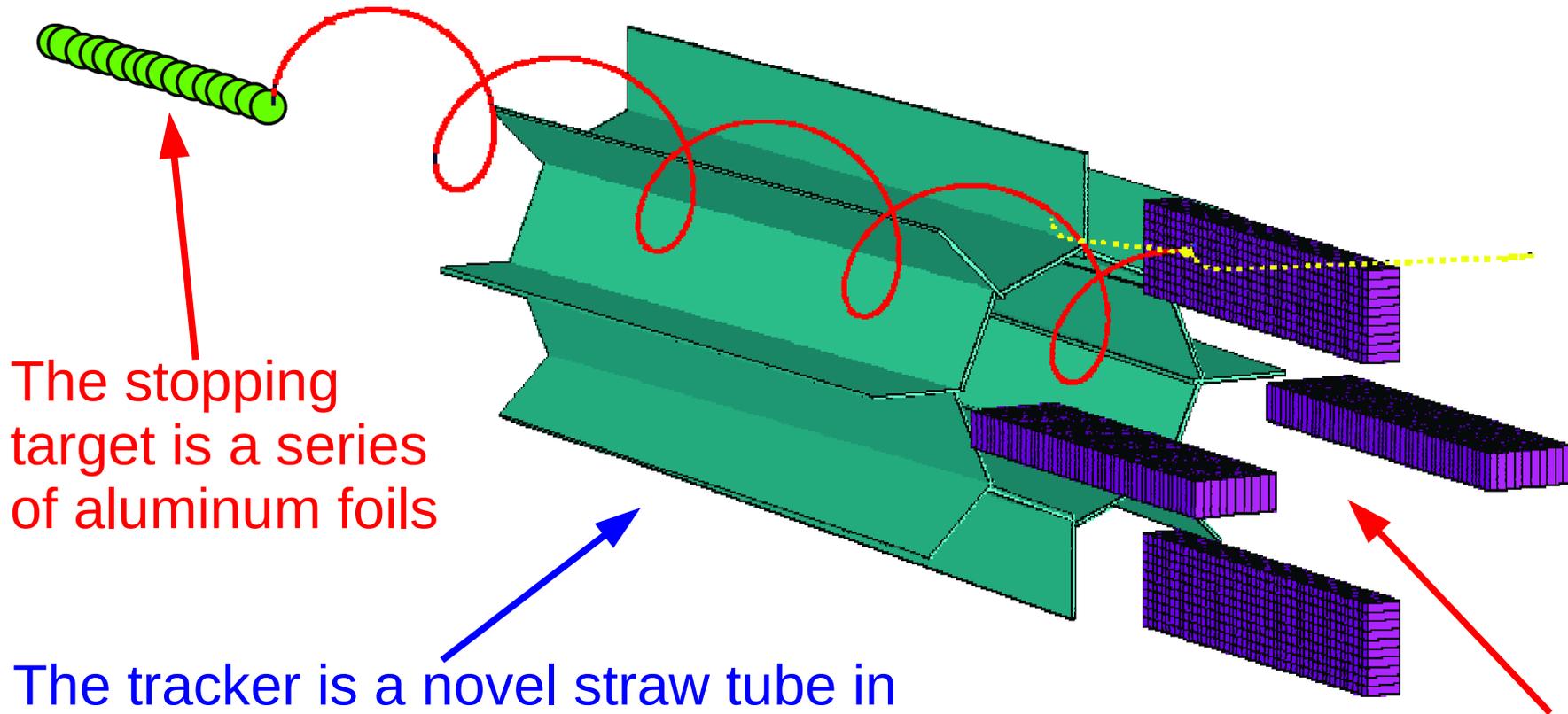


The stopping target is a series of aluminum foils

The remnant muon beam and most DIO electrons pass through the central openings, and are caught by a beam stop.

The target and detectors

A conversion electron track



The stopping target is a series of aluminum foils

The tracker is a novel straw tube in vacuum design. A number of baseline designs are under consideration; all would have or order 20000 readout channels.

The remnant muon beam and most DIO electrons pass through the central openings, and are caught by a beam stop.

The target and detectors

The crystal calorimeter provides direct measures of track position, time, and energy to help protect against catastrophic reconstruction errors.

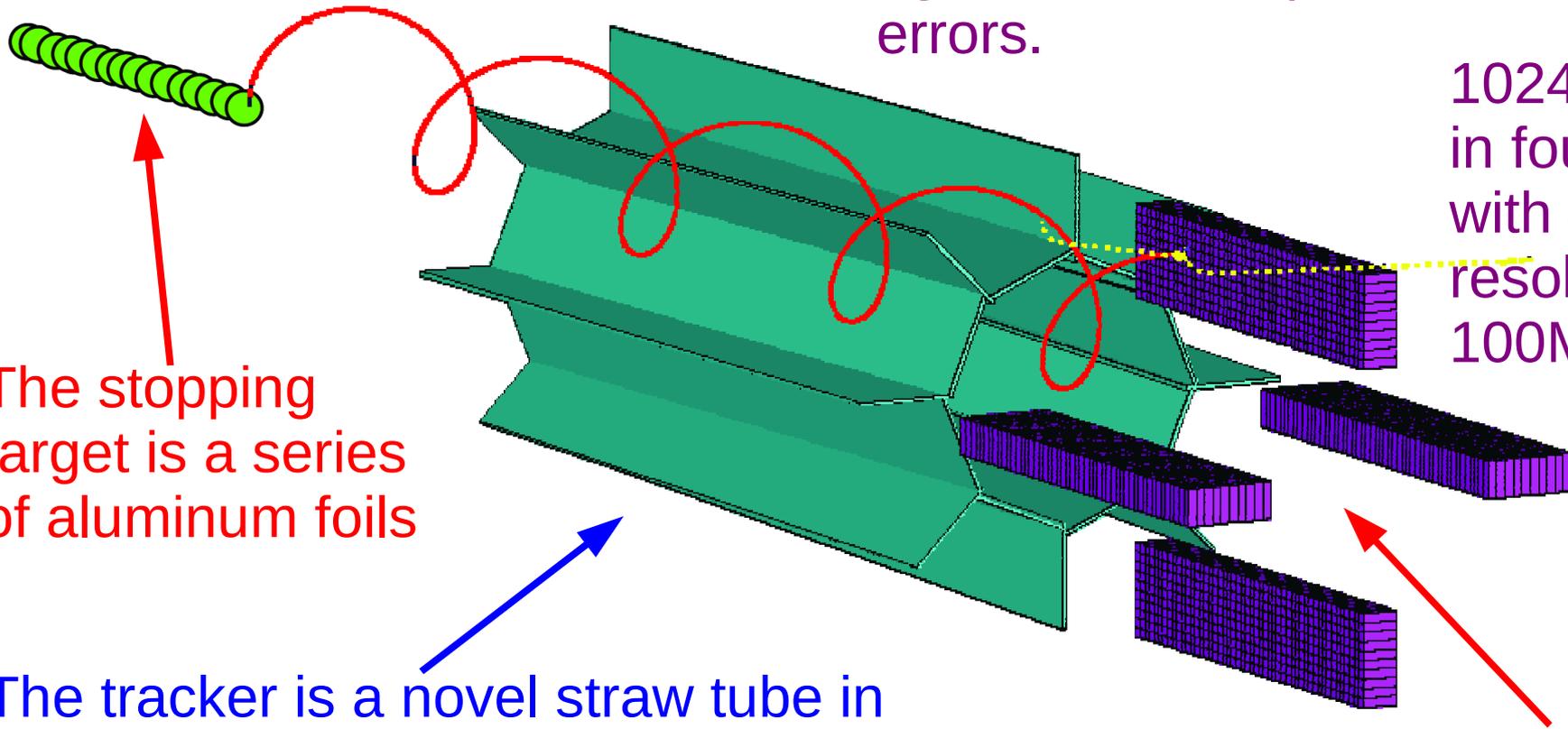
A conversion electron track

1024 crystals in four vanes with 4MeV resolution at 100MeV

The stopping target is a series of aluminum foils

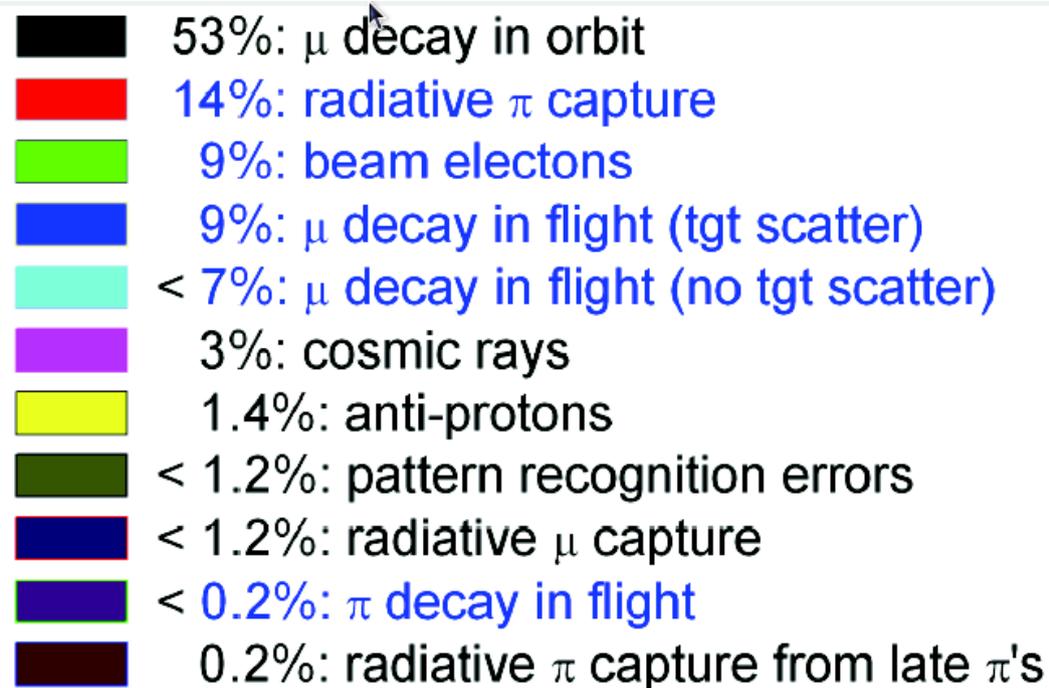
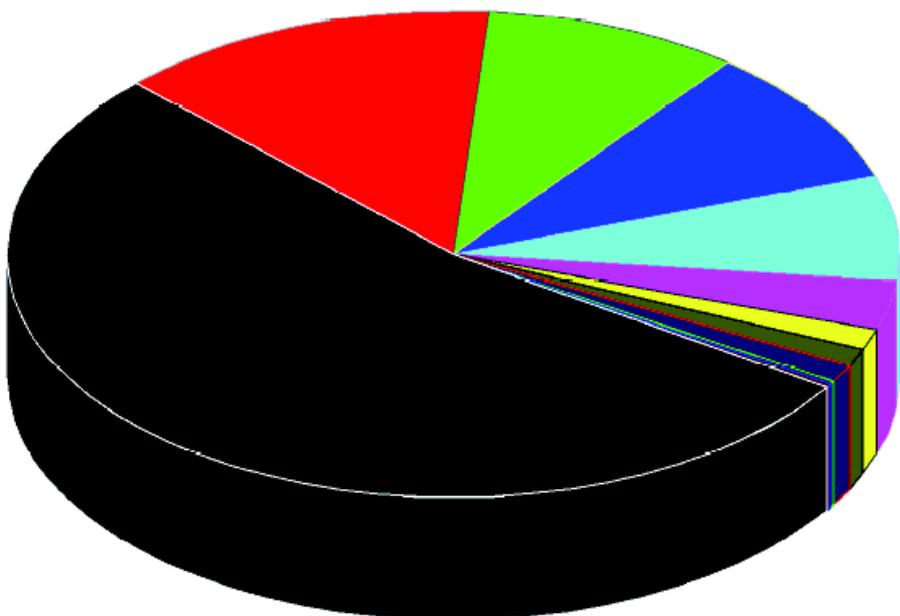
The tracker is a novel straw tube in vacuum design. A number of baseline designs are under consideration; all would have or order 20000 readout channels.

The remnant muon beam and most DIO electrons pass through the central openings, and are caught by a beam stop.



With this design, the background rates are indeed very small!

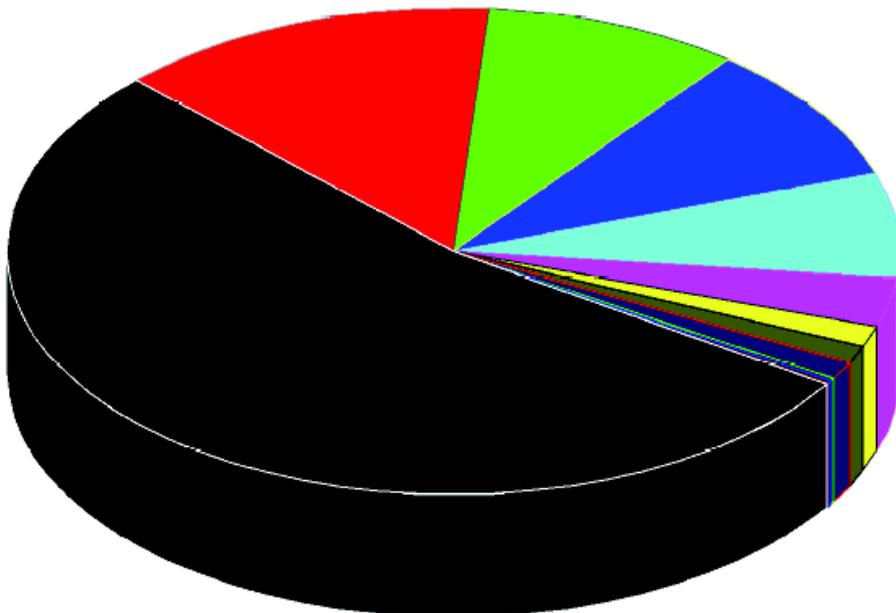
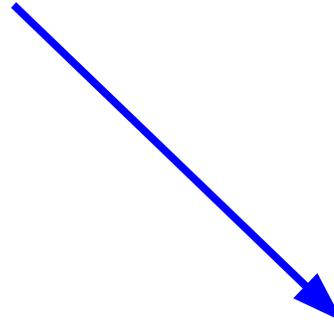
Background	Events per 2×10^7 s
μ DIO	0.225
Radiative π capture	0.072
μ DIF	0.072
Scattered beam e	0.035
Total	0.41



With this design, the background rates are indeed very small!

In two years of running, we expect fewer than one background event in the signal region.

Background	Events per 2×10^7 s
μ DIO	0.225
Radiative π capture	0.072
μ DIF	0.072
Scattered beam e	0.035
Total	0.41



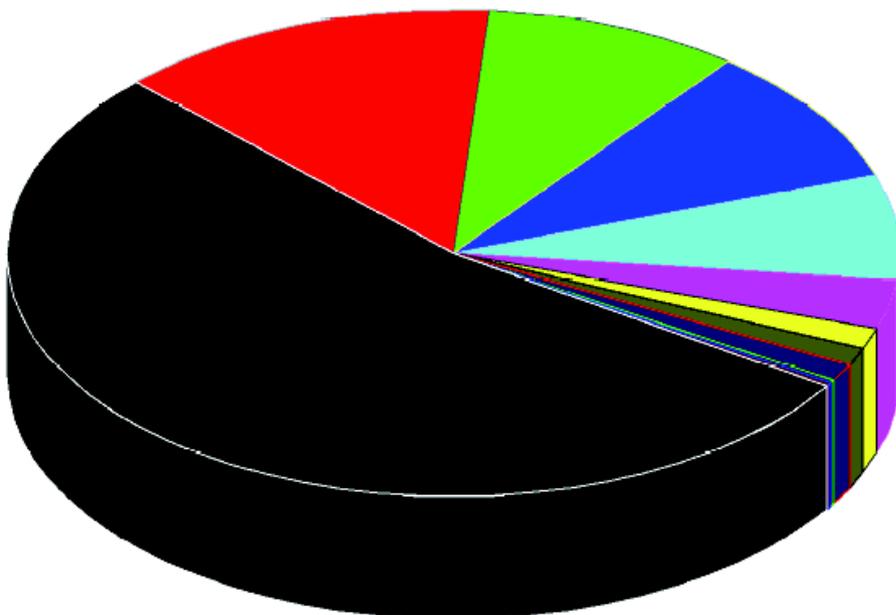
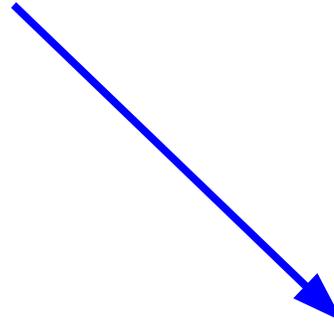
- 53%: μ decay in orbit
- 14%: radiative π capture
- 9%: beam electrons
- 9%: μ decay in flight (tgt scatter)
- < 7%: μ decay in flight (no tgt scatter)
- 3%: cosmic rays
- 1.4%: anti-protons
- < 1.2%: pattern recognition errors
- < 1.2%: radiative μ capture
- < 0.2%: π decay in flight
- 0.2%: radiative π capture from late π 's

With this design, the background rates are indeed very small!

In two years of running, we expect fewer than one background event in the signal region.

If $R_{\mu e} = 10^{-15}$, we will see 40 signal events

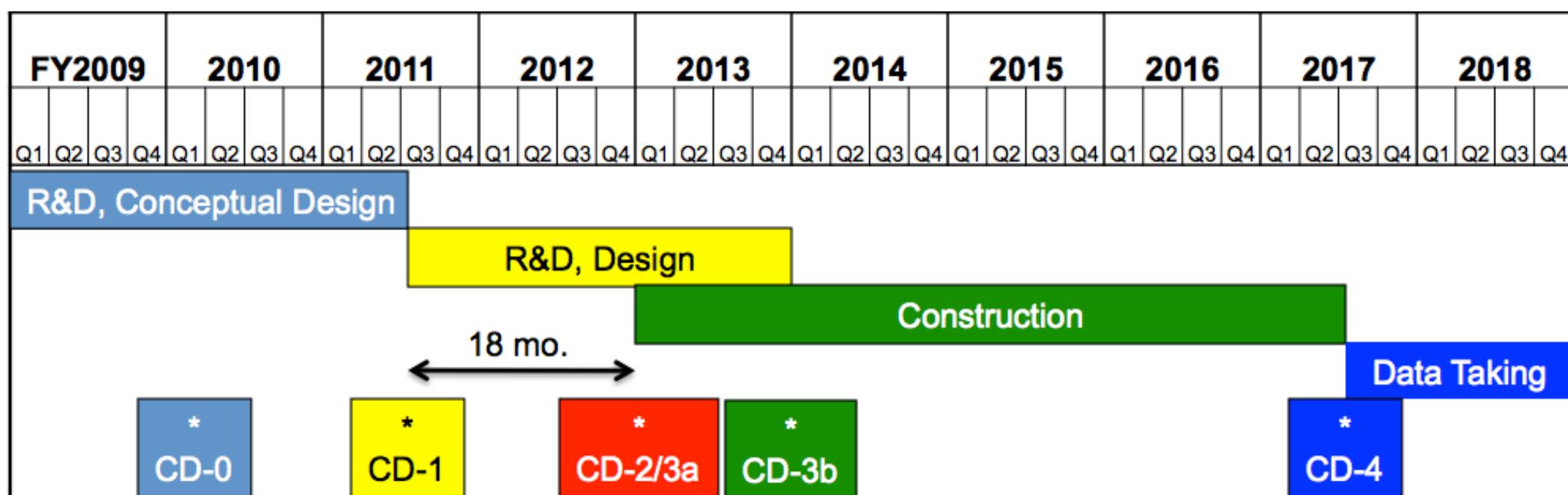
Background	Events per 2×10^7 s
μ DIO	0.225
Radiative π capture	0.072
μ DIF	0.072
Scattered beam e	0.035
Total	0.41



- 53%: μ decay in orbit
- 14%: radiative π capture
- 9%: beam electrons
- 9%: μ decay in flight (tgt scatter)
- < 7%: μ decay in flight (no tgt scatter)
- 3%: cosmic rays
- 1.4%: anti-protons
- < 1.2%: pattern recognition errors
- < 1.2%: radiative μ capture
- < 0.2%: π decay in flight
- 0.2%: radiative π capture from late π 's

The schedule is technically driven

- We have received both Stage 1 approval and CD-0.
- Magnet design and construction remain the schedule drivers.



Please join us!

- Boston University
- Brookhaven National Laboratory
- City University of New York
- College of William & Mary
- Fermi National Accelerator Laboratory
- Idaho State University
- Istituto Nazionale di Fisica Nucleare, Lecce
- Istituto Nazionale di Fisica Nucleare, Pisa
- Institute for Nuclear Research, Moscow
- Joint Institute for Nuclear Research, Dubna
- Laboratori Nazionali di Frascati

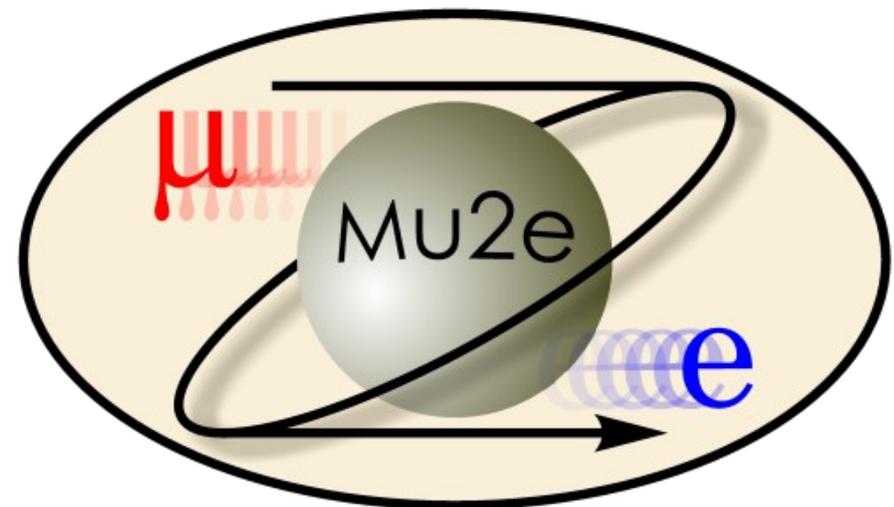
Contact:

Bob Bernstein, rhbob@fnal.gov,

or

Jim Miller, miller@bu.edu.

- Lawrence Berkeley National Laboratory
- Los Alamos National Laboratory
- Muons Inc.
- Northwestern University
- Rice University
- Syracuse University
- University of California, Berkeley
- University of California, Irvine
- University of Houston
- University of Illinois
- University of Massachusetts, Amherst
- University of Virginia
- University of Washington



We normalize the conversions with the captures

$$R_{\mu e} = \frac{\Gamma(\mu^- A \rightarrow e^- A)}{\Gamma(\mu^- A \rightarrow \nu_\mu A')}$$

Events in signal window

Acceptance for signal events

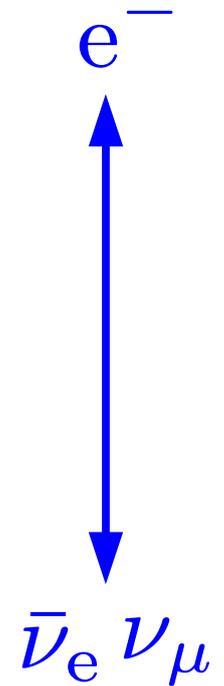
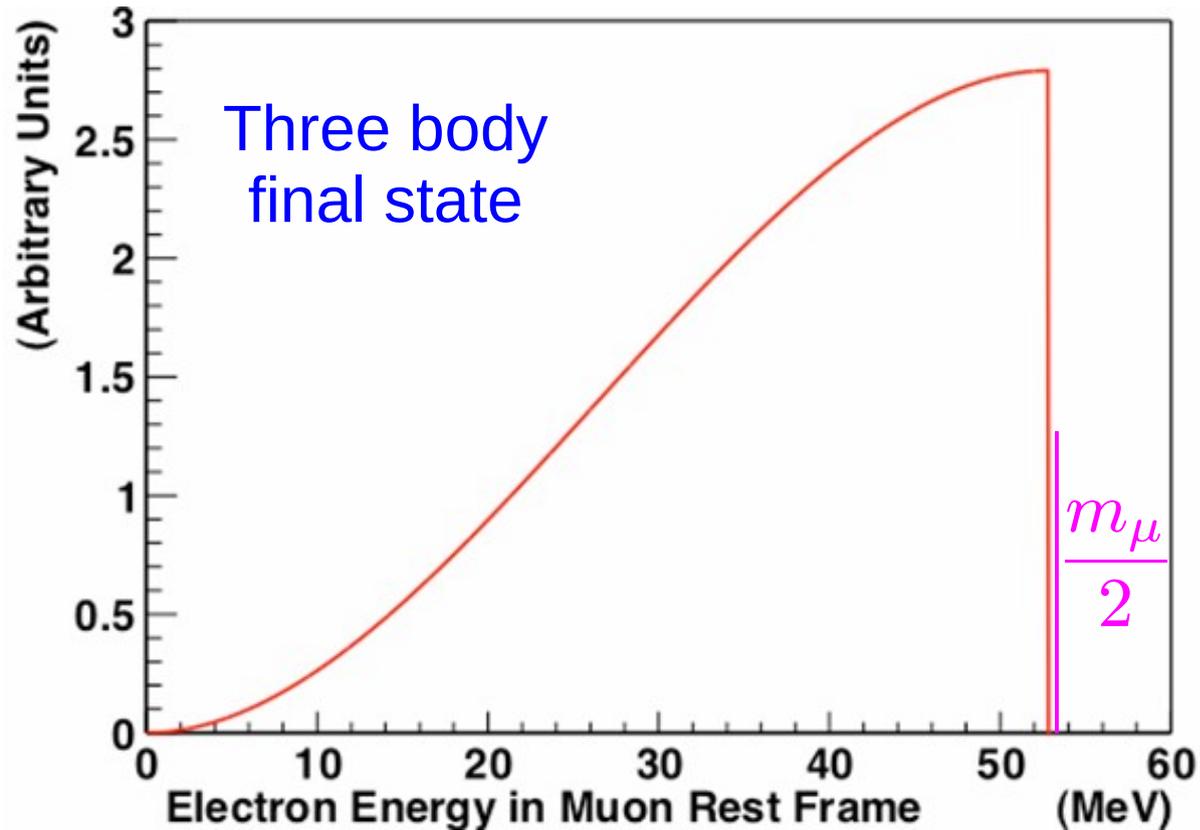
$$= \frac{N_{\mu e} / \epsilon_{\mu e}}{N_{\text{stops}} / \epsilon_{\text{stops}} (\Gamma_{\mu\nu} / \Gamma_{\text{total}})}$$

Directly measured via cascade X-Rays

Well known nuclear capture ratio

Muon Decay-in-Orbit

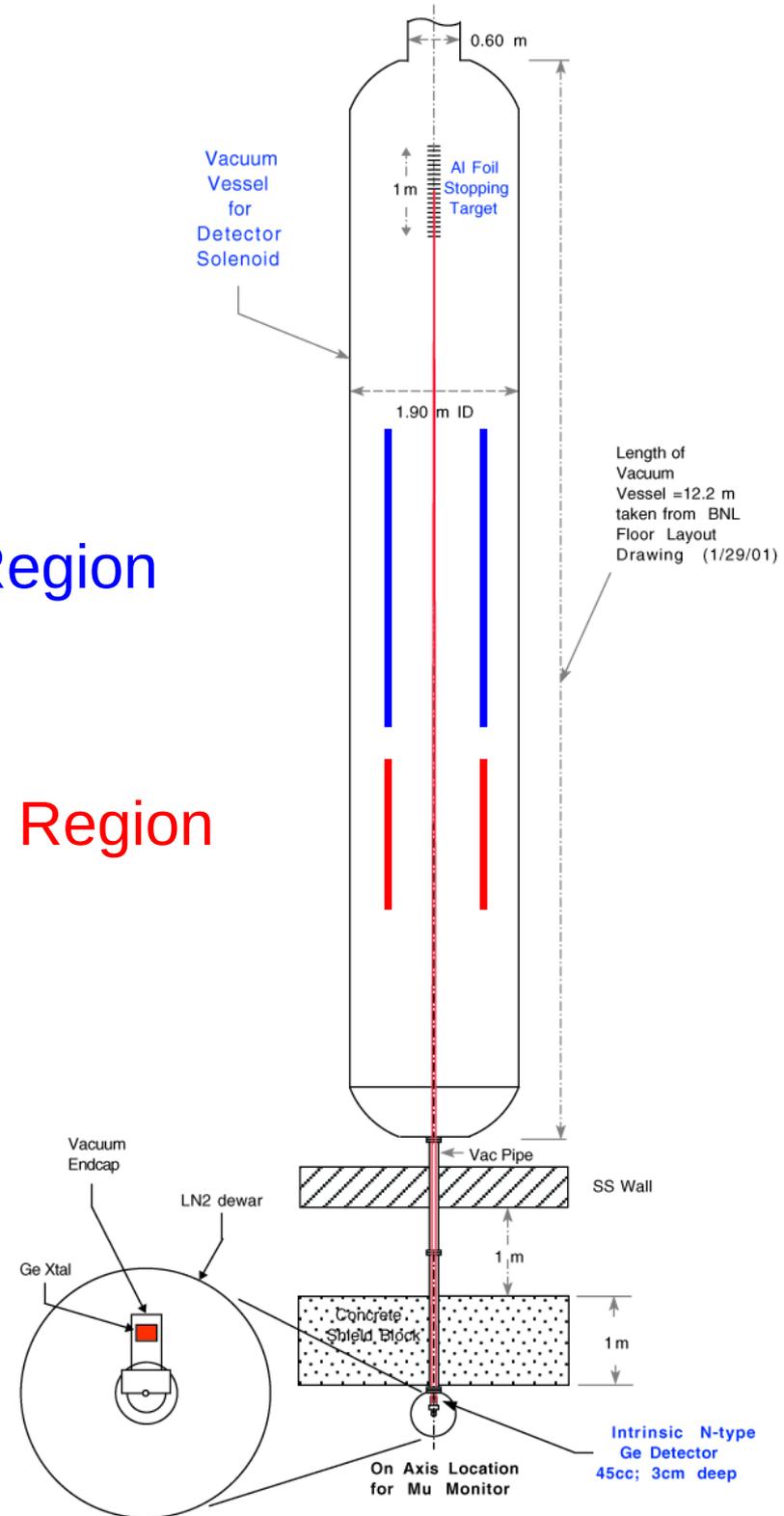
First, consider the free decay



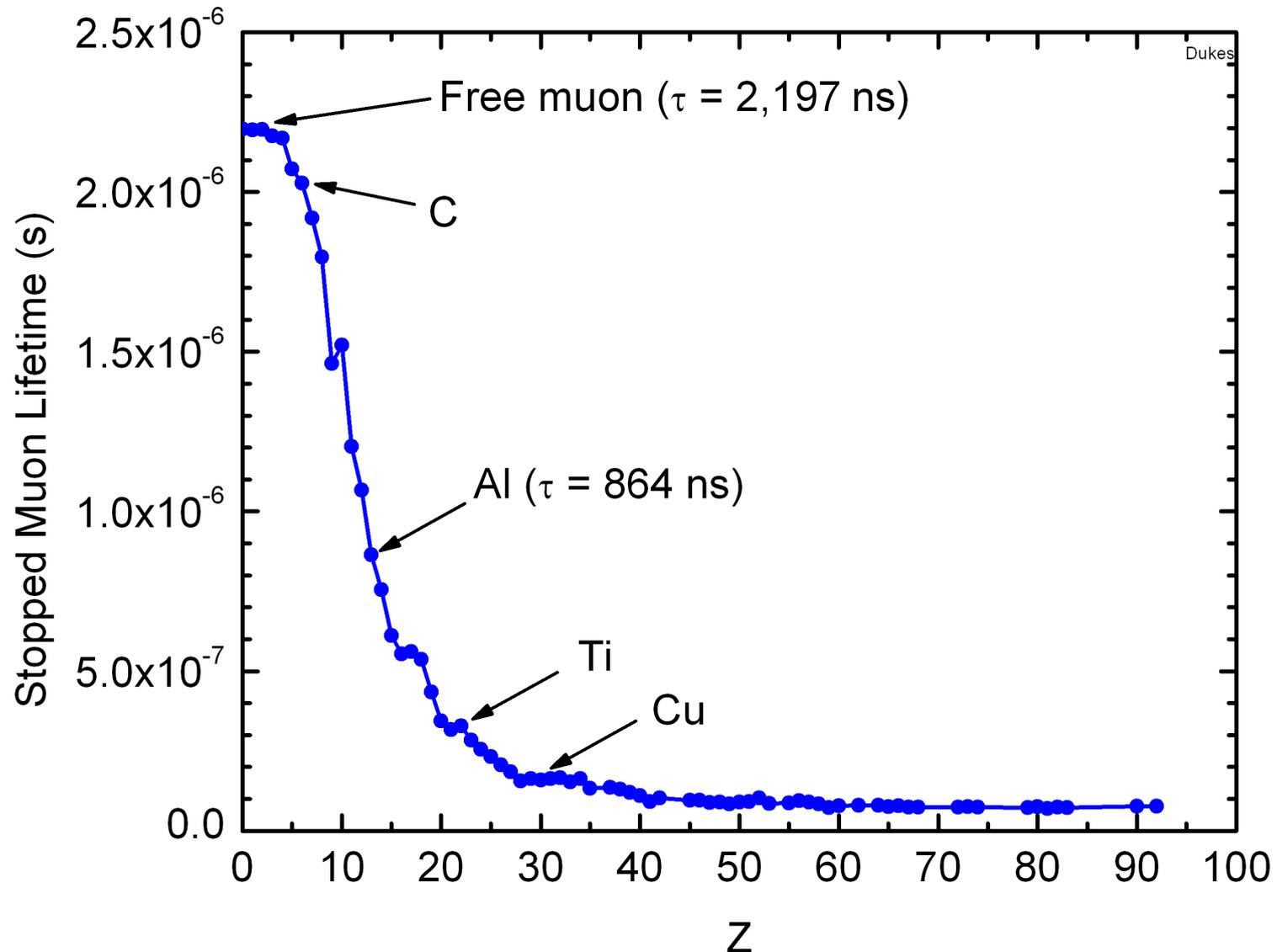
Monitoring the muon stopping rate

Tracker Region

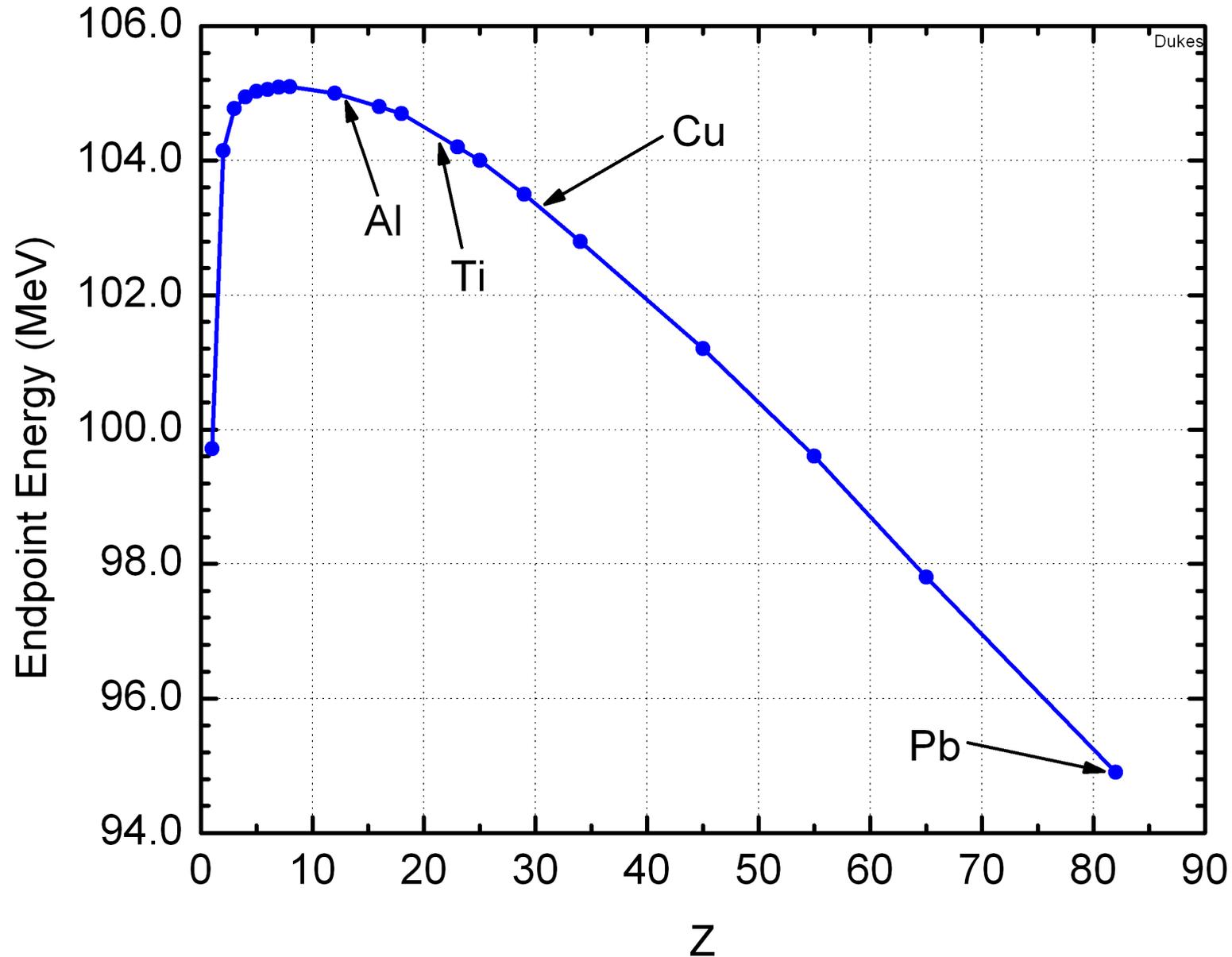
Calorimeter Region



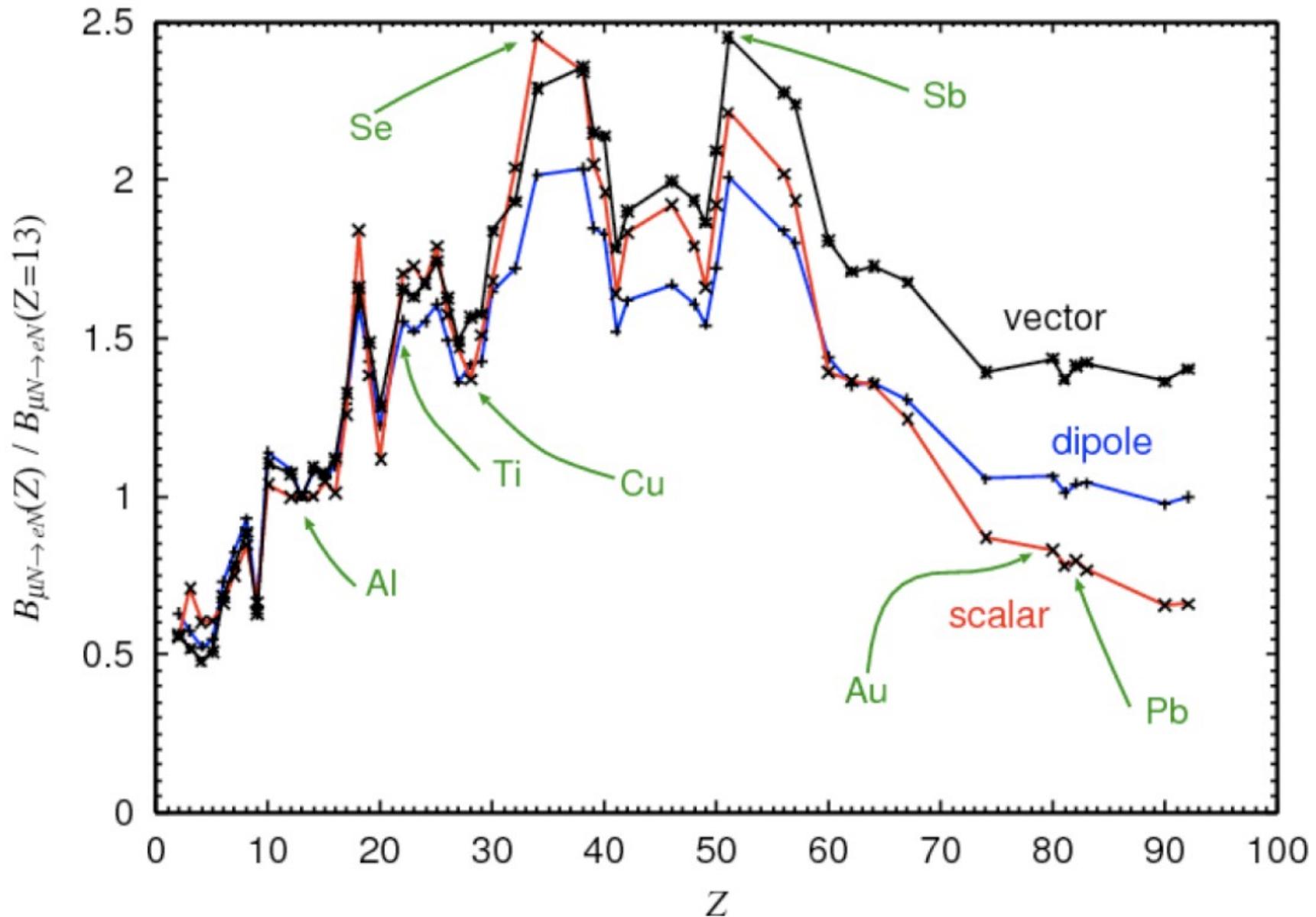
The choice of Al is well matched to the FNAL beam time structure



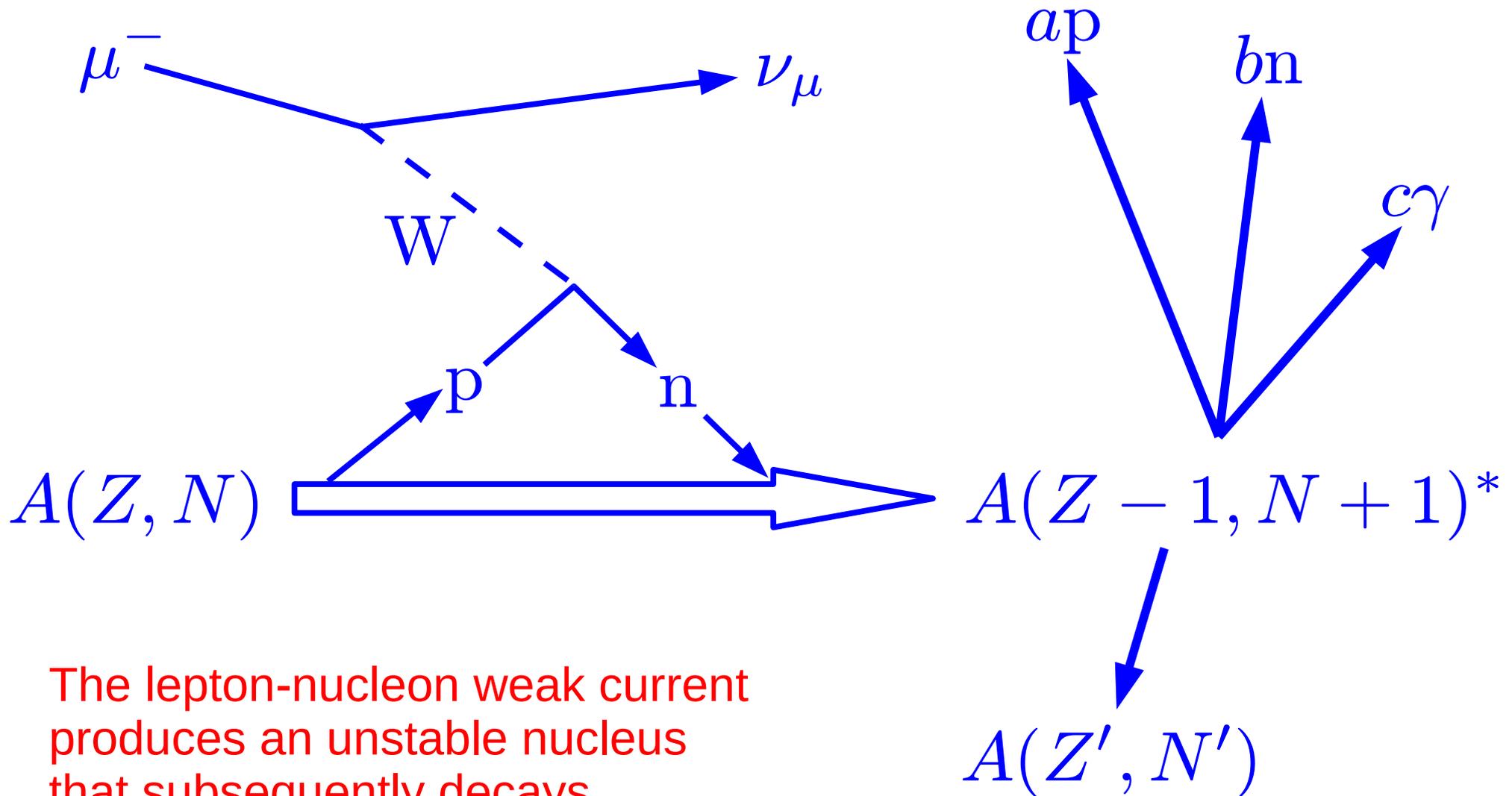
The endpoint energy is material dependent



Different materials are sensitive to different operators

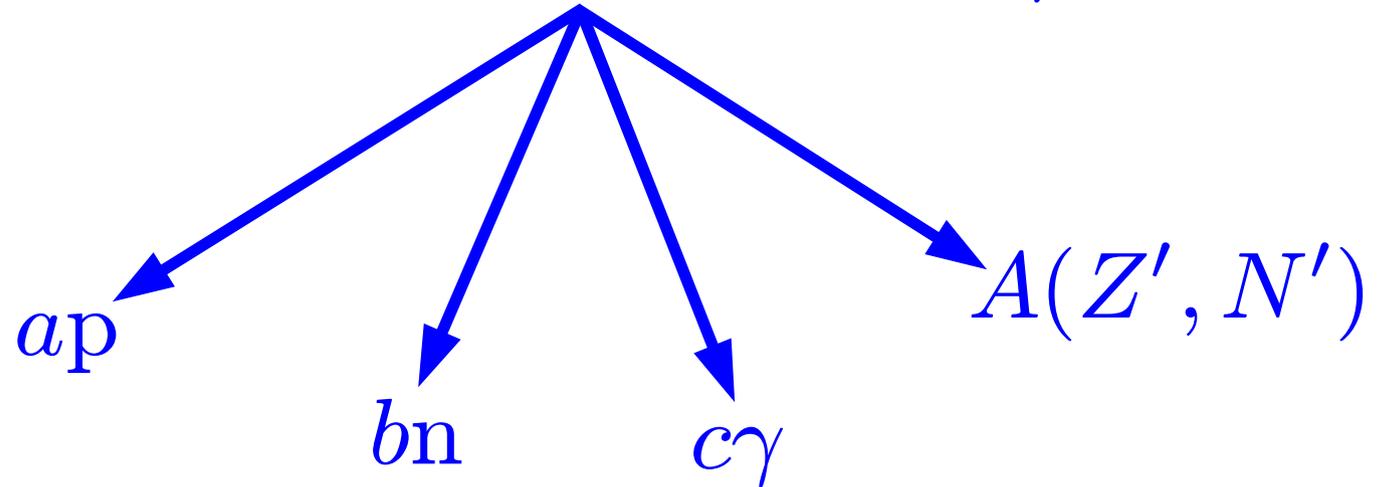


Nuclear capture is both normalization and a source of inefficiency due to extra hits



The lepton-nucleon weak current produces an unstable nucleus that subsequently decays

Nuclear capture is both normalization and a source of inefficiency due to extra hits

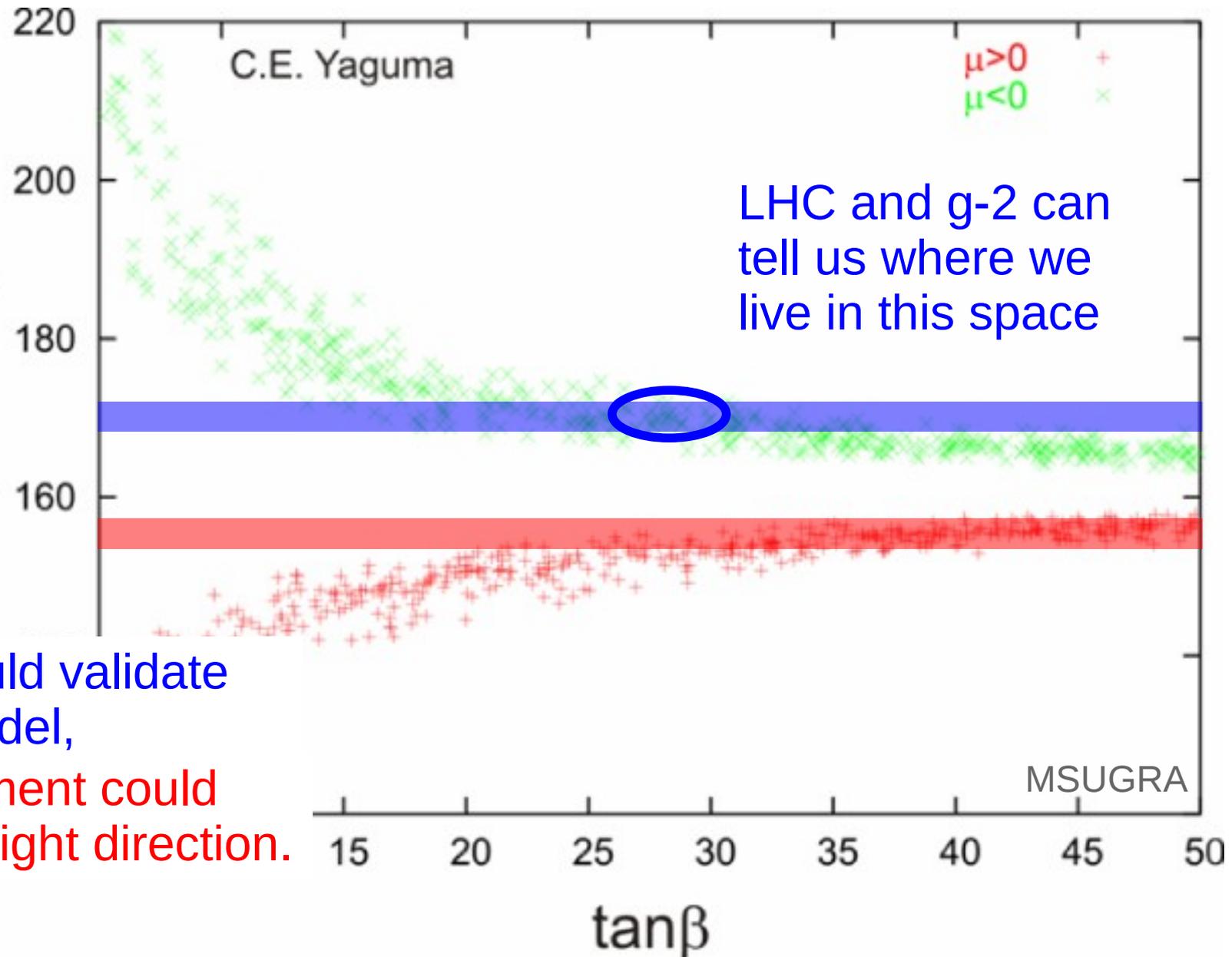


For an ^{27}Al target, the average capture results in $0.1p$, $2n$, and 2γ . These result in additional hits in the detectors, potentially leading to inefficiency and pattern recognition errors.

Conversion is just one complementary piece of the flavor puzzle

The ratio of MEG to Mu2e observations

$$\frac{BR(\mu \rightarrow e\gamma)}{R(\mu\text{Ti} \rightarrow e\text{Ti})}$$



Agreement would validate any leading model,
while disagreement could point us in the right direction.

Different cLFV processes are driven by different operators

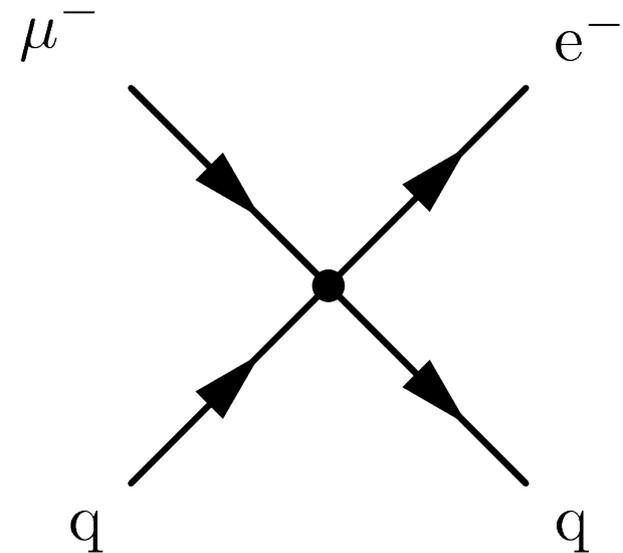
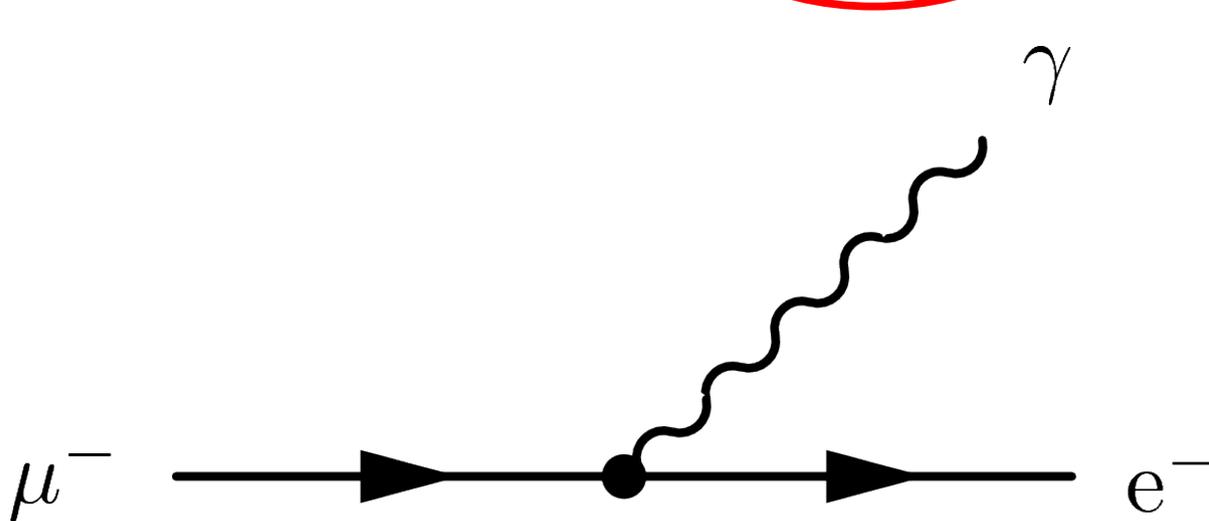
Hence, measurements of the different processes tell us about the underlying physics

$$\mathcal{L}_{\text{cLFV}} = \frac{m_\mu}{(\kappa + 1)\Lambda^2} \bar{\mu}_R \sigma_{\alpha\beta} e_L F^{\alpha\beta} +$$

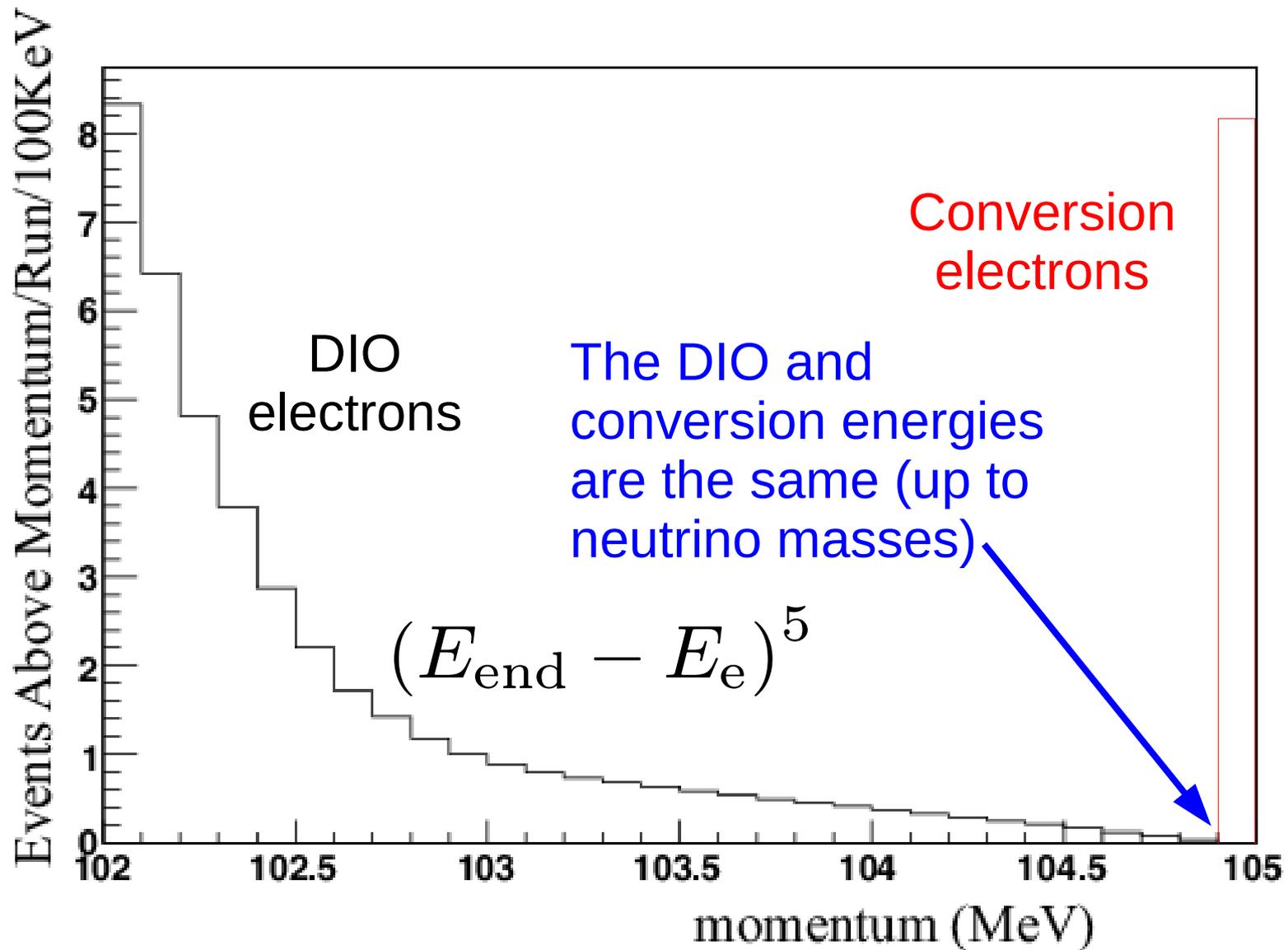
Chirality violating dipole interaction

Chirality preserving contact interaction

$$\frac{\kappa}{(\kappa + 1)\Lambda^2} \bar{\mu}_L \gamma_\alpha e_L (\bar{u}_L \gamma^\alpha u_L + \bar{d}_L \gamma^\alpha d_L)$$



But the DIO endpoint is the conversion energy



In the Project X era ...

... beam power will increase 10-100 fold. Mu2e naturally extends into this era, regardless of the first generation results.

With no signal

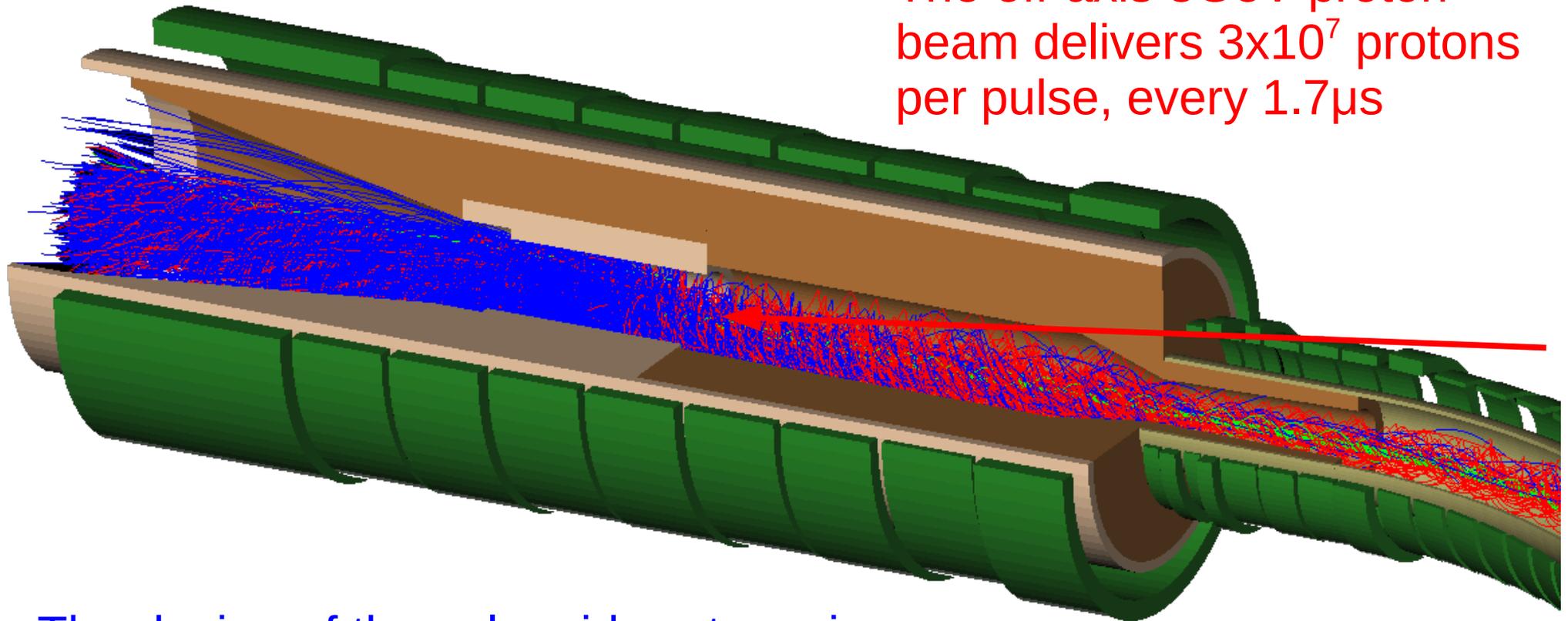
- All backgrounds must be further suppressed to reach $R_{\mu e}$ of 10^{-18}
- All detector and transport systems will need upgrades to handle the higher rates

If we do see a signal

- Change target material to determine structure of new physics amplitudes
- All detector and transport systems will need upgrades to handle the higher rates

Muon production solenoid

The off-axis 8GeV proton beam delivers 3×10^7 protons per pulse, every $1.7 \mu\text{s}$

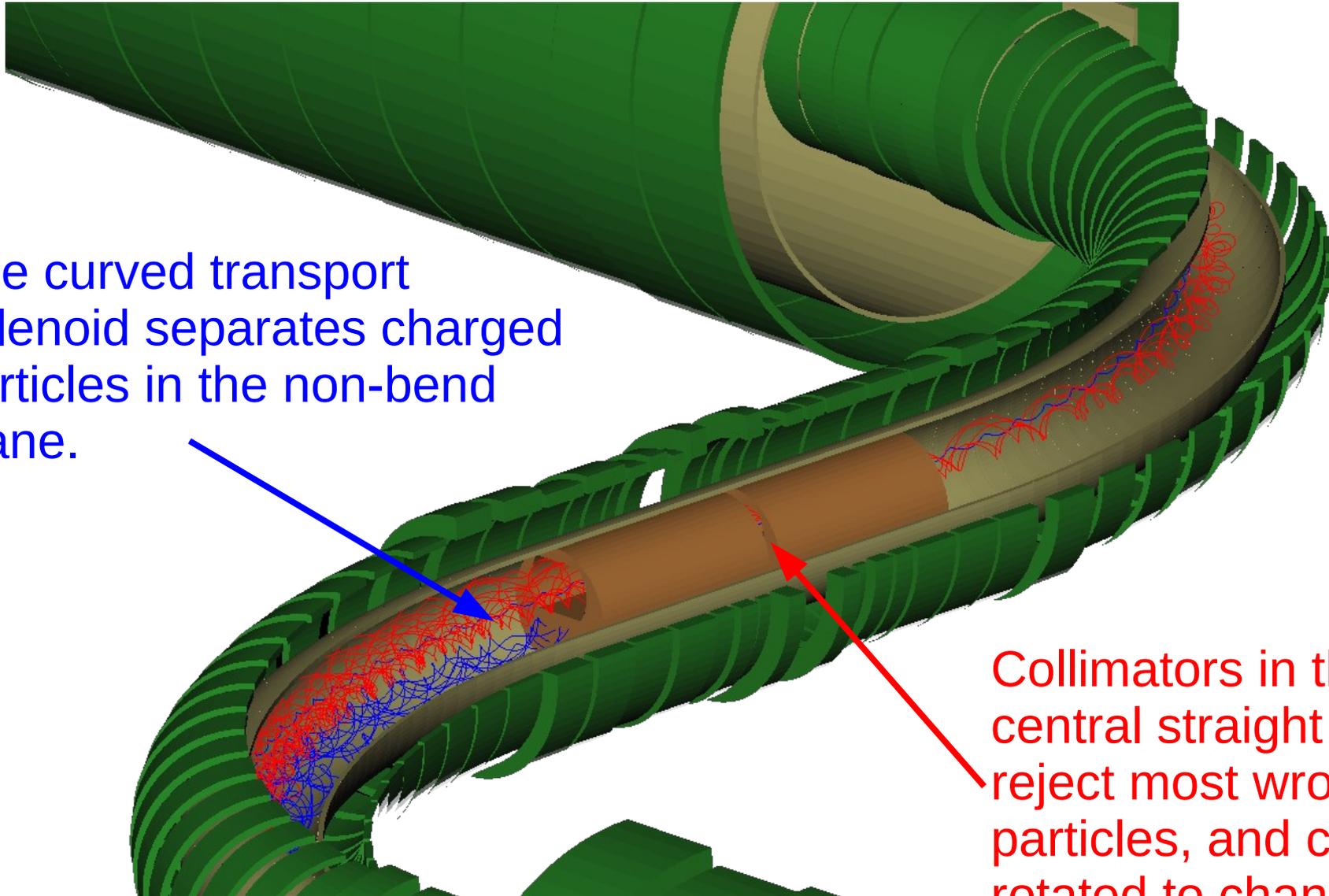


The design of the solenoid systems is a collaborative effort among TD, AD, and Physicists: “It's not just a big gizmo for TD to build ... we need University people to get involved!”

Transport solenoid

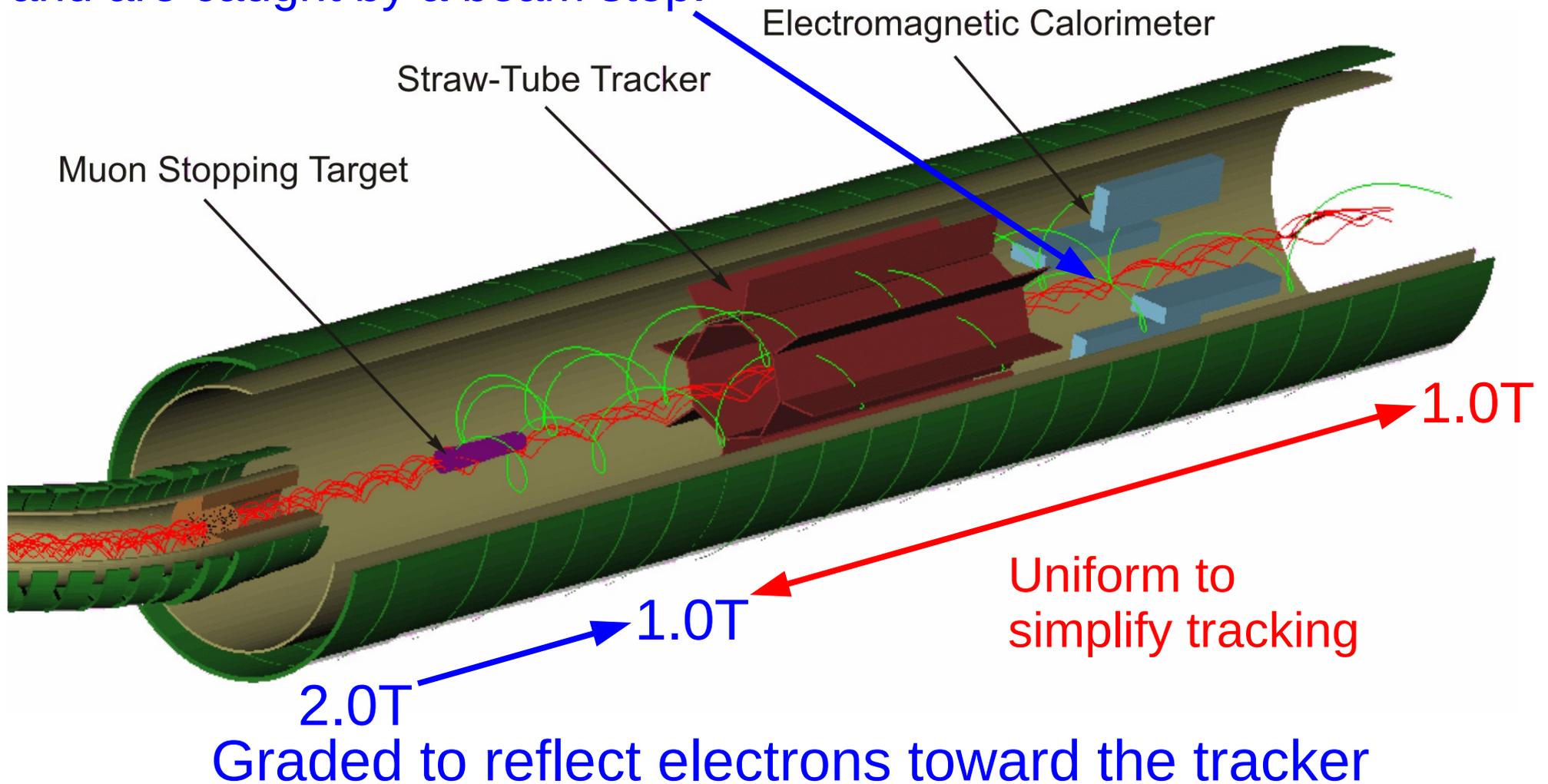
The curved transport solenoid separates charged particles in the non-bend plane.

Collimators in the central straight section reject most wrong sign particles, and can be rotated to change sign for calibration runs.

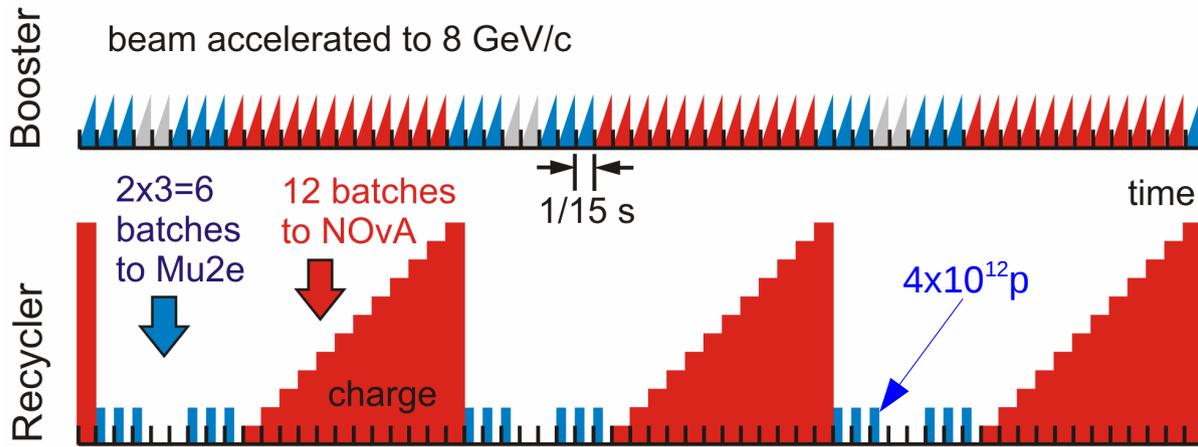


The detector solenoid and stopping target

The remnant muon beam and most DIO electrons pass through the central openings, and are caught by a beam stop.

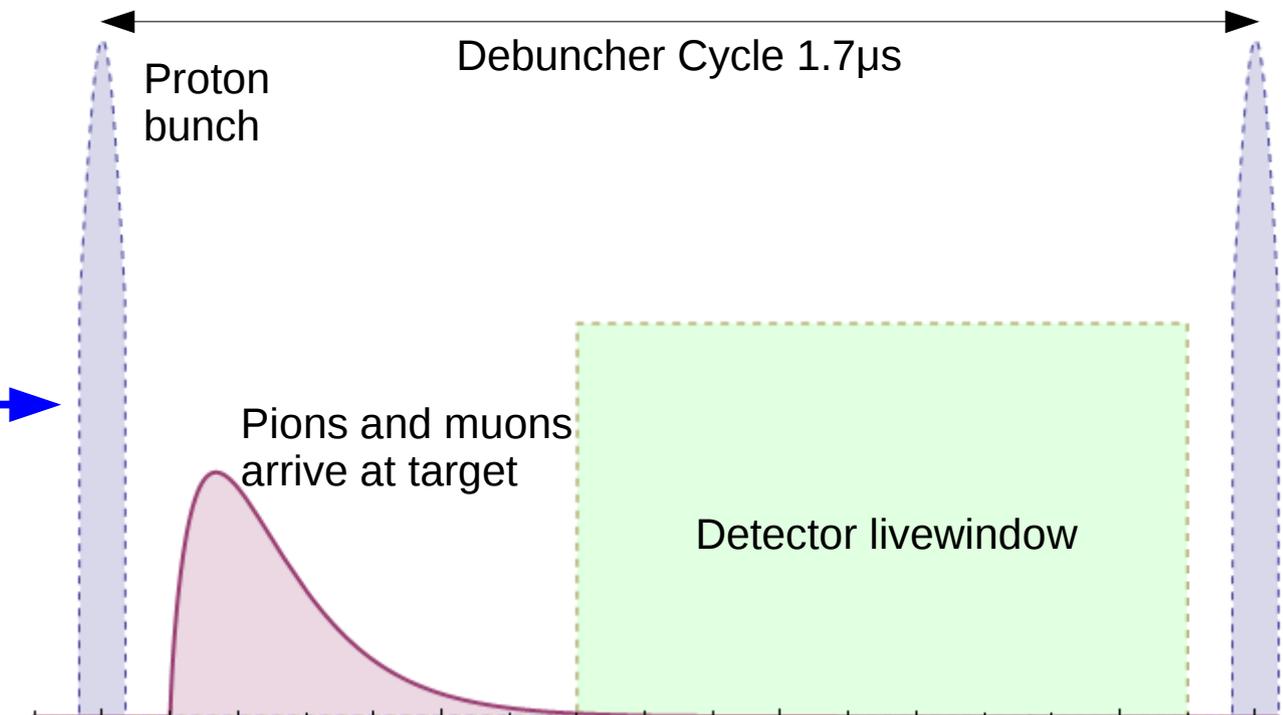


We will utilize *unused* Booster batches during Nova running

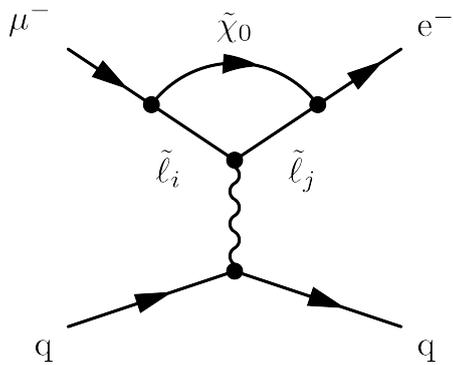


Mu2e running does not impact the neutrino program!

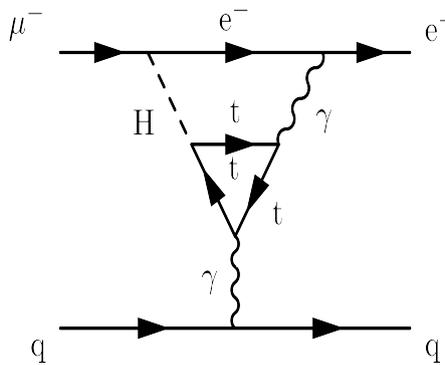
These batches are stacked and bunched in the accumulator, then transferred to the debuncher for slow extraction to Mu2e



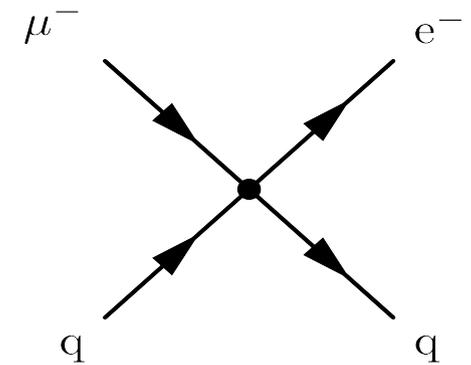
Muon conversion arises in many new physics scenarios



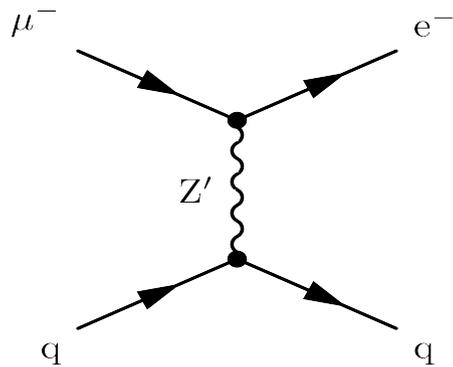
SUSY



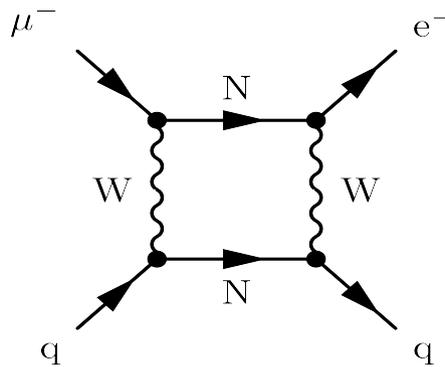
Second Higgs Doublet



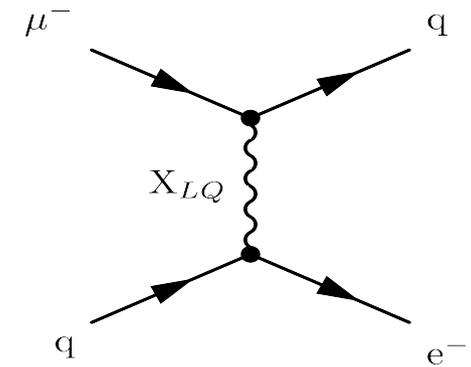
Compositeness



Heavy Gauge Bosons



Heavy Neutrinos

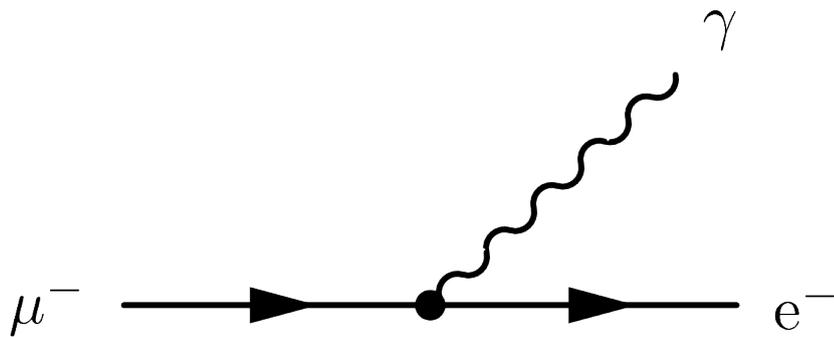


Leptoquarks

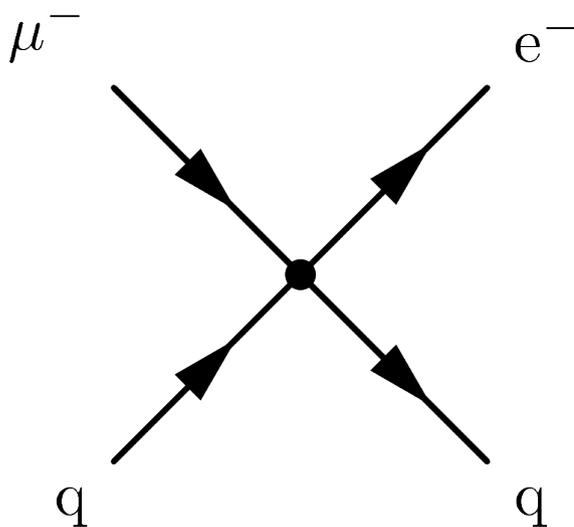
The discovery of Weak scale SUSY at LHC would imply observable cLFV rates

$$R_{\mu e} \sim 10^{-15}$$

We can parameterize all these models in terms of two effective interactions



Active at a common energy scale Λ



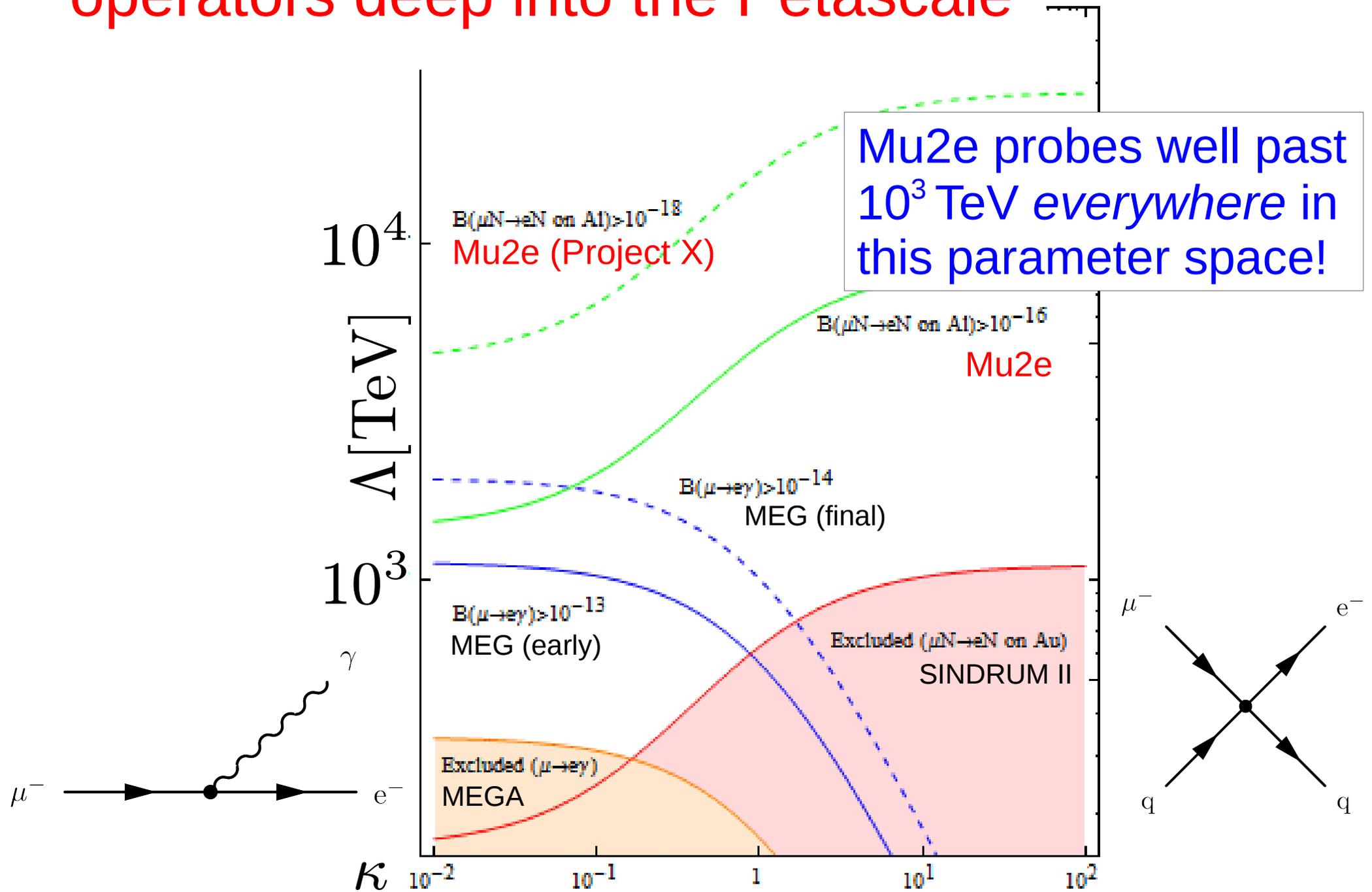
A dimension 5 dipole interaction

And we'll set their relative strength with a dimensionless interpolating factor κ

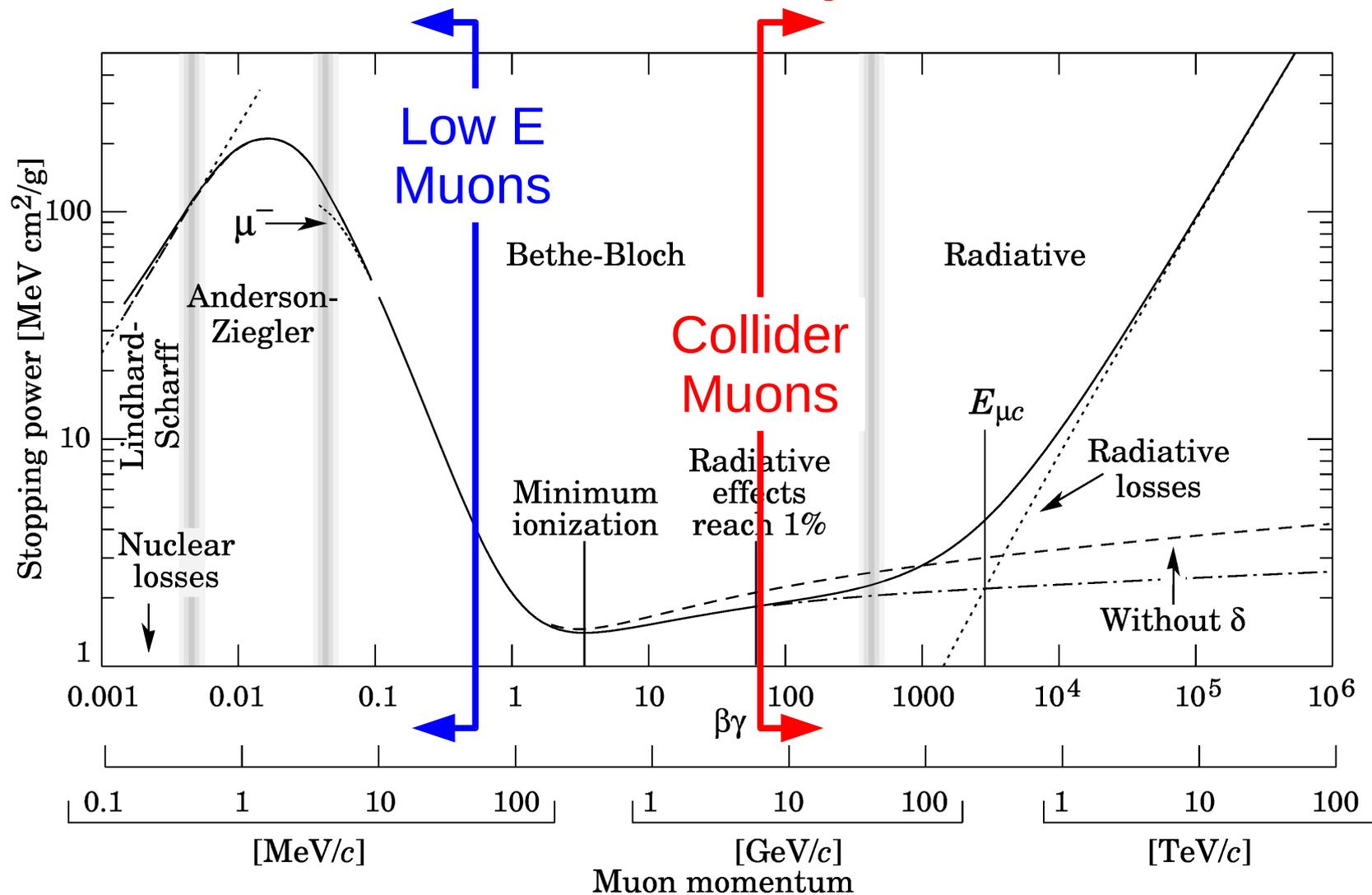
A dimension 6 contact interaction

$$\frac{1}{\kappa + 1}$$
$$\frac{\kappa}{\kappa + 1}$$

Mu2e probes combinations of these operators deep into the Petascale



Our muons are not like your muons...



Negative muons rapidly stop, capture on atoms, and cascade to the 1s state.

$$E_{1s} = -Z^2 2.7 \text{ keV}$$