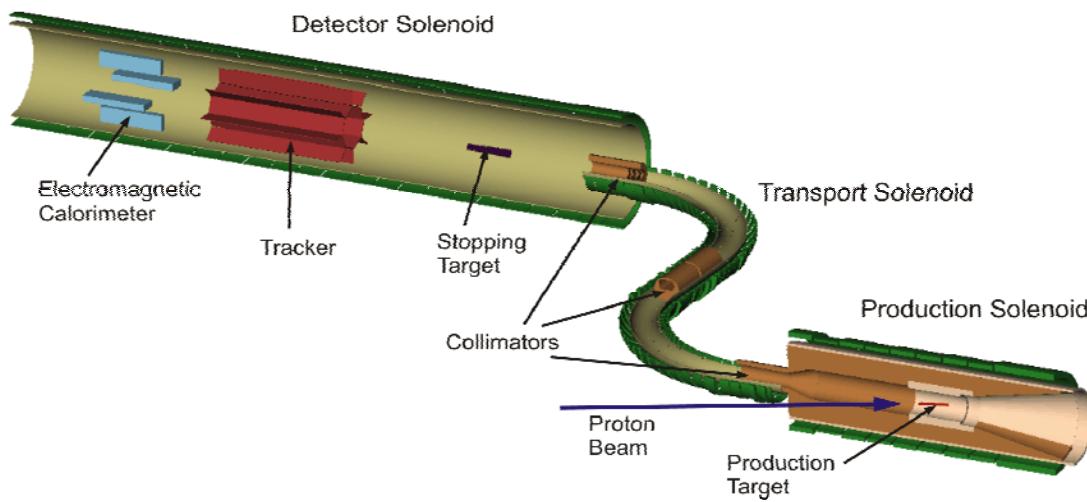


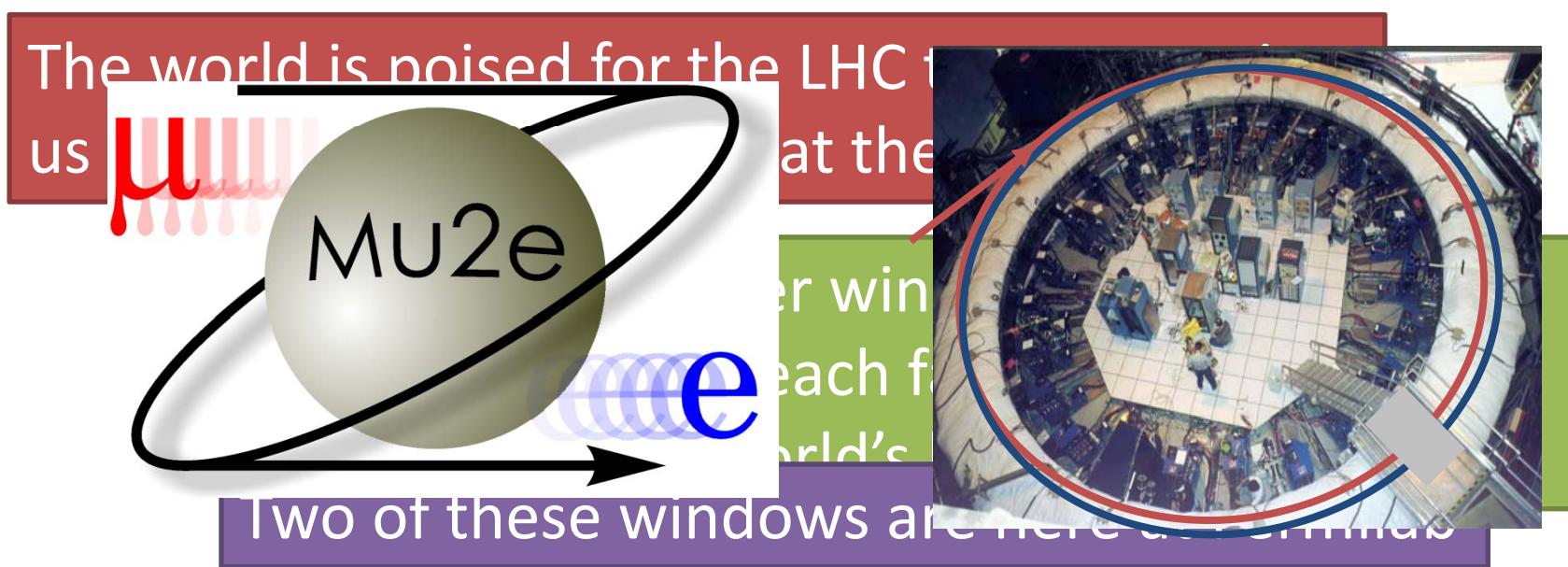
The Mu2e and g-2 Experiments

A precision window into physics
beyond the standard model



FNAL Users Meeting 2009

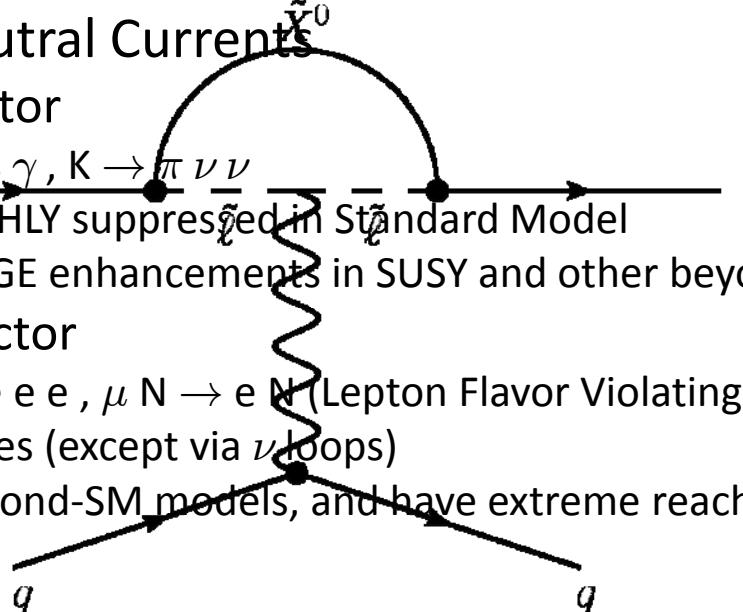
Introduction



Mu2e and g-2 will probe new physics, new models and new energy scales through precision measurements and ultra rare searches

Why look for Ultra-Rare Processes?

- We want to access beyond standard model physics
 - This means access to High and Ultra-High Energy interactions
 - We get to these energies through loops, and this means rare effects
- Ideally we start with processes that are forbidden or highly suppressed in the standard model
 - Any observation becomes proof of non-SM physics
- Flavor Changing Neutral Currents
 - FCNC in quark sector
 - $B_s \rightarrow \mu \mu$, $b \rightarrow s \gamma$, $K \rightarrow \pi \nu \bar{\nu}$
 - Allowed but HIGHLY suppressed in Standard Model
 - Can receive LARGE enhancements in SUSY and other beyond-SM physics
 - FCNC in lepton sector
 - $\mu \rightarrow e \gamma$, $\mu \rightarrow e e e$, $\mu N \rightarrow e N$ (Lepton Flavor Violating)
 - No SM amplitudes (except via ν loops)
 - Permitted in beyond-SM models, and have extreme reach in energy



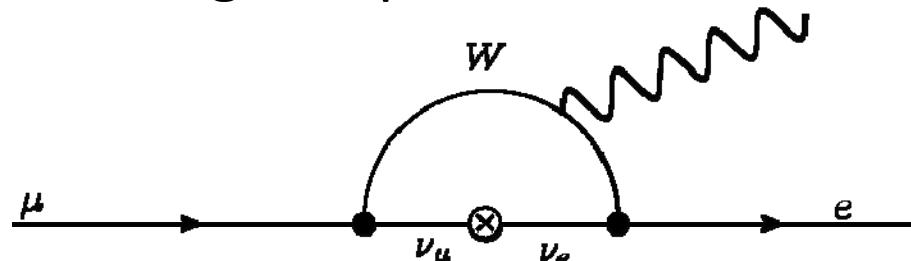
Lepton Mixing in the Standard Model

- We have three generations of leptons:

$$\left(\begin{array}{c} e \\ \nu_e \end{array} \right) \left(\begin{array}{c} \mu \\ \nu_\mu \end{array} \right) \left(\begin{array}{c} \tau \\ \nu_\tau \end{array} \right)$$

No SM couplings between generation!

- In the standard model Lagrangian there is no explicit mixing between generations
- But we have explicitly observed neutrino oscillations, and this means that leptons DO mix.
- Charged leptons must mix through neutrino loops



$$B(\mu \rightarrow e\gamma) = \frac{3\alpha}{32\pi} \left| \sum_\ell V_{\mu\ell}^* V_{e\ell} \frac{m_{\nu_\ell}^2}{M_W^2} \right| \simeq 10^{-54}$$

- But the mixing is so small, it's effectively forbidden

The Three LFV Processes

- There are three basic channels to search for cLFV in:
 $\mu^\pm \rightarrow e^\pm \gamma$
 $\mu^\pm \rightarrow e^\pm e^\pm e^\mp$
 $\mu^- N \rightarrow e^- N$

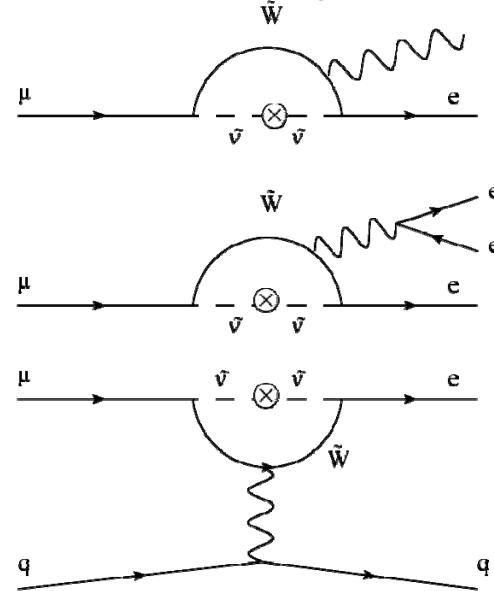
- Where if dipole like interactions dominate we expect a ratio of rates:

389 to 2.3 to 1

Note: $\mu \rightarrow e\gamma$ and $\mu \rightarrow eee$ have experimental limitations from resolution and overlap with accidentals.

Limits the overall $\text{Br}() \sim 10^{-14}$

- New physics for all can come from loop level

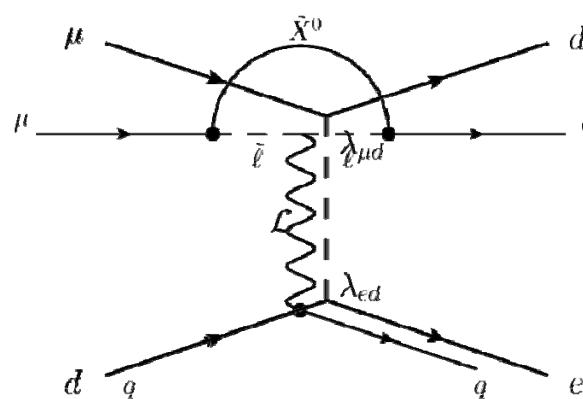


- For $\mu N \rightarrow e N$ and $\mu \rightarrow eee$ we also have contact terms

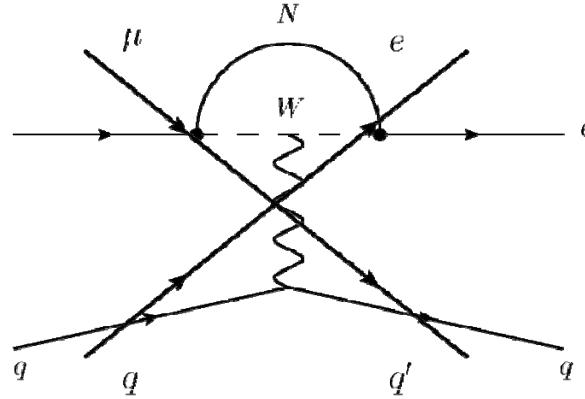
Beyond the Standard Model

- The LFV process can manifest in the $\mu N \rightarrow e N$ channel in many models:

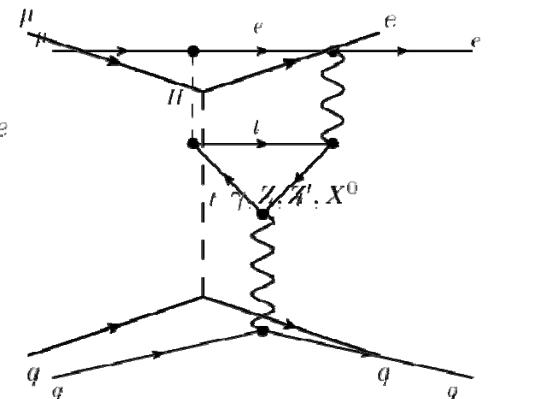
Contact Terms



Leptoquarks



Heavy Neutralinos



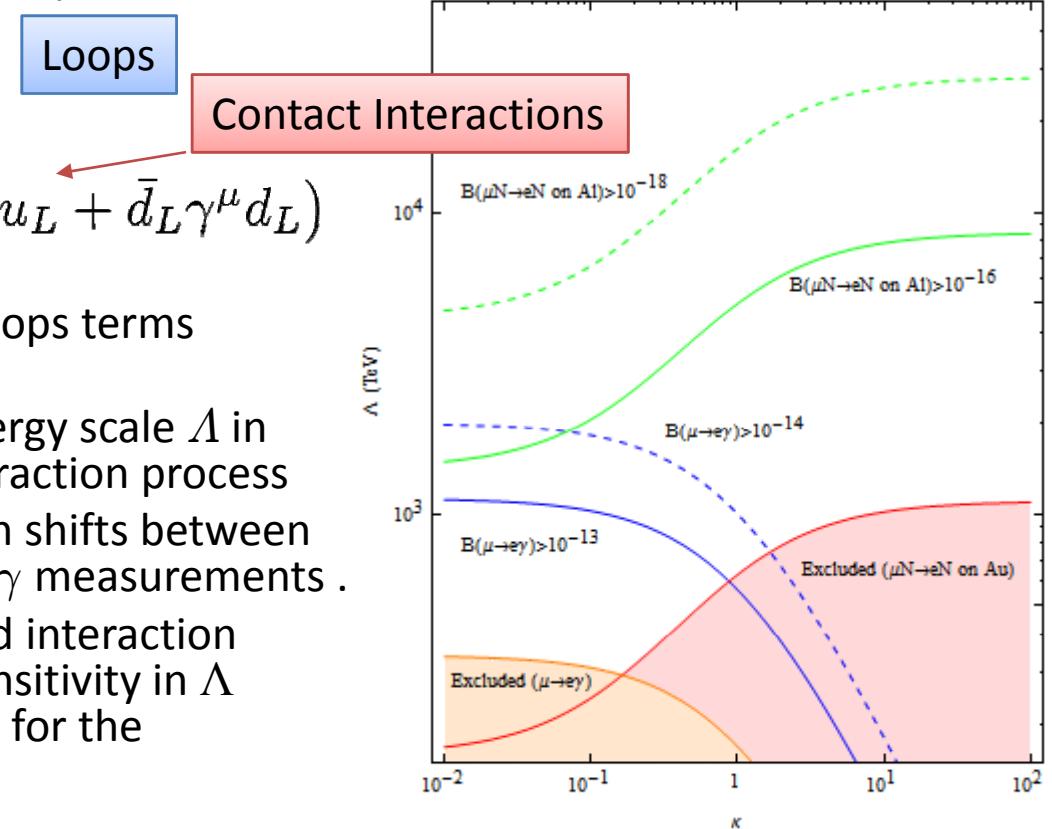
Another Higgs Doublet

General LFV Lagrangian

- Recharacterize as a model independent frame work:

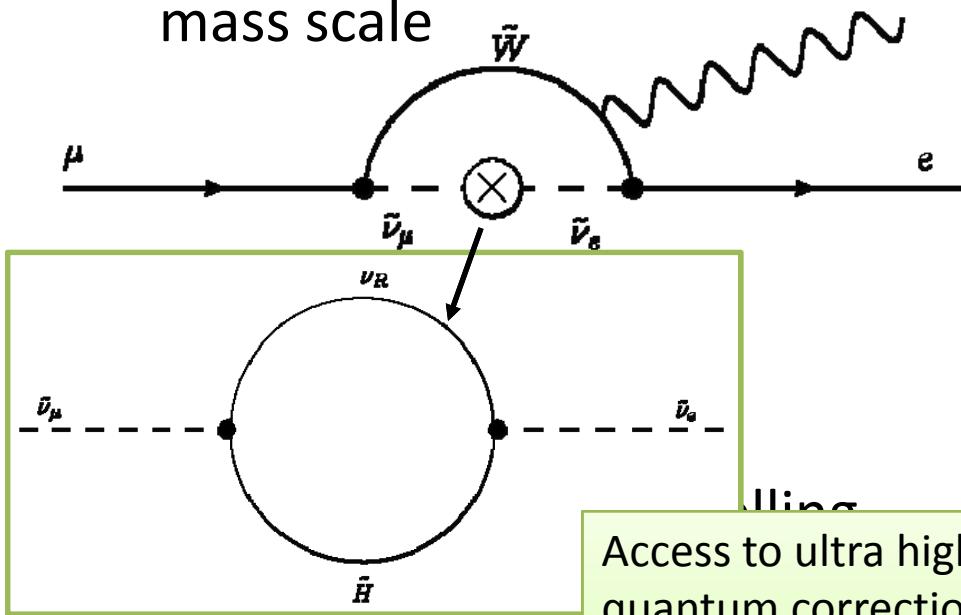
$$\mathcal{L}_{LFV} = \frac{m_\mu}{(\kappa+1)\Lambda^2} \bar{\mu}_R \sigma_{\mu\nu} e_L F^{\mu\nu} + \frac{\kappa}{(1+\kappa)\Lambda^2} \bar{\mu}_L \gamma_\mu e_L (\bar{u}_L \gamma^\mu u_L + \bar{d}_L \gamma^\mu d_L)$$

- Splits LFV sensitivity into Loops terms and Contact terms
- Allows us to look at the energy scale Λ in terms of the dominant interaction process
- The balance in energy reach shifts between favoring $\mu N \rightarrow e N$ and $\mu \rightarrow e \gamma$ measurements .
- For contact term dominated interaction (large κ) the energy the sensitivity in Λ reaches upwards of 10^4 TeV for the conversion process



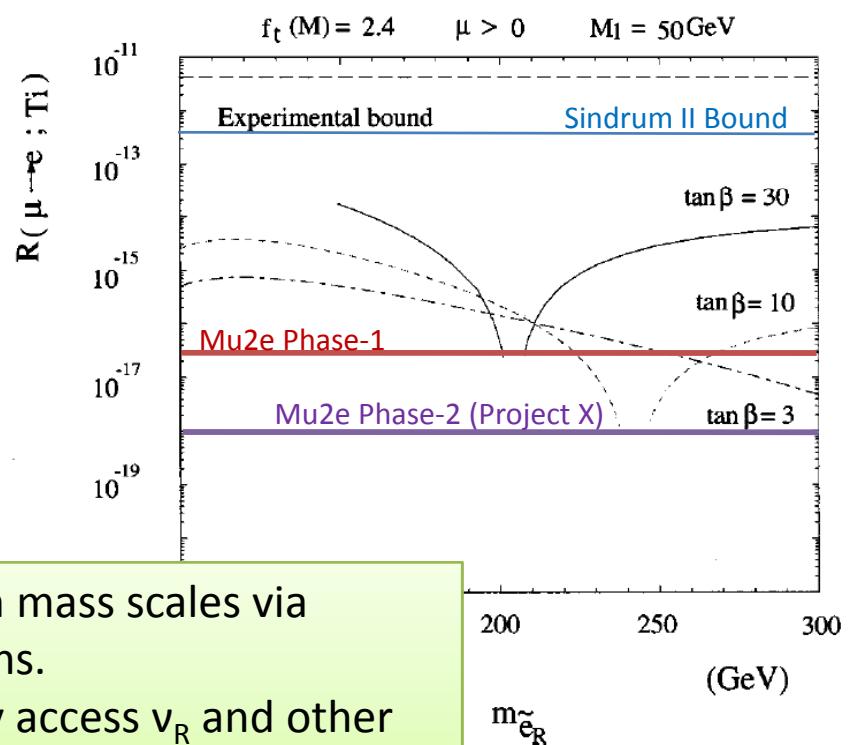
Sensitivity to SUSY

- Rates are not small because they are set by the SUSY mass scale



mean observation
 $\approx \mathcal{O}(40)$ events

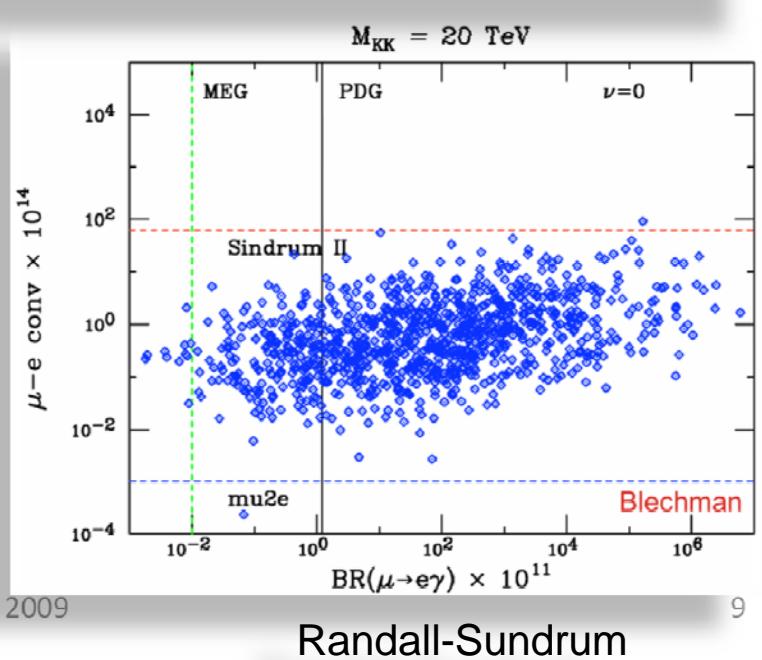
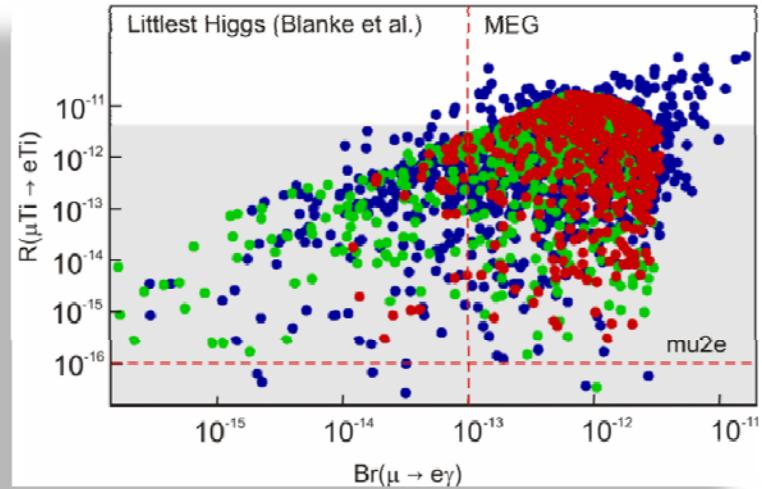
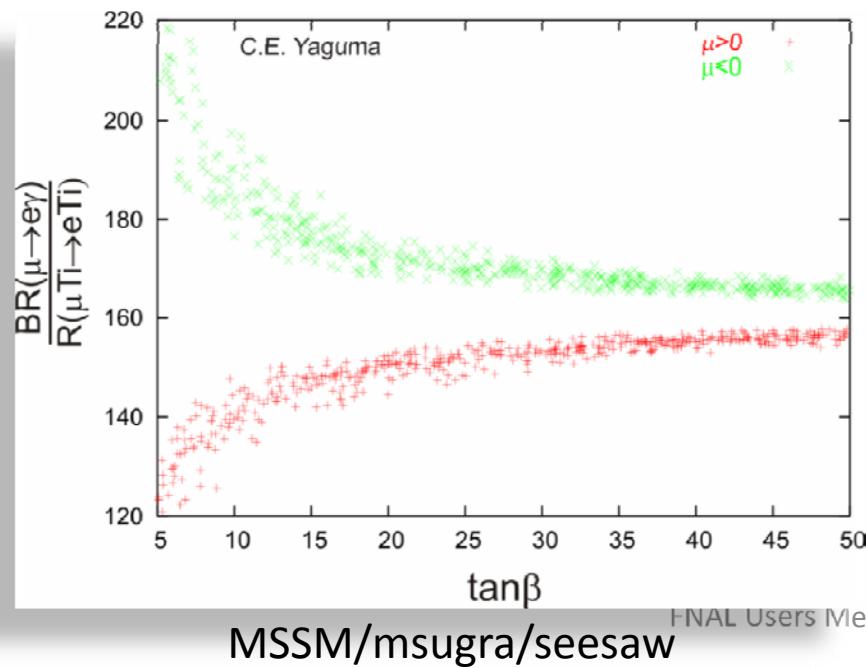
Access to ultra high mass scales via quantum corrections.
 Can access possibly access v_R and other processes at scales 10^{12} - 10^{14} GeV/c 2



Hisano et al. 1997

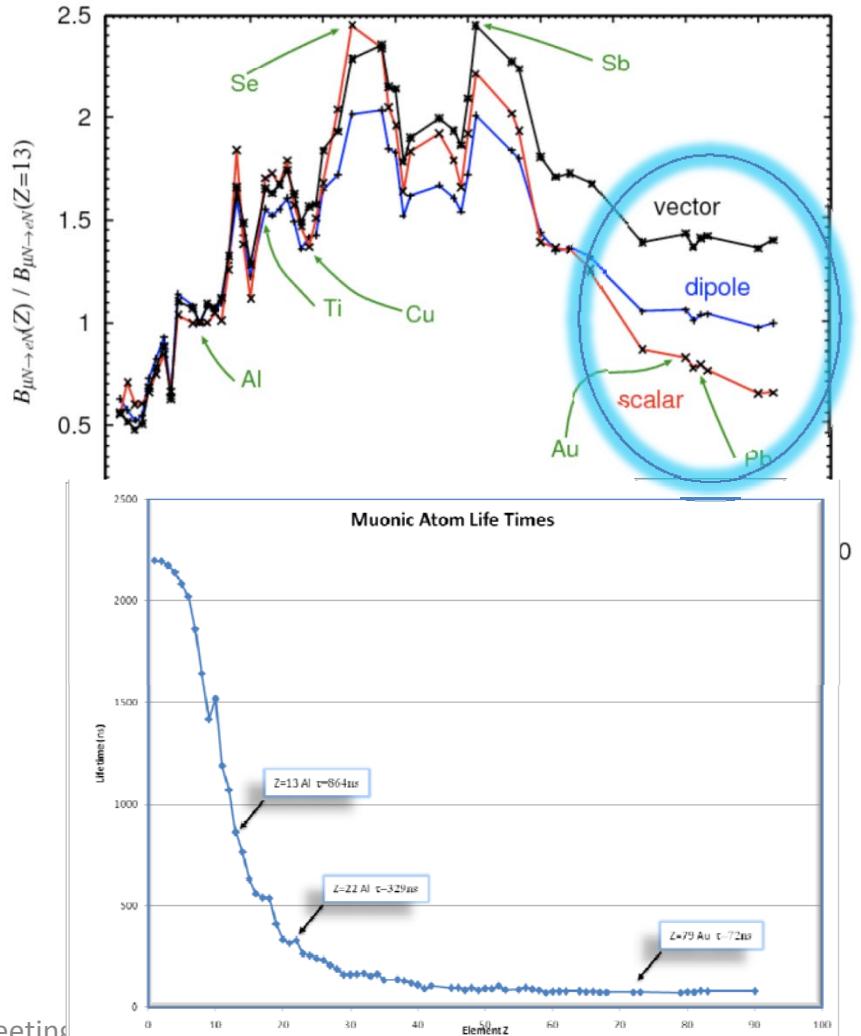
We Need Both $\mu N \rightarrow e N$ and $\mu \rightarrow e\gamma$

- Knowing both mu2e and mu2e gamma gives us knowledge about the structure of SUSY



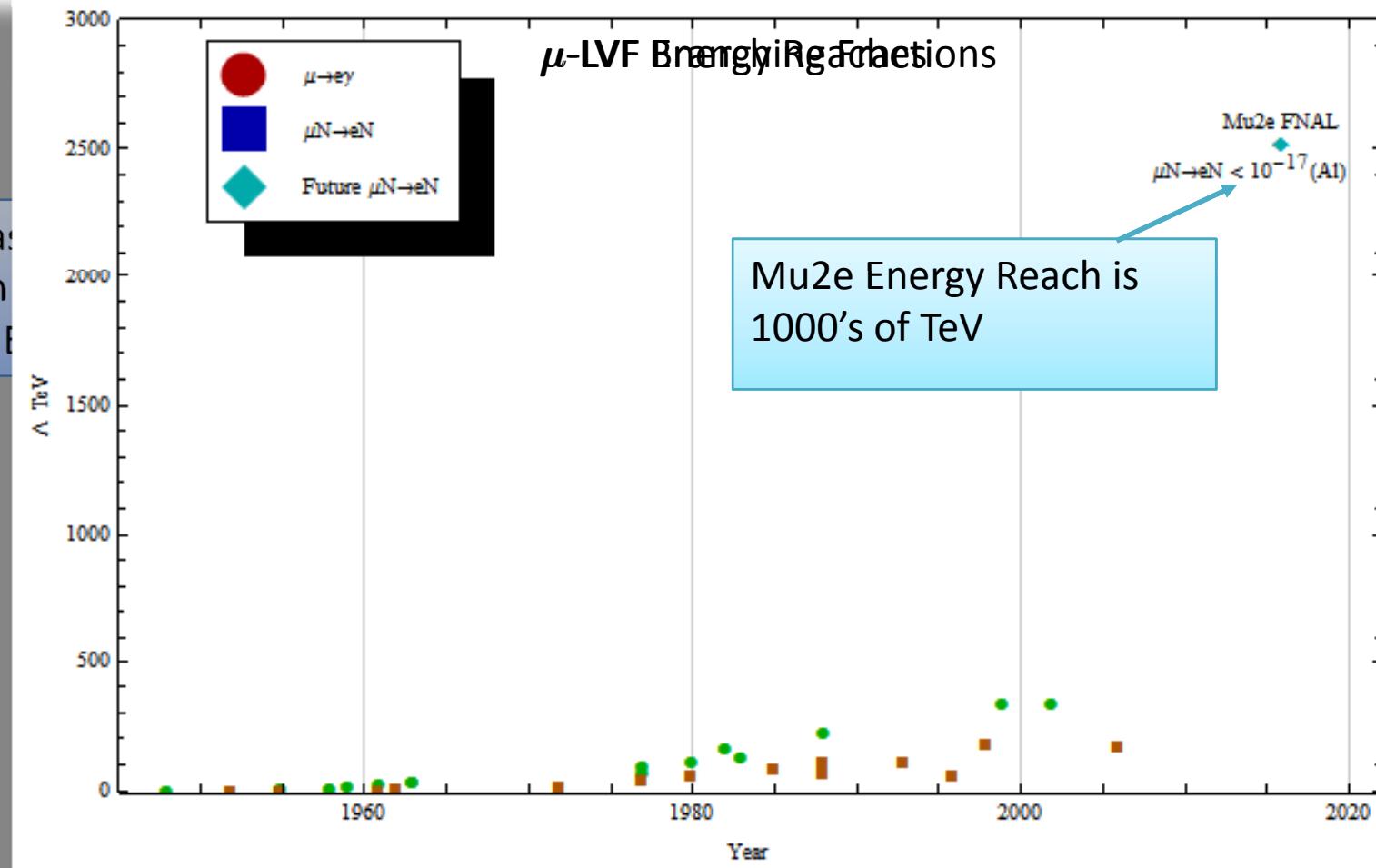
$\mu N \rightarrow eN$ & SUSY Models

- Assuming we see a signal:
 - By changing target, we gain sensitivity to the scalar, vector or dipole nature of the interaction
 - Need to go to high Z
 - Hard because τ small for large Z ($\tau_{Au} = 72\text{ns}$)
 - But DIO backgrounds are suppressed and Conversion/OMC ratio scales as Z
- This is a unique feature of the $\mu N \rightarrow eN$ measurements



A Brief History of μ -LFV

First Measurement
 6×10^{-2} in
Effective Width



Part I

MU2E AT FNAL

The $\mu N \rightarrow e N$ measurement (in a nutshell)

- Stop $\sim \mathcal{O}(10^{18}) \mu^-$ on a target (Al, Ti, Au)
- Wait 700ns (to let prompt backgrounds clear)
- Look for the coherent conversion of a muon to a mono-energetic electron:

$$E_e = M_\mu - N_{recoil} - (B.E.)_\mu^{1S} = 104.96 \text{ MeV (on } {}^{27}\text{Al)}$$

- Report the rate relative to nuclear capture

$$\mathcal{R} = \frac{\Gamma(\mu^- N \rightarrow e^- N)}{\Gamma(\mu^- N^Z \rightarrow \nu_\mu N^{Z-1})}$$

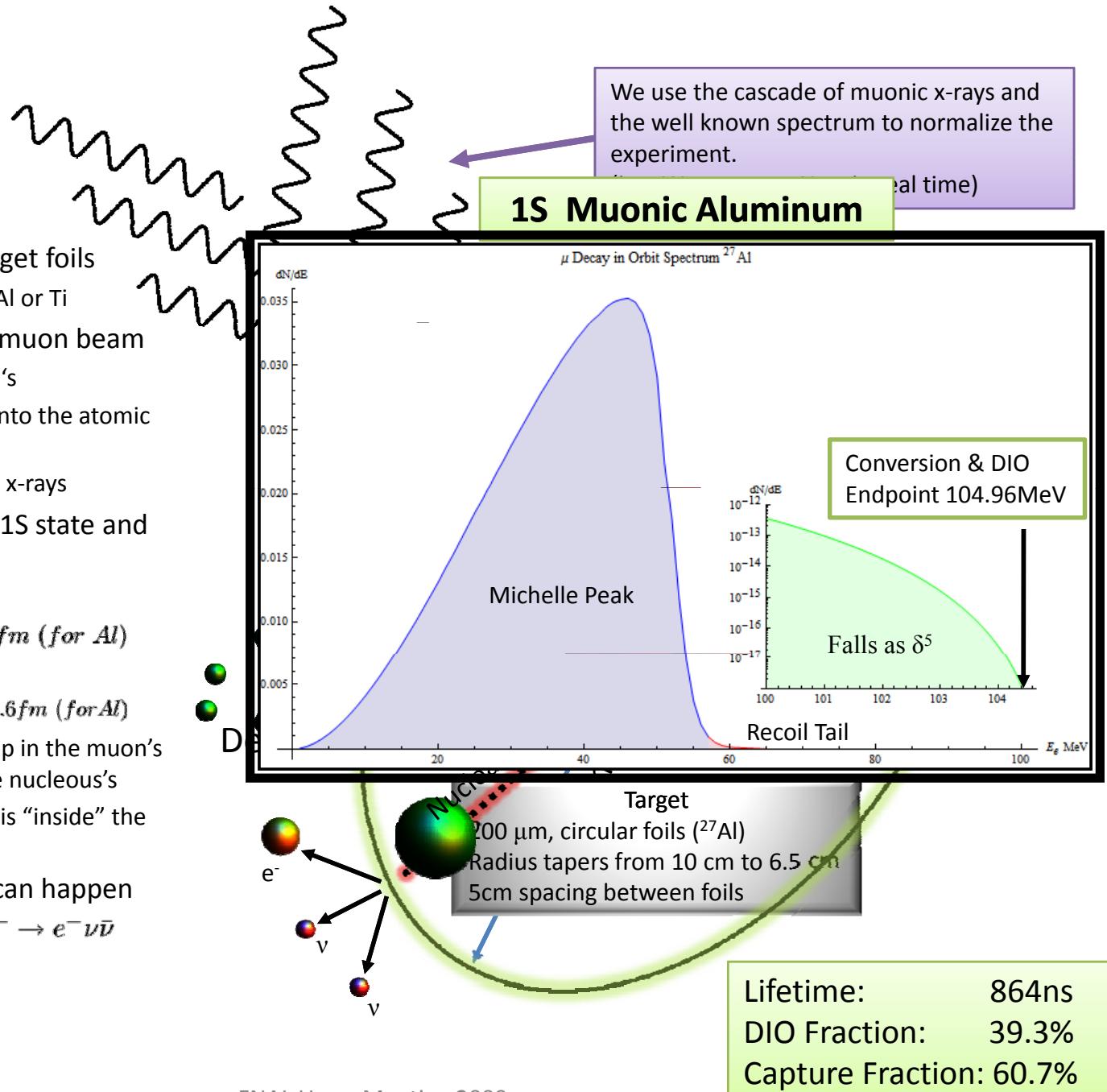
- *If we see a signal, it's compelling evidence for physics beyond the standard model!*

Muonic Atom

- Start with a series of target foils
 - For Mu2E these are Al or Ti
- Bring in the low energy muon beam
 - We stop $\approx 50\%$ of μ^- 's
 - Stopped muons fall into the atomic potential
 - As they do they emit x-rays
- Muons fall down to the 1S state and are captured in the orbit
 - Muonic Bohr Radius

$$\langle r_\mu \rangle = \frac{n^2 \hbar}{m_\mu z e^2} \approx 19.6 \text{ fm (for Al)}$$
 - Nuclear Size

$$R \approx 1.2 A^{1/3} \text{ fm} = 3.6 \text{ fm (for Al)}$$
 - Provides large overlap in the muon's wavefunction with the nucleous's
 - For $Z > 25$ the muon is "inside" the nucleous
- Once captured 3 things can happen
 - Decay in Orbit: $\mu^- \rightarrow e^- \nu \bar{\nu}$



Muonic Atom

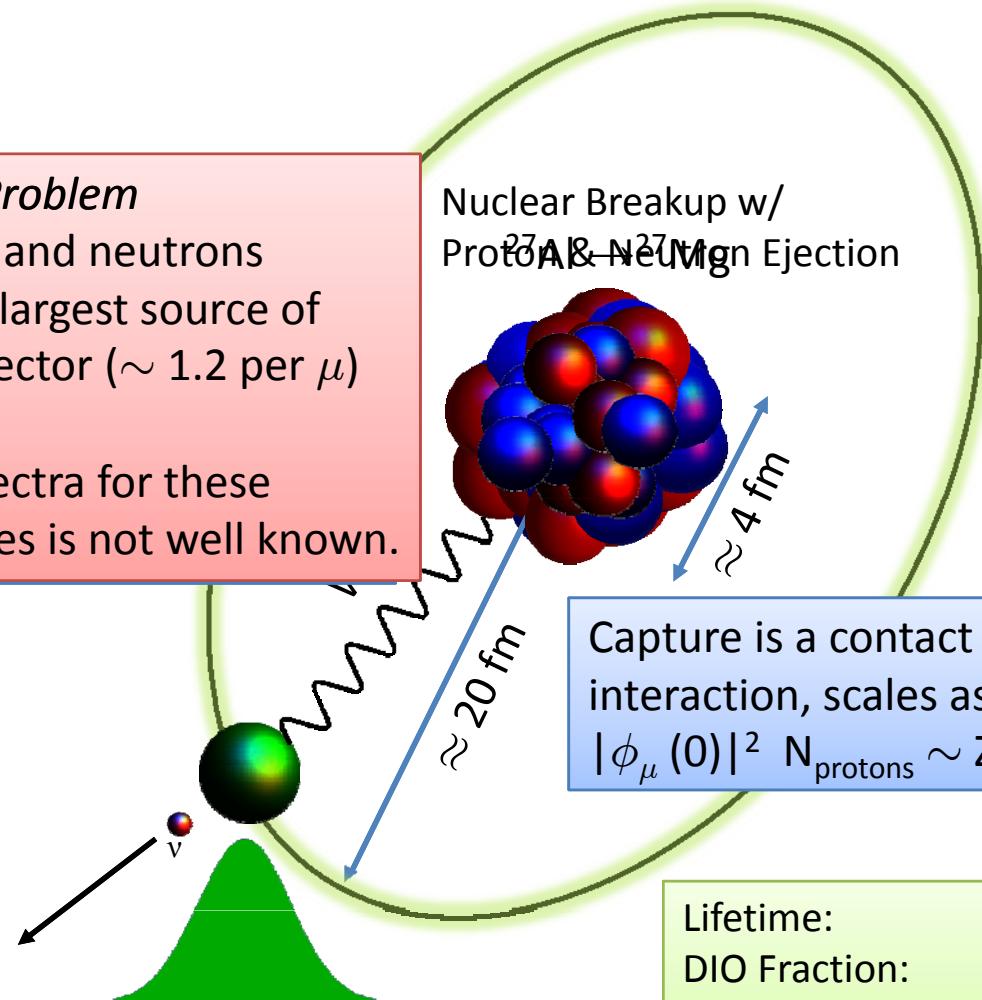
- Start with a series of target foils
 - We stop $\approx 50\%$ of μ^- 's
- Bring in the low energy muon beam
 - We stop $\approx 50\%$
 - Stopped muons have potential
 - As they do they fall down to a captured in the orbit
- Muons fall down to a captured in the orbit
 - Muonic Bohr Radius: $\langle r_\mu \rangle = \frac{n^2 \hbar}{m_\mu z e^2} \approx 1 fm$
 - Nuclear Size: $R \approx 1.2 A^{1/3} fm$

Problem
 These protons and neutrons constitute the largest source of rate in the detector (~ 1.2 per μ^-)

The energy spectra for these ejected particles is not well known.

- Once captured 3 things can happen
 - Decay in Orbit: $\mu^- \rightarrow e^- \nu \bar{\nu}$
 - Nuclear Capture: $\mu^- N^Z \rightarrow \nu N^{Z-1}$

Ordinary Muon Capture (OMC)



Muonic Atom

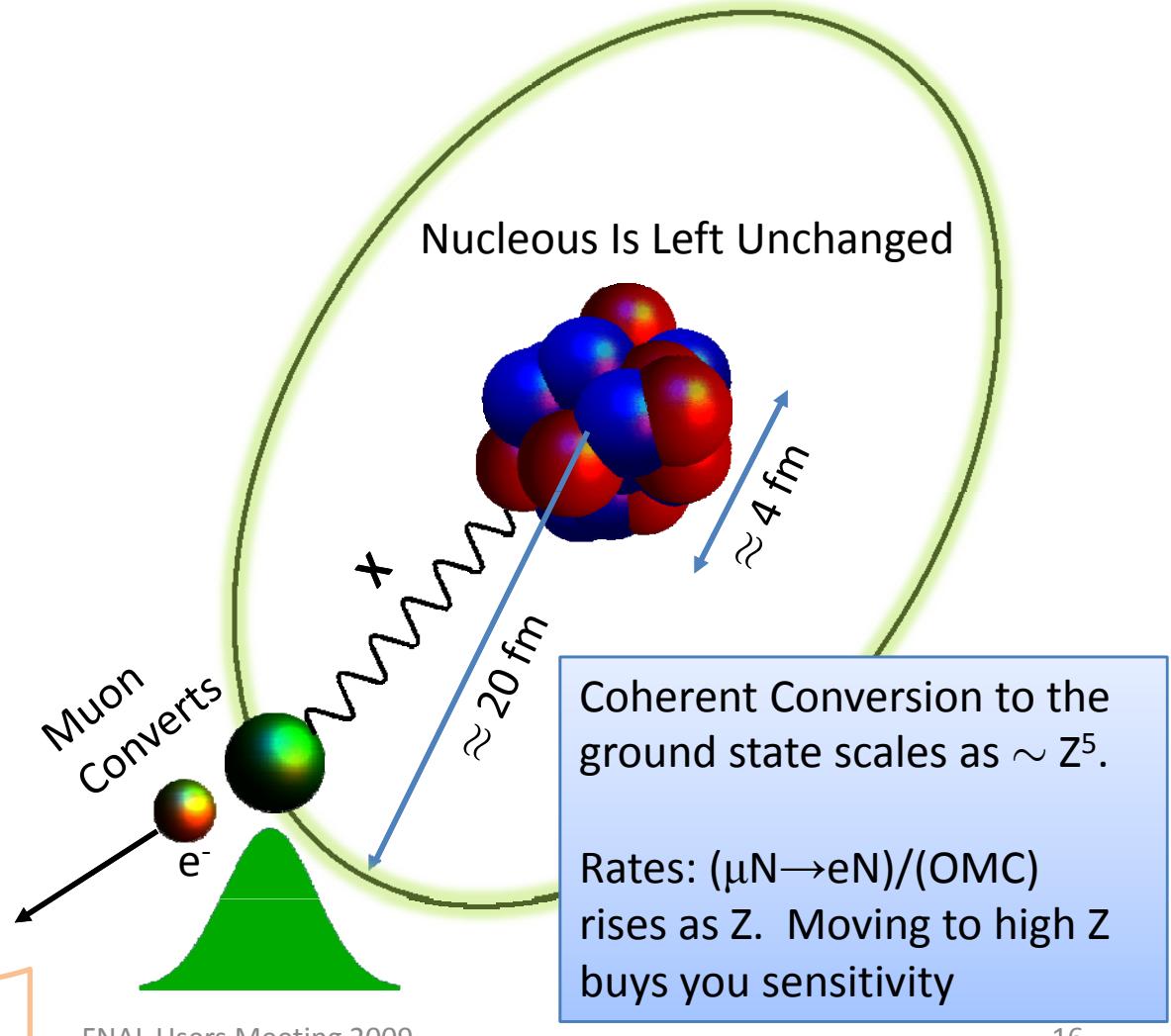
- Start with a series of target foils
 - We stop $\approx 50\%$ of μ^- 's
- Bring in the low energy muon beam
 - We stop $\approx 50\%$ of μ^- 's
 - Stopped muons fall into the atomic potential
 - As they do they emit x-rays
- Muons fall down to the 1S state and are captured in the orbit
 - Muonic Bohr Radius

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 - Provides large overlap in the muon's wavefunction with the nucleous's
 - For $Z > 25$ the muon is "inside" the nucleous
- Once captured 3 things can happen
 - Decay in Orbit: $\mu^- \rightarrow e^- \nu \bar{\nu}$
 - Nuclear Capture: $\mu^- N^Z \rightarrow \nu N^{Z-1}$
 - New Physics! i.e. $\mu^- N \rightarrow e^- N$

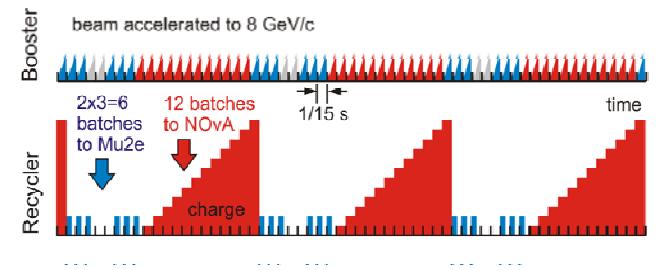
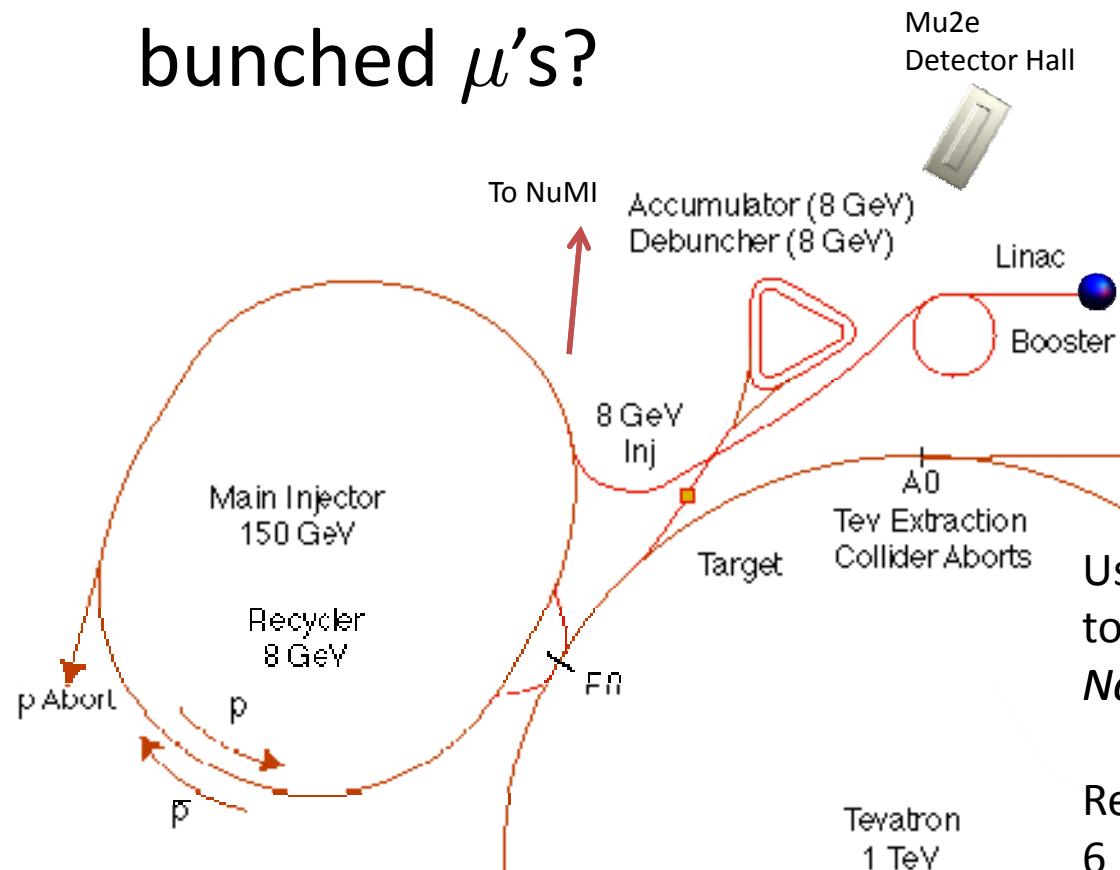
$E_e \approx 105 \text{ MeV}$

Coherent Conversion ($\mu^- \rightarrow e^-$)



Mu2E & NOvA/NuMI

- How do we deliver $\mathcal{O}(10^{18})$ bunched μ 's?



Use NuMI cycles in the Main injector
to slow spill to Mu2e.

No Impact on NOvA

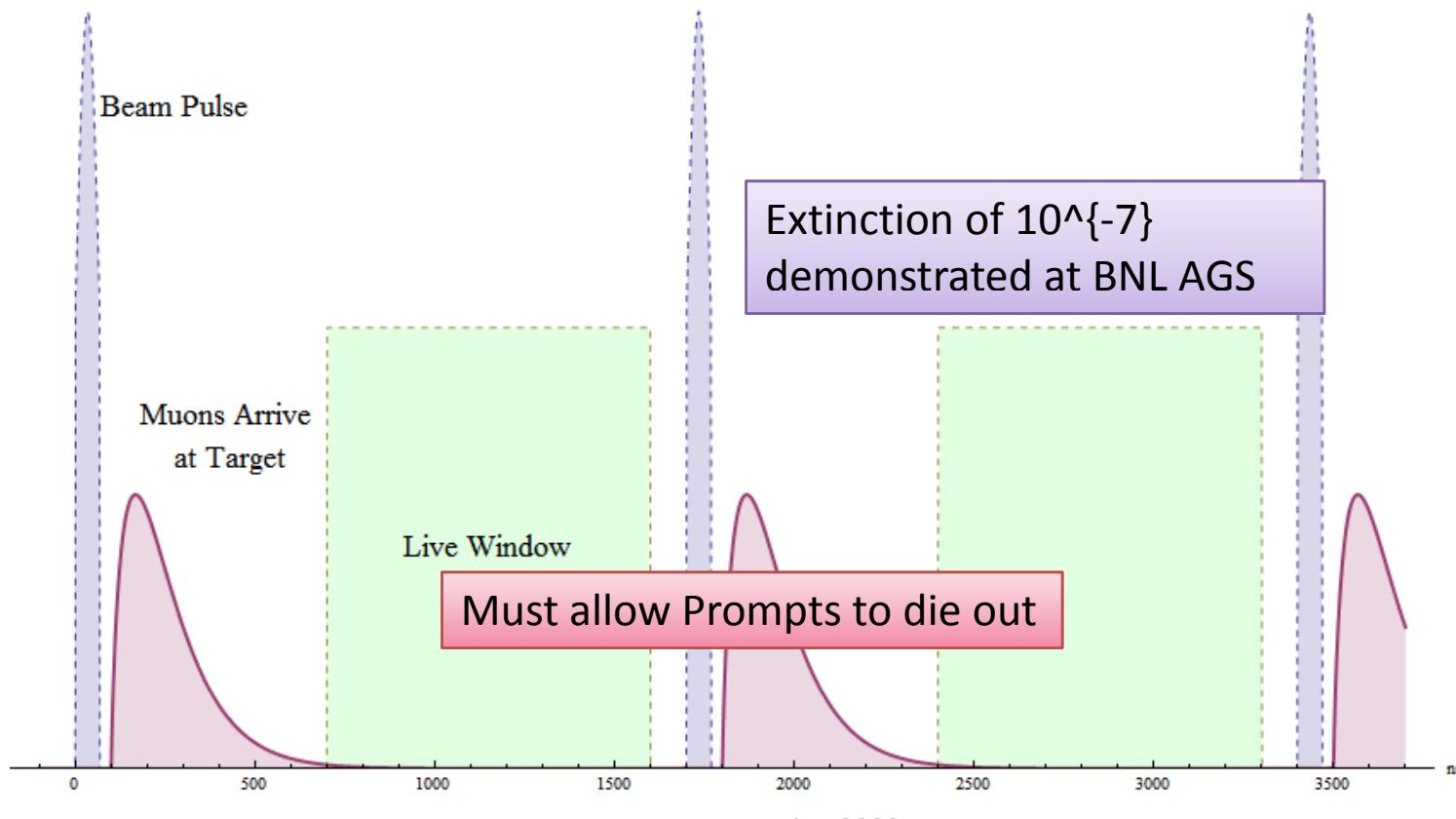
Results in:

6 batches $\times 4 \times 10^{12} / 1.33 \text{ s} \times 2 \times 10^7 \text{ s/yr}$

= **3.6×10^{20} protons/yr**

Beam Structure

- μ 's are accompanied by “prompt” e, π ,
- These cause real background
- Must limit our beam extinction, and detector live window



Backgrounds

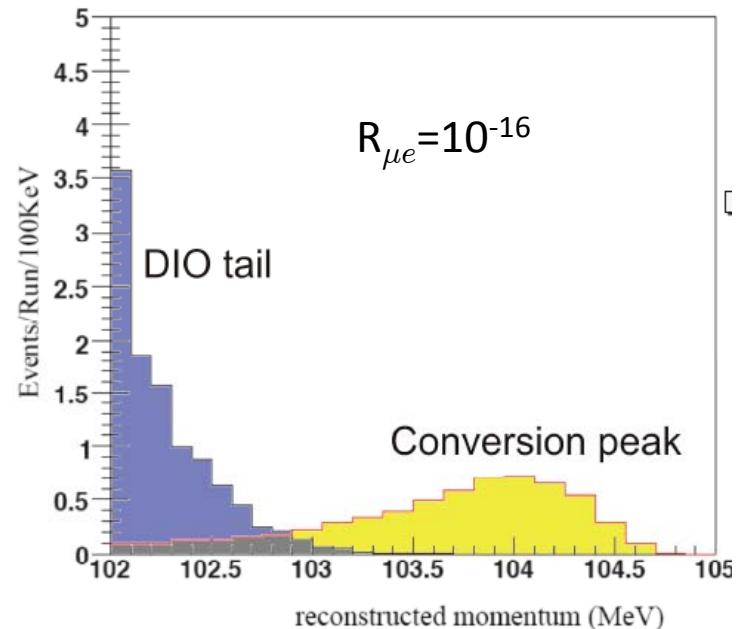
- DIO, RMC, RPC
- Dominant Backgrounds
 - Prompt backgrounds drive the extinction
 - In particular, Rad π Cap. drives the extinction requirement
 - Estimates are for an extinction of 10^{-9}
 - Monte Carlo limited on estimates

Background	Evts (2×10^{-17})
μ Decay in Orbit (DIO) Tail	0.225
μ Decay in flight w/ scatter	0.036
Beam Electrons	0.036
Cosmic Ray	0.016
μ Decay in flight (no scatter)	< 0.027
Anti-proton	0.006
Radiative μ capture	<0.002
Radiative π capture	0.001
π Decay in flight	<0.001
Pat. Recognition Errors	<0.002
Total	0.415

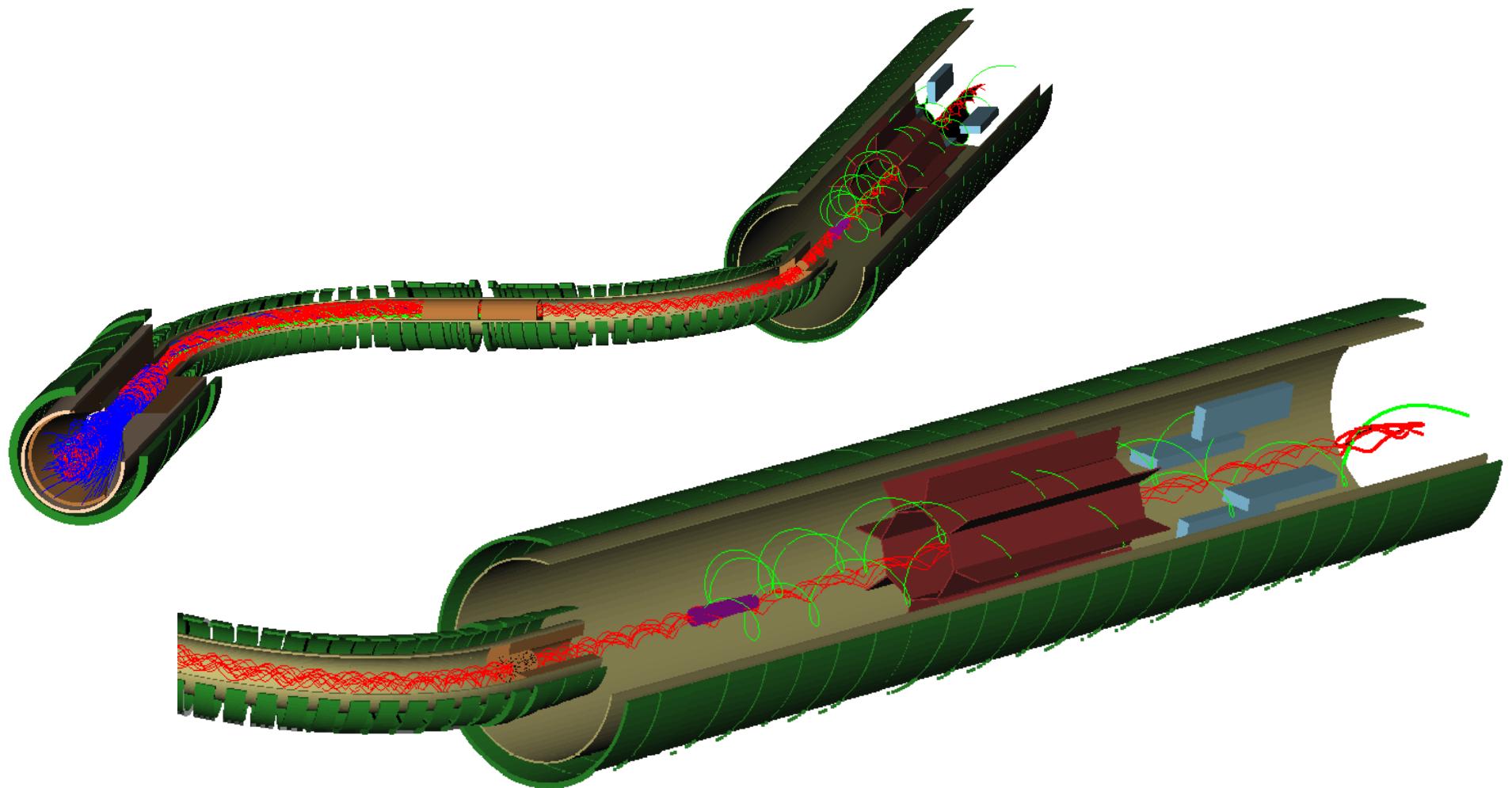
Signal to All Backgrounds

- Signal significance
 - If we assume SUSY like the LHC will see:
 - Mu2E will see $\sim \mathcal{O}(40)$ events
 - On 0.5 event background
 - Even at 10^{-15}
 - Mu2E sees ~ 5 events
 - on 0.5 event background
 - This is a Strong Signature

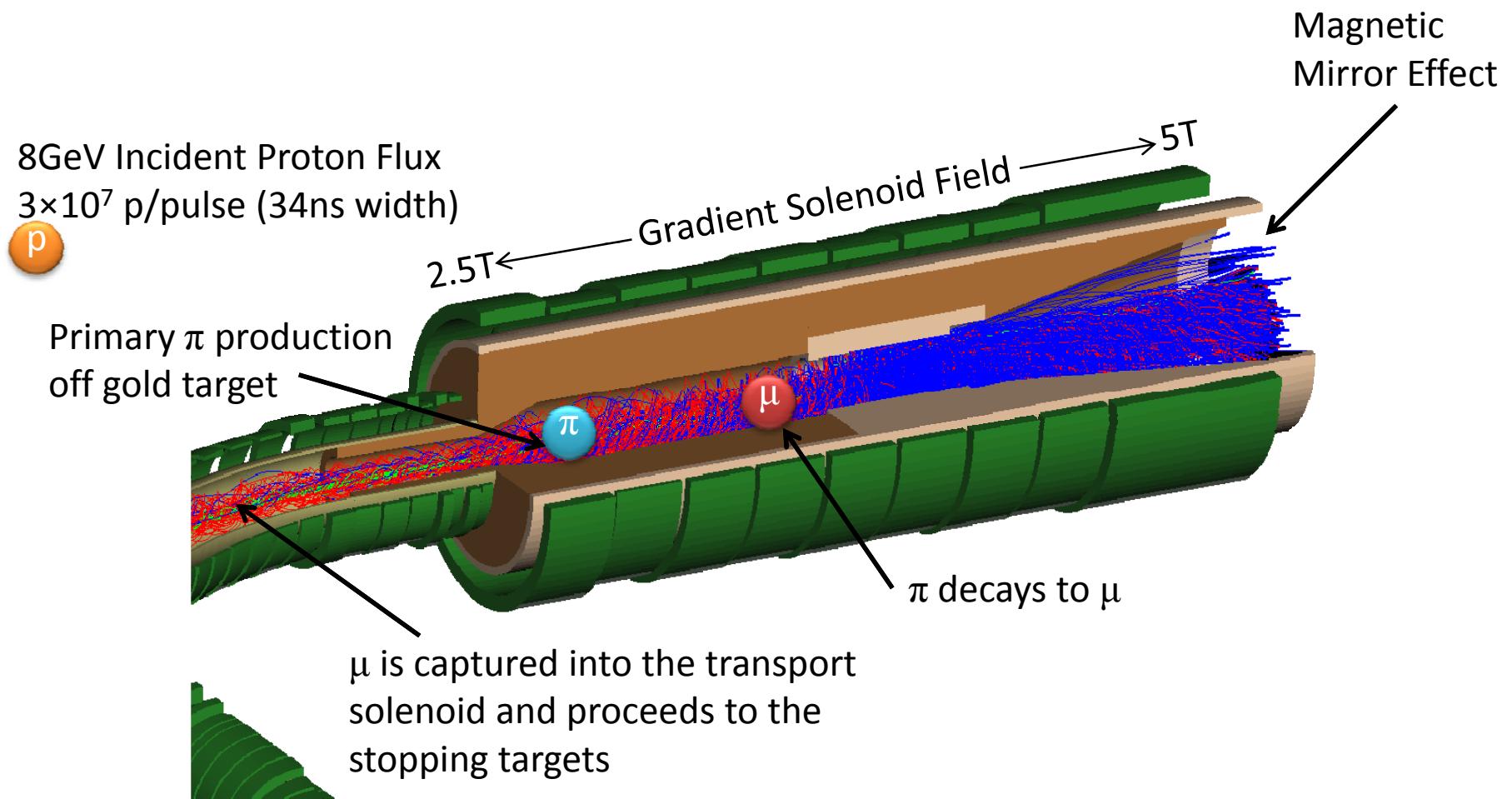
$$\frac{S}{\sqrt{B}} \sim 5.5$$



The Mu2e Detector in Detail

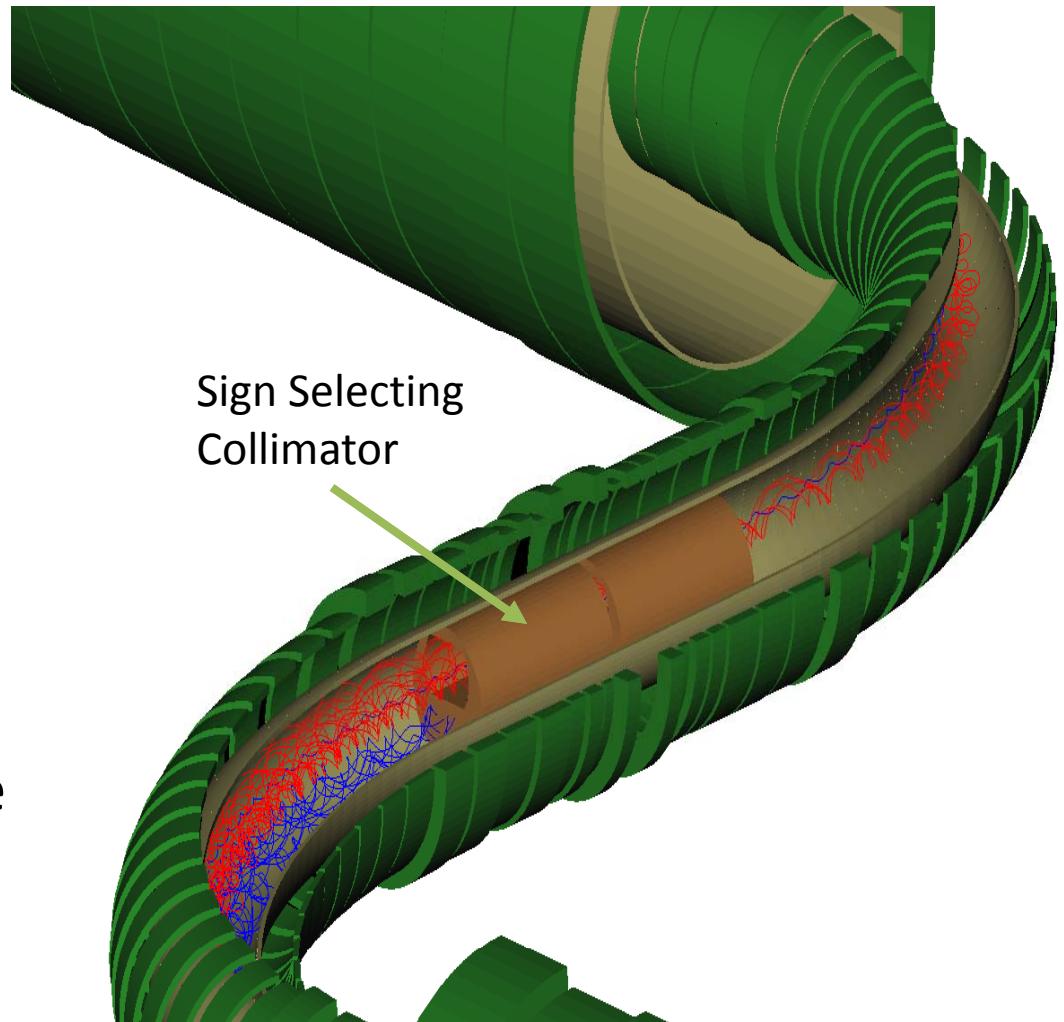


Production Solenoid

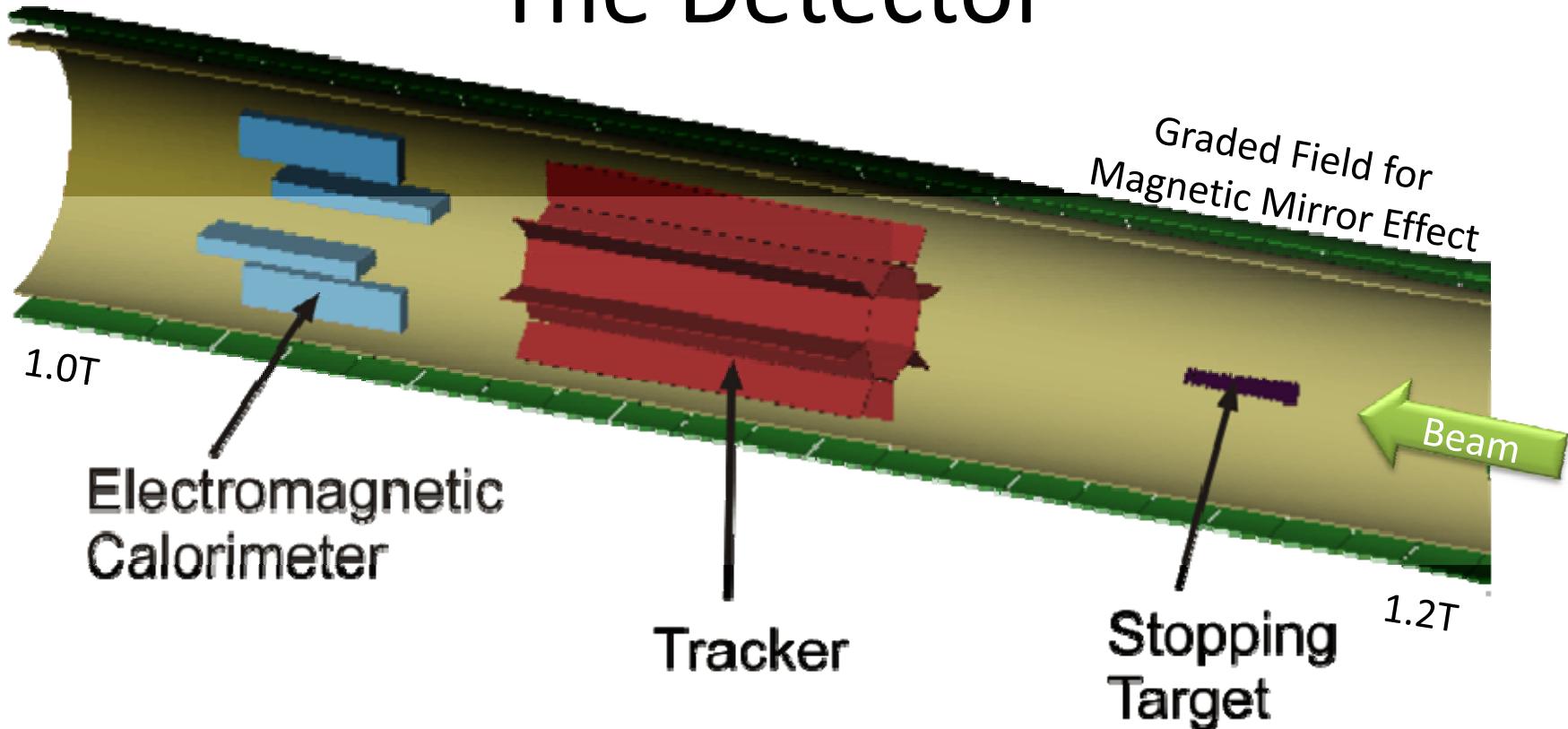


Transport Solenoid

- Designed to sign select the muon beam
 - Colimator blocks the positives after the first bend
 - Negatives are brought back on axis by the second bend
- Gradient along solenoid prevents particles getting “trapped” to reduce long transit time trajectories



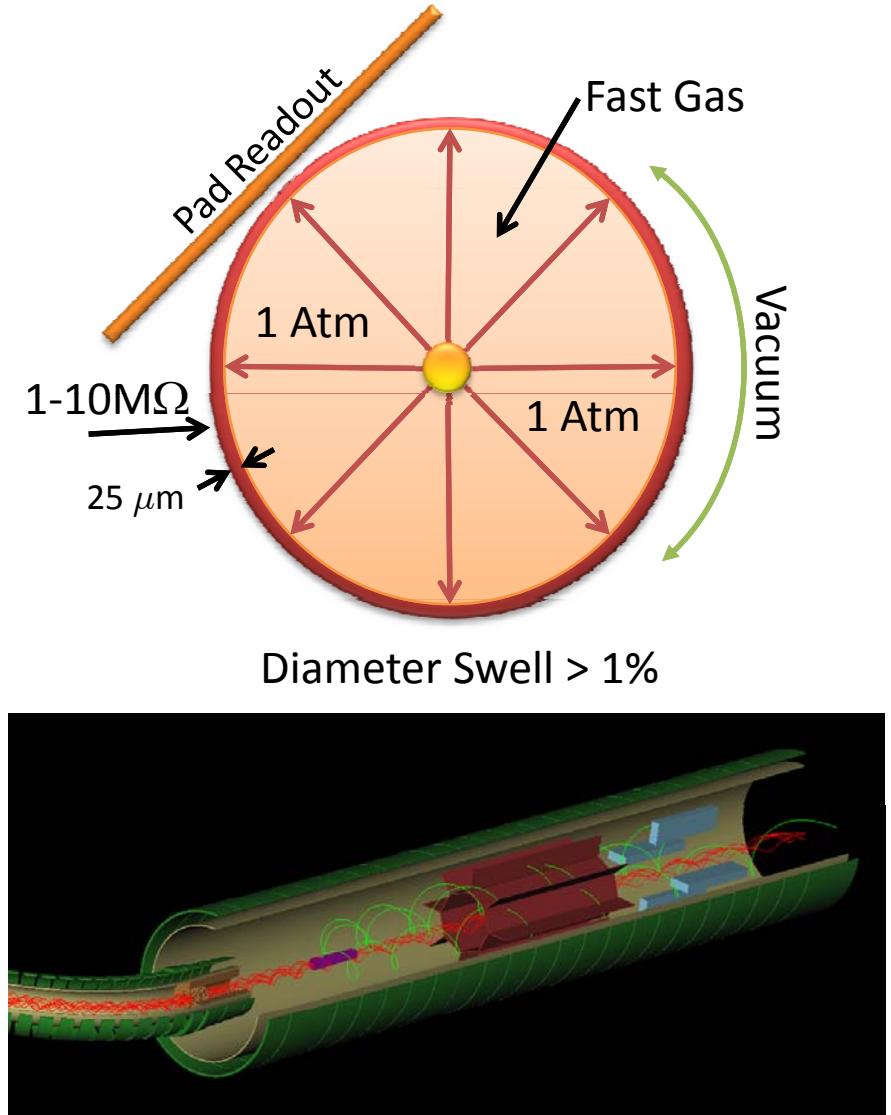
The Detector



- The detector is specifically design to look for the helical trajectories of 105 MeV electrons
- Each component is optimized for maximizing signal detection and suppressing *Decay in Orbit* Backgrounds

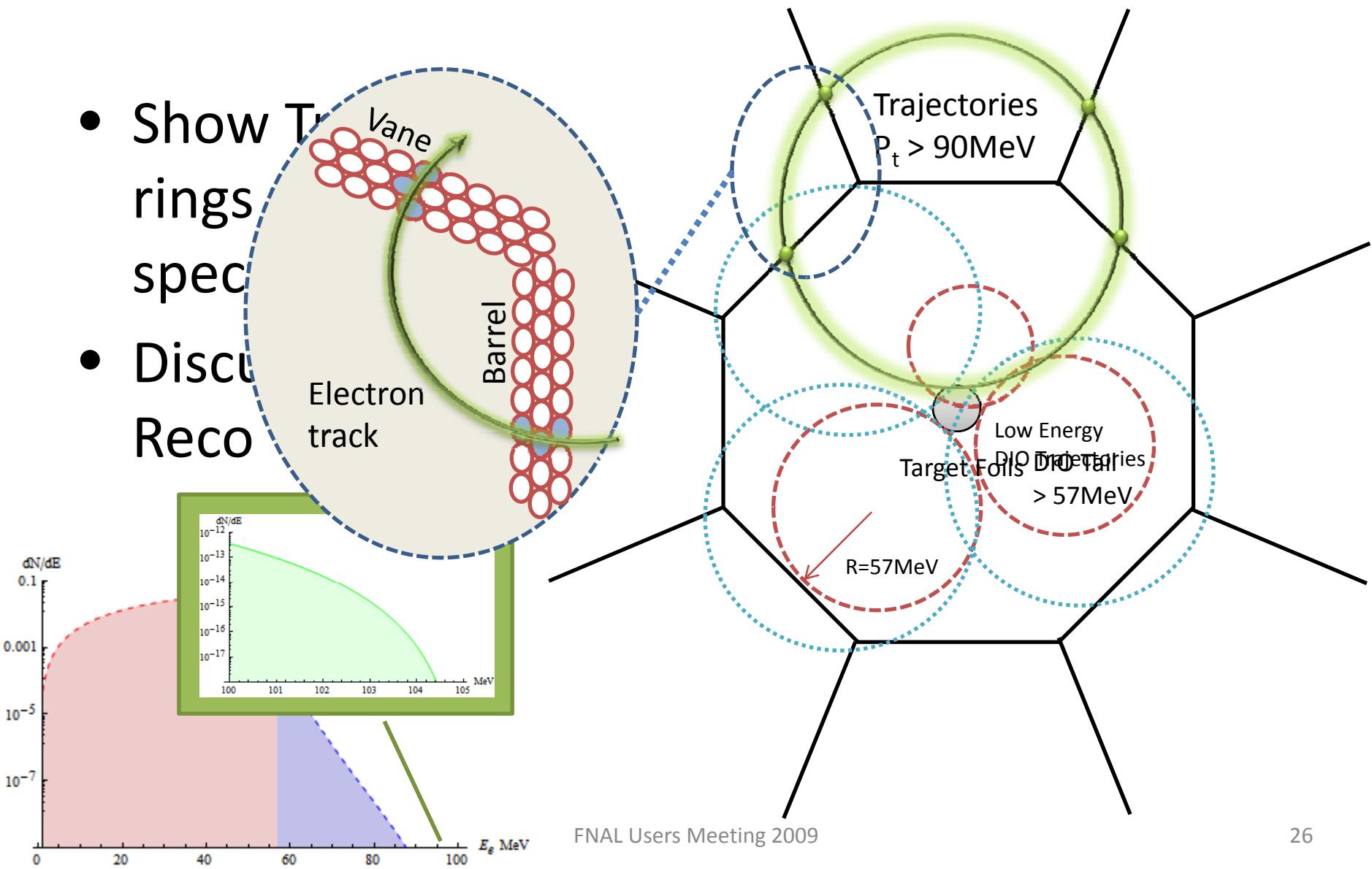
Straw Tracker

- Longitudinal Tracker
Features:
 - 2800 straw tubes in vacuum
 - Utilize 17,000 pad readouts
 - 50% Geometric acceptance to signal ($90^\circ \pm 30^\circ$)
 - Intrinsic resolution 200keV
 - Virtually Immune to DIO



Straw Tracker

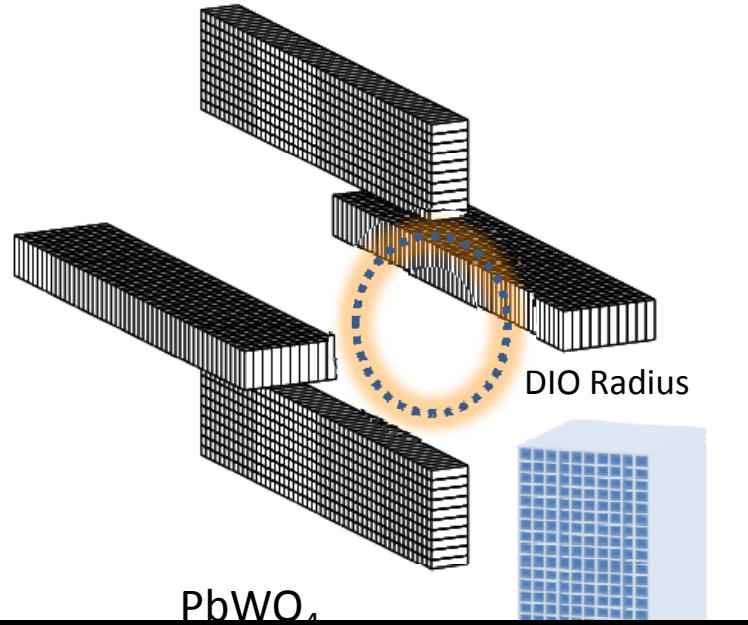
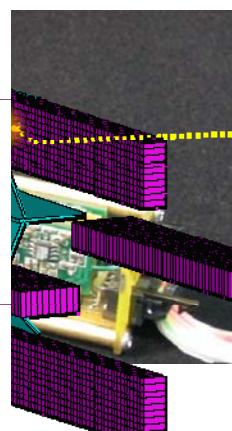
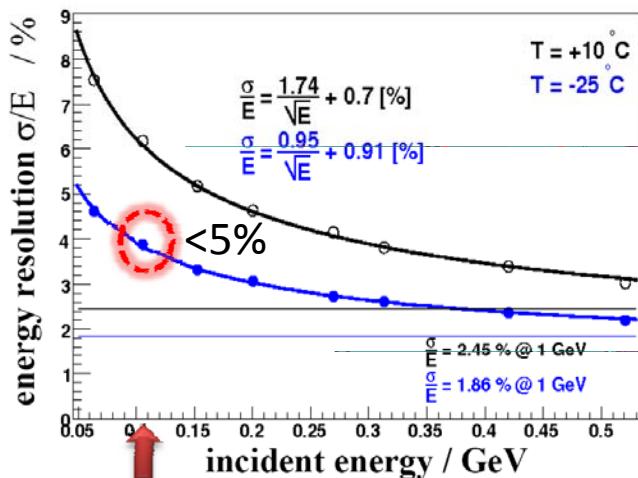
- Show Tr
rings
spec
- Discu
Reco



Crystal Calorimeter

Original Design:

- 5% energy measure for trigger decision (1Hz rate)
- Timing edge for event reconstruction
- Provide PID verification (E/P)
- Spatial match to tracker trajectory
- Immune to DIO rates



PbWO ₄ Calorimeter Properties	
Resolution	5%
Material	$13.6 \times$ PbWO ₄
Readout	Dual APD
Blocks	500 per fin, 4 fins
Segmentation	$30 \times 30 \times 120 \text{ mm}^3$
Trigger Rate	1kHz
Light yield	20-30 p.e./MeV

Cost and Schedule

- Total Project Cost Est. \$200M (fully loaded, escalated, appropriate contingencies)
- Received Stage-1 Approval and CD-0 anticipated shortly
- Technically Driven Schedule (magnet wholly magnet driven) results in 2016 start of data taking
- Signification R&D, Auxiliary Measurements and Test Beam work proceed in parallel to construction

Mu2e Experiment Technically Driven Schedule									
2009		2010		2011		2012		2013	
Q1	Q2	Q3	Q4	Q1	Q2	Q3	Q4	Q1	Q2
Mag. CDR									
Final Magnet Design									
PSI Test Beam		Magnet Construction, Installation, Commissioning							
First Physics Run 2E-17 SES									

Part II

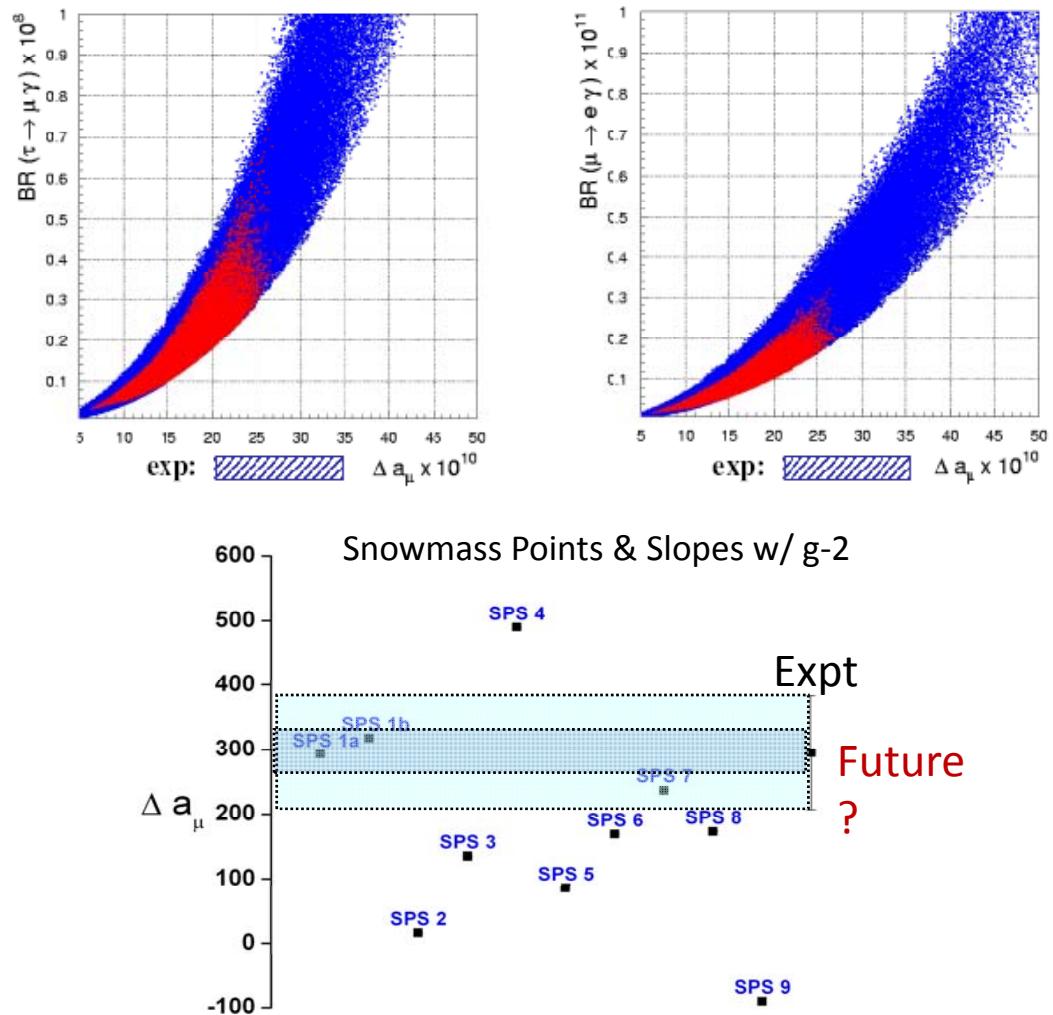
G-2 AT FNAL

Intro & Theory

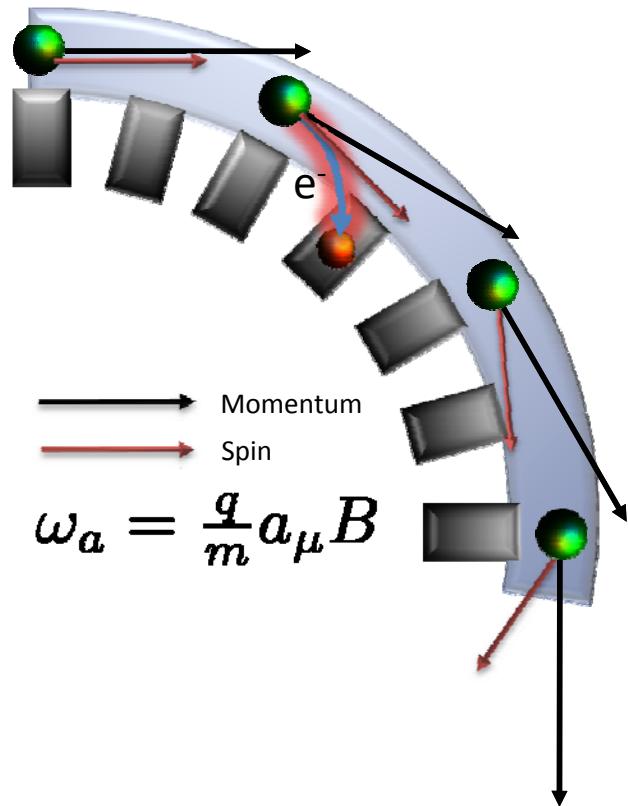
- I owe you a good theory slide

g-2 and SUSY

- g-2 is extremely sensitive to SUSY through the same dipole like interactions that cLFV channels are
- These are amplified by $\tan\beta$ and can be used with LFV results to constrain many models

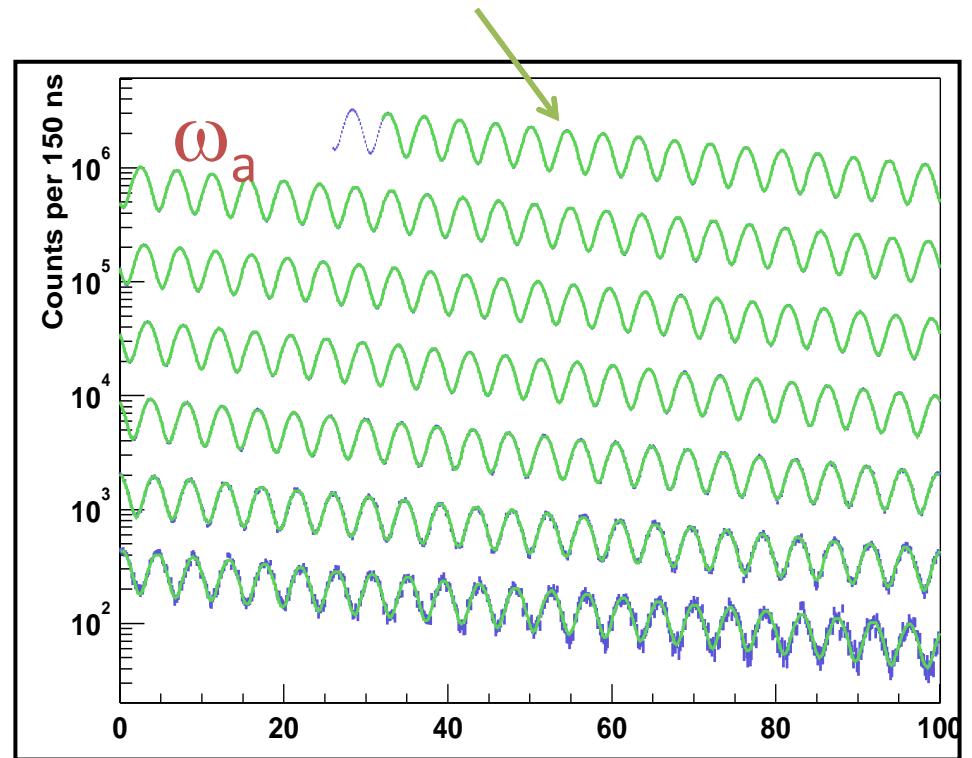


The g-2 Measurement



This method requires extremely precise knowledge of the B field

The muon is self analyzing, and the precession frequency is directly obtained



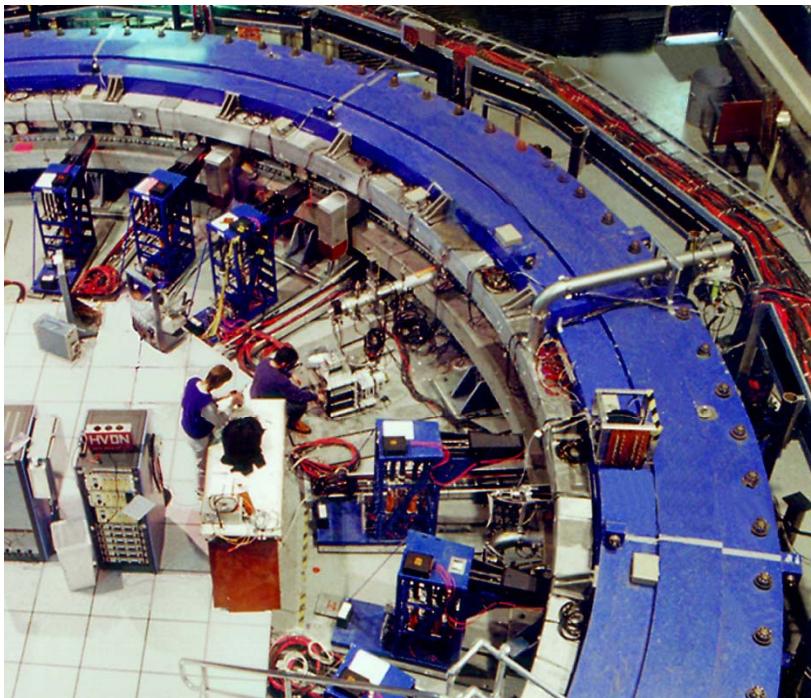
FNAL Beam



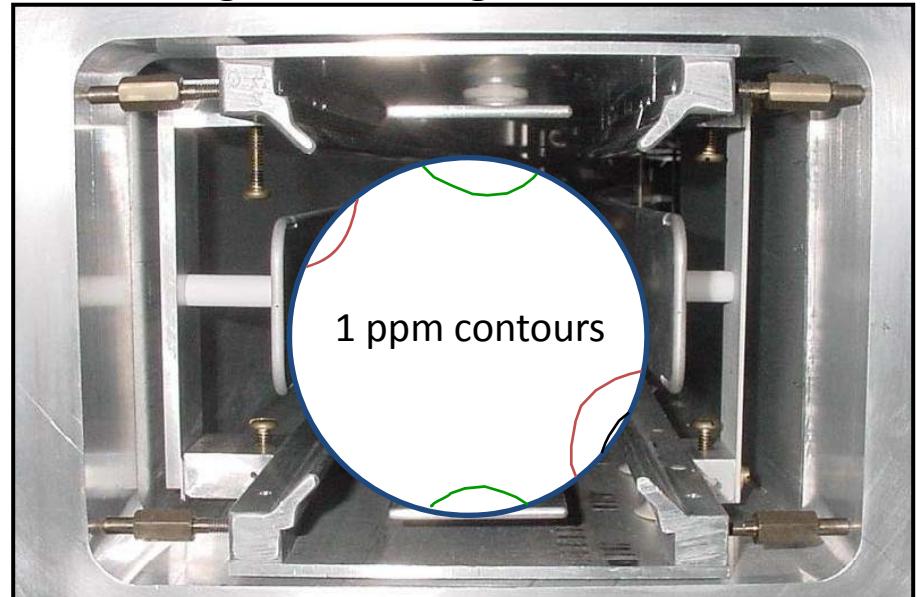
10
High Rep Rate
84 fills/1.4s → 60Hz → $14.5 \times$ BNL
10 20 the statistics in one year of running!

upgrade

- BNL Storage Ring Flies to Chicago!



- Magnetic Field is improved through shimming and calibration

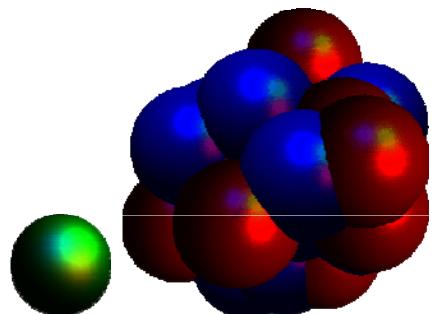


- FNAL beams offer more μ

BACKUP SLIDES

Coherent $\mu^- N \rightarrow e^- N$

- Initial State is a target nucleous N
- First Stop a low energy μ^- beam in the material
- The muon sees the nucleous and cascades down to the 1S orbital with a Bohr Radius and nuclear size: $\langle r_\mu \rangle = \frac{n^2 \hbar}{m_\mu z e^2} \approx 19.6 \text{ fm}$ (for Al)
 $R \approx 1.2 A^{1/3} \text{ fm} = 3.6 \text{ fm}$ (for Al)



Mu2e Collaboration

- 17 Institutions

- Other

<i>Boston University</i> J.Miller, R.Carey, K.Lynch, B. L.Roberts	E.Prebys, R.Ray, V.Rusu, P.Shanahan, M.Syphers, H.White,B.Tschirhart, K.Yonehara,C.Yoshikawa	<i>Northwestern University</i> A.De Gouvea
<i>Brookhaven National Laboratory</i> P.Yamin, W.Marciano, Y.Semertzidis	<i>Idaho State University</i> K.Keeter, E.Tatar	<i>Instituto Nazionale di Fisica Nucleare Pisa, Universita Di Pisa, Pisa, Italy</i> L.Ristori, R.Carosi, F.Cervelli, T.Lomtadze, M.Incagli, F.Scuri, C.Vannini
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<i>University of California, Irvine</i> W.Molzon	<i>Institute for Nuclear Research, Moscow, Russia</i> V.Lobashev	<i>Syracuse University</i> P.Souder, R.Holmes
<i>City University of New York</i> J.Popp	<i>University of Massachusetts, Amherst</i> K.Kumar, D.Kawall	<i>University of Virginia</i> E.C.Dukes, M.Bychkov, E.Frlez, R.Hirosky, A.Norman, K.Paschke, D.Pocanic
<i>Fermi National Accelerator Laboratory</i> C.Ankenbrandt, R.Bernstein, D.Bogert, S.Brice, D.Broemmelsiek, R.Coleman, D.DeJongh, S.Geer, D.Glenszinski, D.Johnson, R.Kutschke, M.Lamm, P.Limon, M.Martens, S.Nagaitsev, D.Neuffer, M.Popovic,	<i>Muons, Inc.</i> T.Roberts, R.Abrams, M.Cummings R.Johnson, S.Kahn, S.Korenev, R.Sah	<i>College of William and Mary</i> J.Kane